

Energy, density matrices and reduced density matrices in molecular systems: Grand-canonical electron states

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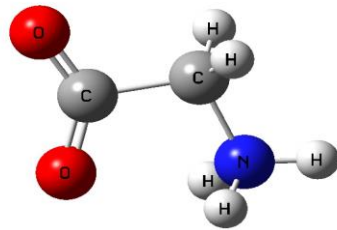
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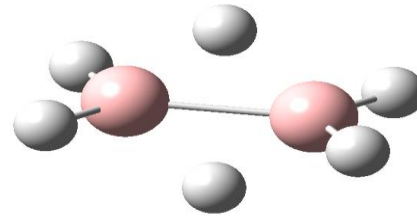
What is a molecule?

Views

- Chemical view \approx set of atoms and bonds



Zwitterionic glycine



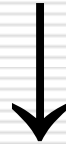
Diborane B₂H₆

- Physical view \approx set of interacting particles: nuclei and electrons

Conciliate both views

Starting point: solving the
Schrödinger equation

Many-body problem
(really few)



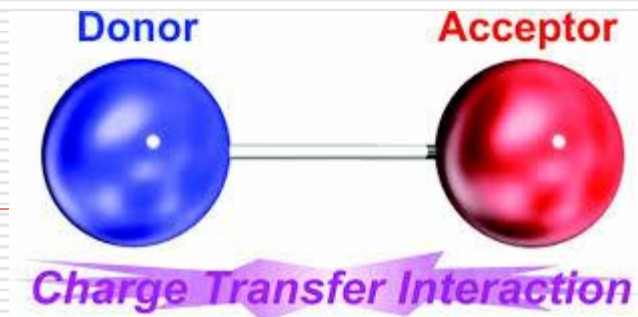
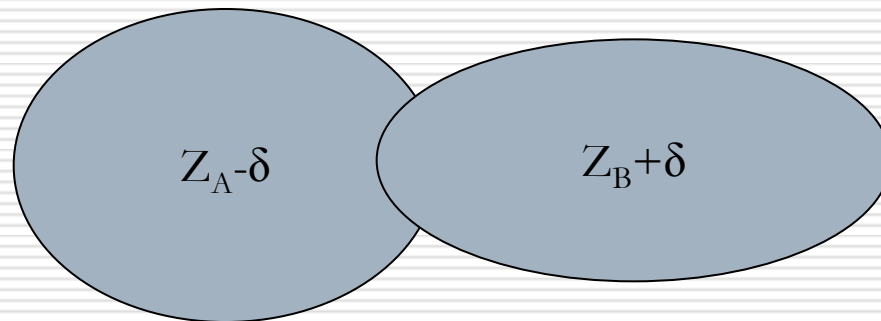
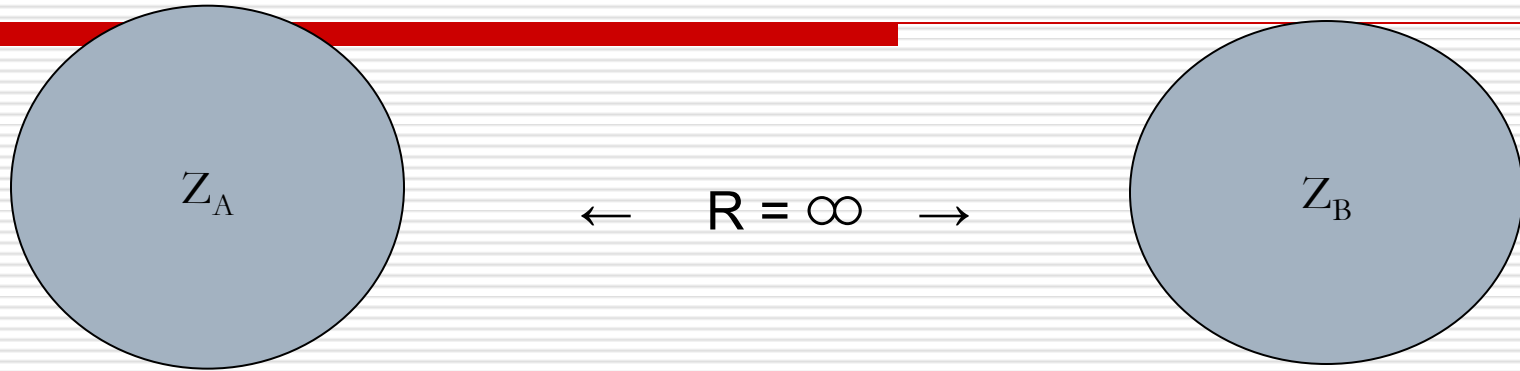
$$H\Psi = E\Psi$$

