



Transiciones de fase en la corteza de las estrellas de neutrones ?

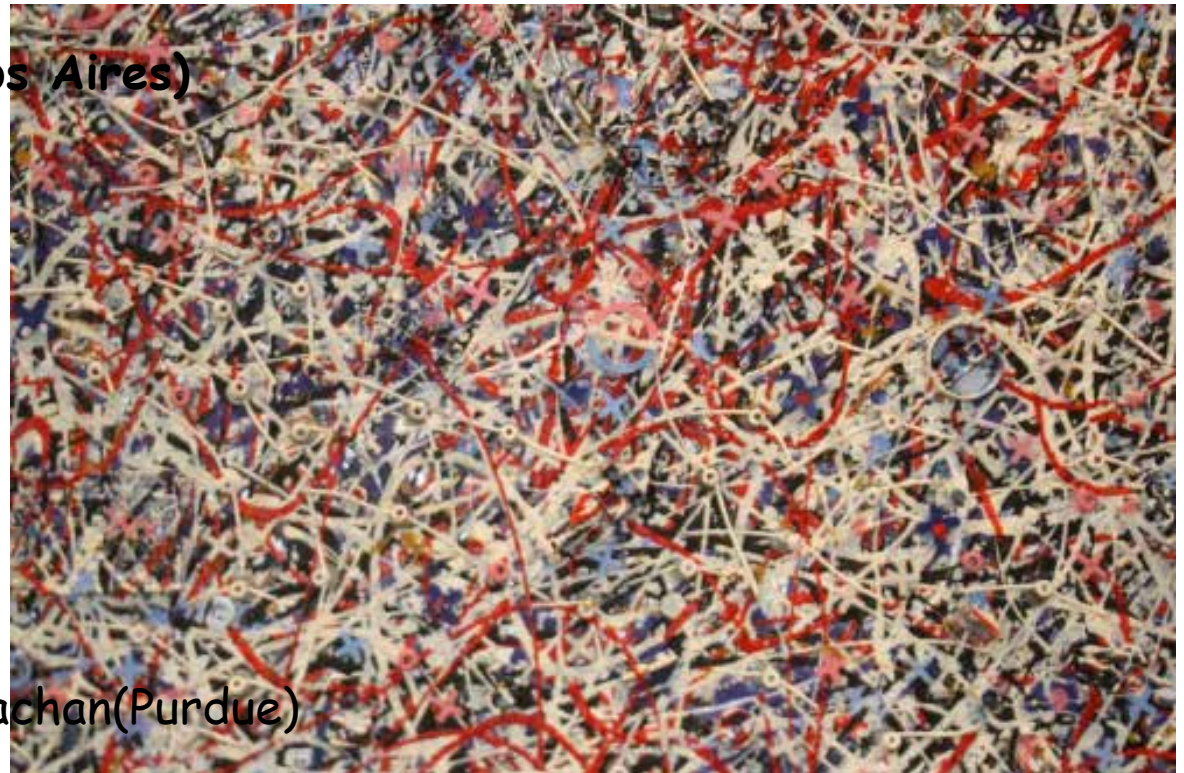
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Astrofísica Nuclear

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SocioFísica

Dinámica Peatonal (pánico)

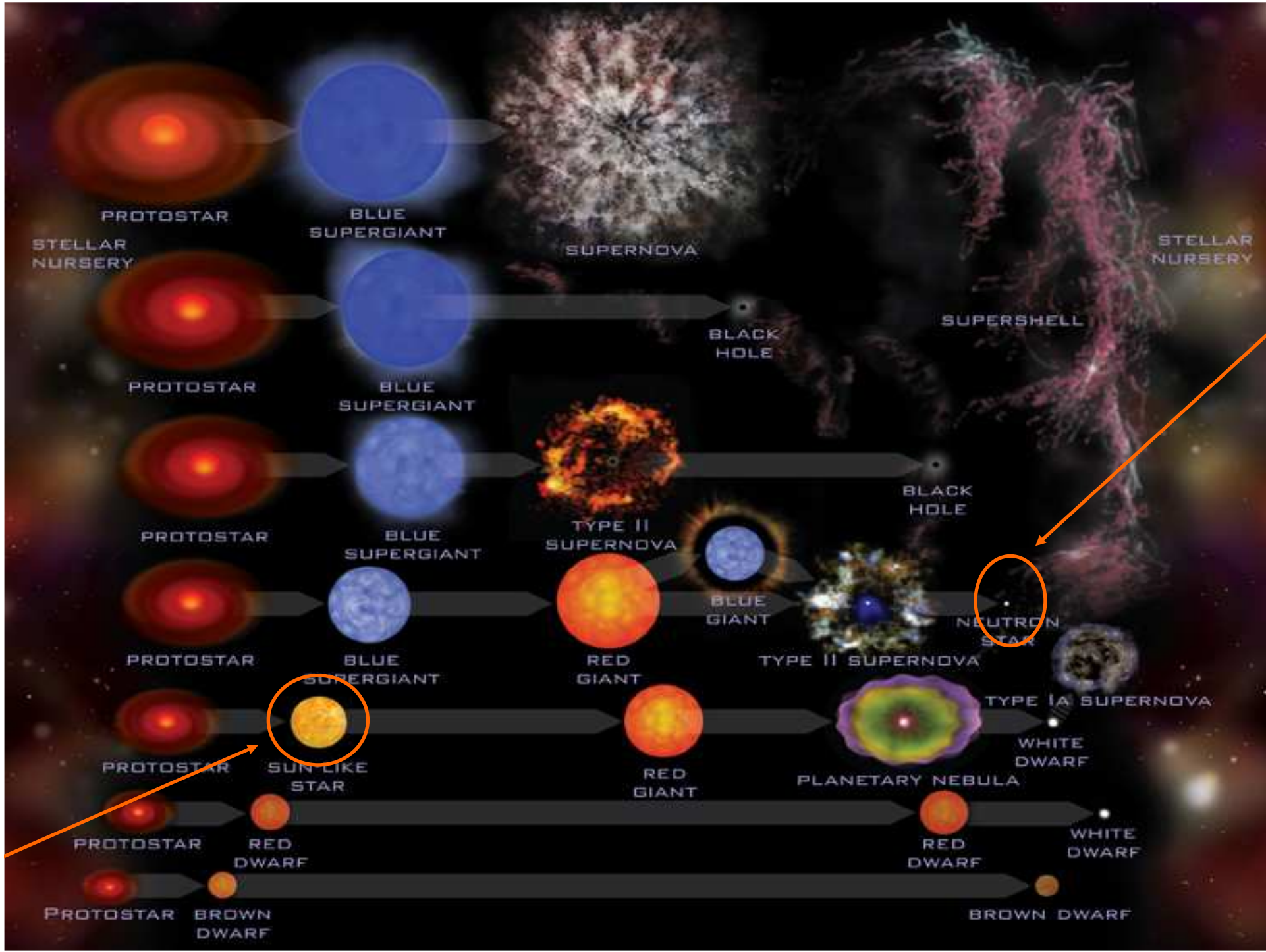
C.O.Dorso
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Complex Networks Formación de Opiniones

C.O.Dorso
P.Balenzuela
S.Pinto
G.Pasqualetti
A. Medus
D. Barmak