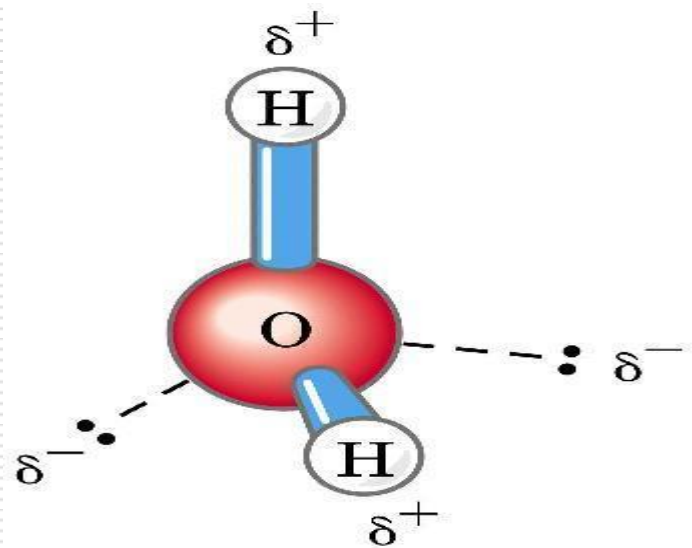
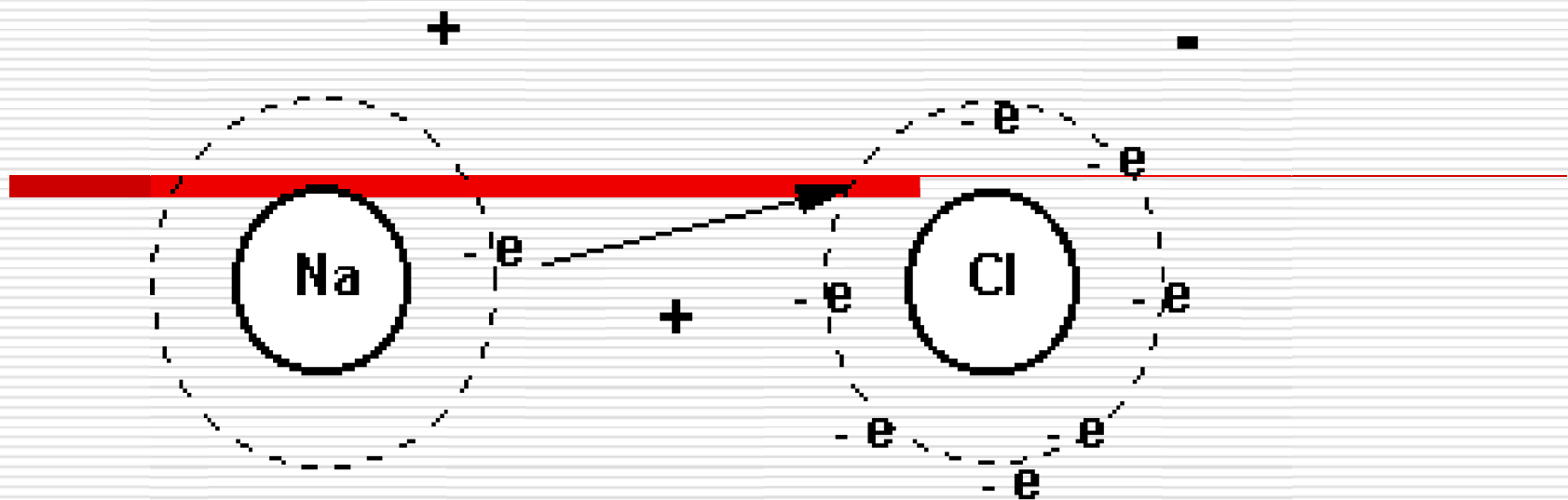
**To extract the information contained within
the state function**

Schrödinger equation at the Born-Oppenheimer approach

$$\begin{aligned}
 \hat{H} = & \cancel{-\frac{\hbar^2}{8\pi^2} \sum_A^{\text{nuclei}} \frac{1}{M_A} \nabla_A^2} - \frac{\hbar^2}{8\pi^2 m} \sum_a^{\text{electrons}} \nabla_a^2 - e^2 \sum_A^{\text{nuclei}} \sum_a^{\text{electrons}} \frac{Z_A}{r_{Aa}} \\
 & + \cancel{e^2 \sum_{A > B}^{\text{nuclei}} \sum_{B} \frac{Z_A Z_B}{r_{AB}}} + e^2 \sum_a^{\text{electrons}} \sum_{b > a} \frac{1}{r_{ab}} \\
 = & \sum_a h(a) + \sum_{a,b} v(a,b)
 \end{aligned}$$

Todo proceso fisicoquímico a nivel microscópico se caracteriza por una **TRANSFERENCIA de CARGA**





Chemical descriptors

Empirical ideas

- Covalent bond orders
(shared electrons)
- Covalency – ionicity
(complementary)
- Atomic charges
(concept of particles)
- Valence
- Free valence
(concept of holes)

associated with the electron density
(population analysis)



partition of the number of electrons in
the system, **N** (an integer number)

Among others ... the list is not complete

Other fundamental descriptors

- Reactivity (vacant density sites)
- Electronegativity
- Chemical potential
- Others

associated with energy and electron density derivatives



a non integer number of particles, \mathcal{N}

Problems: energy and electron distribution

- System
 - States
 - Energy determination
-

The system

Se necesita una definición y su caracterización de lo que sea un **Átomo en una Molécula (AIM)**

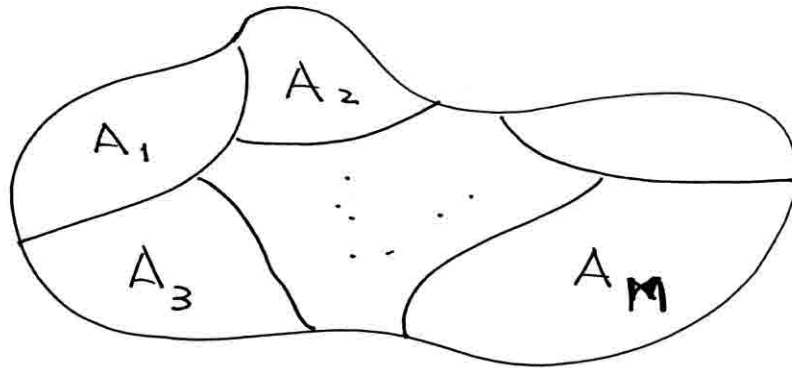
Formulaciones

- 1. Mulliken (base de orbitales atómicos, espacio de Hilbert)*
 - 2. Modelos "Fuzzy": dominios en el espacio físico (regiones) de soporte empírico*
 - 3. La aproximación Lagrangiana para la definición de los dominios (Bader).
Define una topología en forma rigurosa.*
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ATOMS IN MOLECULES

(arbitrary definitions)

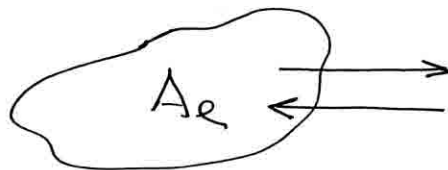
\mathcal{M}
(molecule)



• A_1, A_2, \dots, A_M
(DOMAINS)
PHYSICAL SPACE

• $\mathcal{M} = \bigcup_{k=1}^M A_k$ $A_i \cap A_\ell = \begin{cases} \phi & \text{(empty)} \\ \neq \phi \end{cases}$

• each $\{A_\ell\}$ is an OPEN SYSTEM



$$N = N + \nu$$

$$N \in \mathbb{Z}^+; \quad 0 \leq \nu \leq 1$$

The states

DENSITY MATRICES (DM)

$${}^N D_{\Phi_k^N} = |\Phi_k^N\rangle\langle\Phi_k^N|$$

$$D = \sum_M \sum_{\Phi_k^M} \omega_{\Phi_k^M} |\Phi_k^M\rangle\langle\Phi_k^M|; \quad \sum_M \sum_{\Phi_k^M} \omega_{\Phi_k^M} = 1; \quad \omega_{\Phi_k^M} \geq 0$$

$$D = \bigoplus_{\{\Phi_k^M\}}^{\infty} \omega_{\Phi_k^M} {}^M D_{\Phi_k^M}$$

**Grand-canonical
state**

- finite trace, $tr(\mathcal{D}) < \infty$; normalization: $tr(\mathcal{D}) = 1$
- $\mathcal{D} = \mathcal{D}^\dagger$; hermiticity
- $\mathcal{D} \geq 0$; positive semi-definite

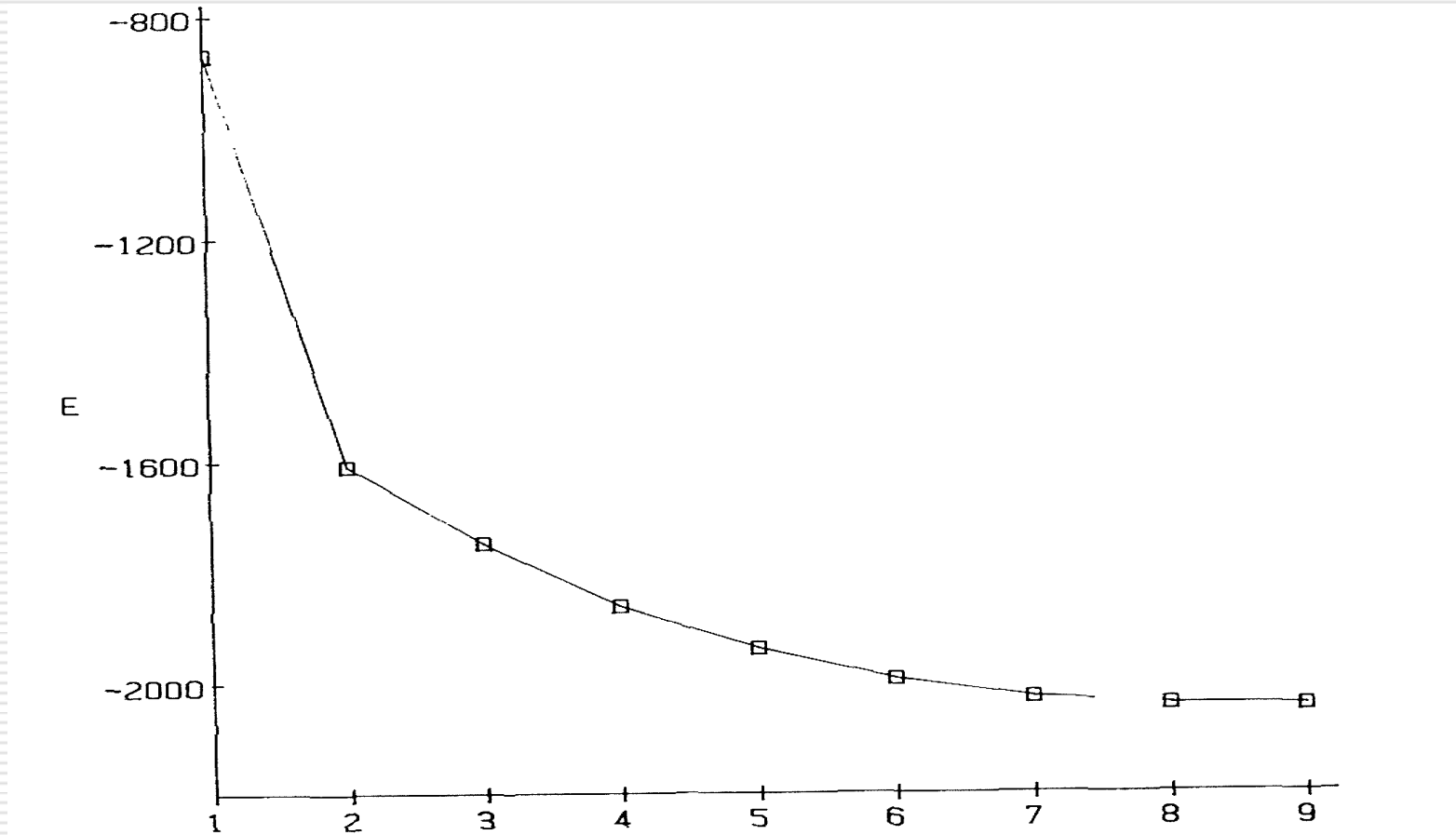
Probabilistic interpretation

$$\int dx_1 \int dx_2 \dots \int dx_N {}^N \mathcal{D} (x_1, x_2, \dots, x_N | x_1, x_2, \dots, x_N) = 1$$

System energy

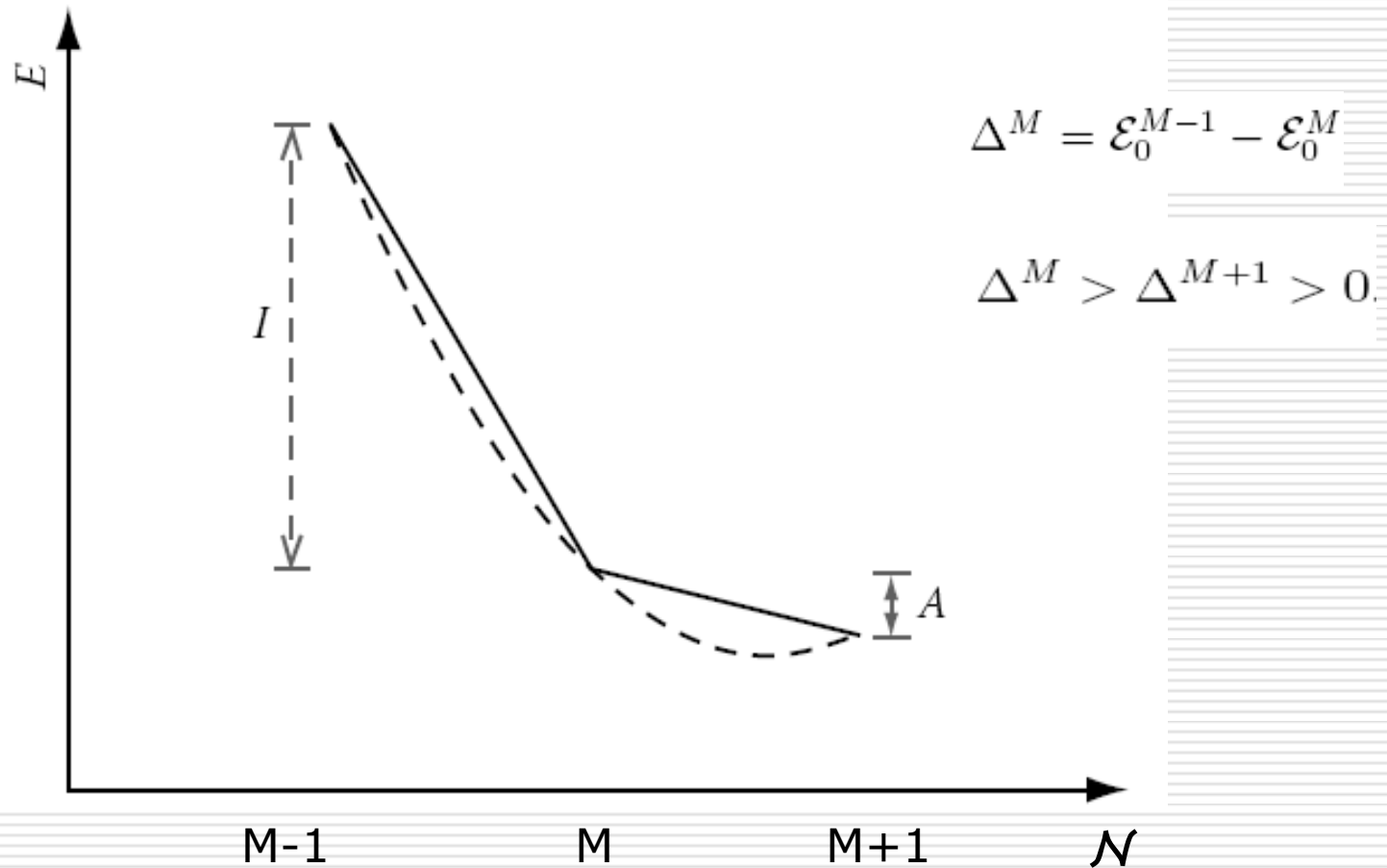
$$\mathcal{E} = \text{Tr}(D \mathcal{H}) = \sum_{\{\Phi_k^M\}} \omega_{\Phi_k^M} \text{Tr} \left({}^M D_{\Phi_k^M} \mathcal{H} \right)$$

The energy of an atomic or molecular system as a function of the number of electrons (arbitrary units)
Experimentally and numerically



Example: Oxygen atom and its ions

- *The energy is convex (Hypothesis)*



***Then it follows,
(Lemma and Proposition (de la Oveja Rial))***

Lemma 1. *For the sequence $\{\mathcal{E}_0^M\}_{M \in \mathbb{N}}$ verifying the above inequality, and arbitrary $N, M \in \mathbb{N}$ numbers, such that $M \neq N, N + 1$, then*

$$\mathcal{E}_0^M \geq (N + 1 - M) \mathcal{E}_0^N + (M - N) \mathcal{E}_0^{N+1}$$

The equality holds iff $M = N, N + 1$.

mathematical basis of the energy convexity

The ground state energy for the system with non integer number of particles is obtained from

$$\left\{ \begin{array}{l} \min \mathcal{E} = \sum_{\{\Phi_k^M\}} \omega_{\Phi_k^M} \text{Tr} \left({}^M D_{\Phi_k^M} \mathcal{H} \right) \\ \sum_{\{\Phi_k^M\}}^{\infty} \omega_{\Phi_k^M} M = N + \nu \end{array} \right.$$

With solution

Proposition 1. *The solution of this problem, i.e., $D = \sum_M \sum_{\Phi_k^M} \omega_{\Phi_k^M} D_{\Phi_k^M}$ where the variational parameters are the statistical weights $\{\Phi_k^M\}$, is unique and become expressed by*

$$D = (1 - \nu) {}^N D_0 + \nu {}^{N+1} D_0$$

A convex combination of only two pure state distributions: N and $(N+1)$ - particles, leading to the energy

$$\mathcal{E}_0^{N+\nu} = (1 - \nu) \mathcal{E}_0^N + \nu \mathcal{E}_0^{N+1}.$$

Las propiedades del sistema dependen exclusivamente de las densidades de partículas y de pares debido a que los correspondientes operadores son de 1- y 2-partículas.



Por ello son necesarias las distribuciones marginales de la distribución D

p-order Reduced Density Matrices (p-RDM)

- Marginal distributions
- Contraction mapping from the DM

p-reduced density matrix (p-RDM)

$$\overbrace{\Gamma(\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_p | \mathbf{x}'_1, \mathbf{x}'_2, \dots, \mathbf{x}'_p)} =$$

contraction mapping → average → marginal distribution

$$\overbrace{\int d\mathbf{x}_{p+1} \dots \int d\mathbf{x}_N} {}^N \mathcal{D}(\mathbf{x}_1, \dots, \mathbf{x}_p \mathbf{x}_{p+1}, \dots, \mathbf{x}_N | \mathbf{x}'_1, \dots, \\ \dots, \mathbf{x}'_p \mathbf{x}_{p+1}, \dots, \mathbf{x}_N)$$

$$\equiv \text{tr}_{p+1, \dots, N}({}^N \mathcal{D})$$

- finite trace, $tr({}^pD) < \infty$; normalization: $tr({}^pD) = \binom{N}{p}$
- ${}^pD = {}^pD^\dagger$; hermiticity
- ${}^pD \geq 0$; positive semi-definite
- N-representability

$$\text{Physical meaning} \rightarrow {}^pD = \binom{N}{p} {}^p\Gamma; \quad tr({}^pD) = \binom{N}{p}$$

$$\downarrow$$

Densities: populations (diagonal elements) and coherences (off-diagonal)

$$p = 1, 2, \dots,$$

$${}^1D = N {}^1\Gamma; \quad tr({}^1D) = N$$

$${}^2D = \binom{N}{2} {}^2\Gamma; \quad tr({}^2D) = \binom{N}{2}$$

Contraction mapping in Fock space

$${}^p D = \hat{L}_p\{D\} = \bigoplus_{\{\Phi_k^M, M \geq p\}} \omega_{\Phi_k^M} \binom{M}{p} \hat{L}_M^p\{{}^M D_{\Phi_k^M}\}$$

$$\hat{L}_p^p\{{}^p D_{\Phi_k^p}\} = \mathcal{I}$$

$$\hat{L}_M^p\{{}^M D_{\Phi_k^M}\} = \mathbf{O} \text{ for } M < p,$$

Identity superoperator

Null superoperator

Consequences

- **Energy functionals**

$$\mathcal{E}_0^{N+\nu} = \mathcal{F}({}^2D_o^{N+1}, {}^2D_o^N, \nu)$$

not

$${}^2D = (1 - \nu) {}^2D_o^N + \nu {}^2D_o^{N+1}$$

- **Derivatives**

1. Chemical potential

$$\mu = \left(\frac{\partial \mathcal{E}_0^N}{\partial N} \right)_v$$

$$\mu_{\pm} = \pm \left(\frac{\partial \mathcal{E}_0^N}{\partial \nu} \right)_v = \pm \left(\mathcal{E}_0^{N \pm 1} - \mathcal{E}_0^N \right)$$

$$\mu_+ = \mathcal{E}_0^{N+1} - \mathcal{E}_0^N = -EA$$

$$\mu_- = \mathcal{E}_0^N - \mathcal{E}_0^{N-1} = -IP$$

Average ν electrons are removed or attached to the system

$$IP^{\nu} = \nu IP$$

$$EA^{\nu} = \nu EA$$

2. Local descriptors: Fukui functions

$$f^\pm(\mathbf{r}) = \left(\frac{\partial \rho(\mathbf{r})}{\partial \mathcal{N}} \right)_v^\pm$$

$$F^\pm(\mathbf{r}|\mathbf{r}') = \pm \left(\frac{\partial}{\partial \omega} {}^1D^{N\pm\omega}(\mathbf{r}|\mathbf{r}') \right)_v$$

$${}^1D^{N\pm\nu} = \nu {}^1D^{N\pm 1} + (1 - \nu) {}^1D^N$$

$$F^\pm(\mathbf{r}|\mathbf{r}') = \pm \left({}^1D^{N\pm 1}(\mathbf{r}|\mathbf{r}') - {}^1D^N(\mathbf{r}|\mathbf{r}') \right)$$

Final remarks

1. GC ideas: few particle open systems in their ground states.
 2. Only two pure states contribute to the electron distribution of $N+\nu$ particle systems. Convex combination of N y $N+1$ states.
 3. Contraction mapping for density matrices: p-RDM
 4. Chemical local and non local descriptors in terms of p-RDMs valid for any type of state function.
 5. Second order derivatives are coming
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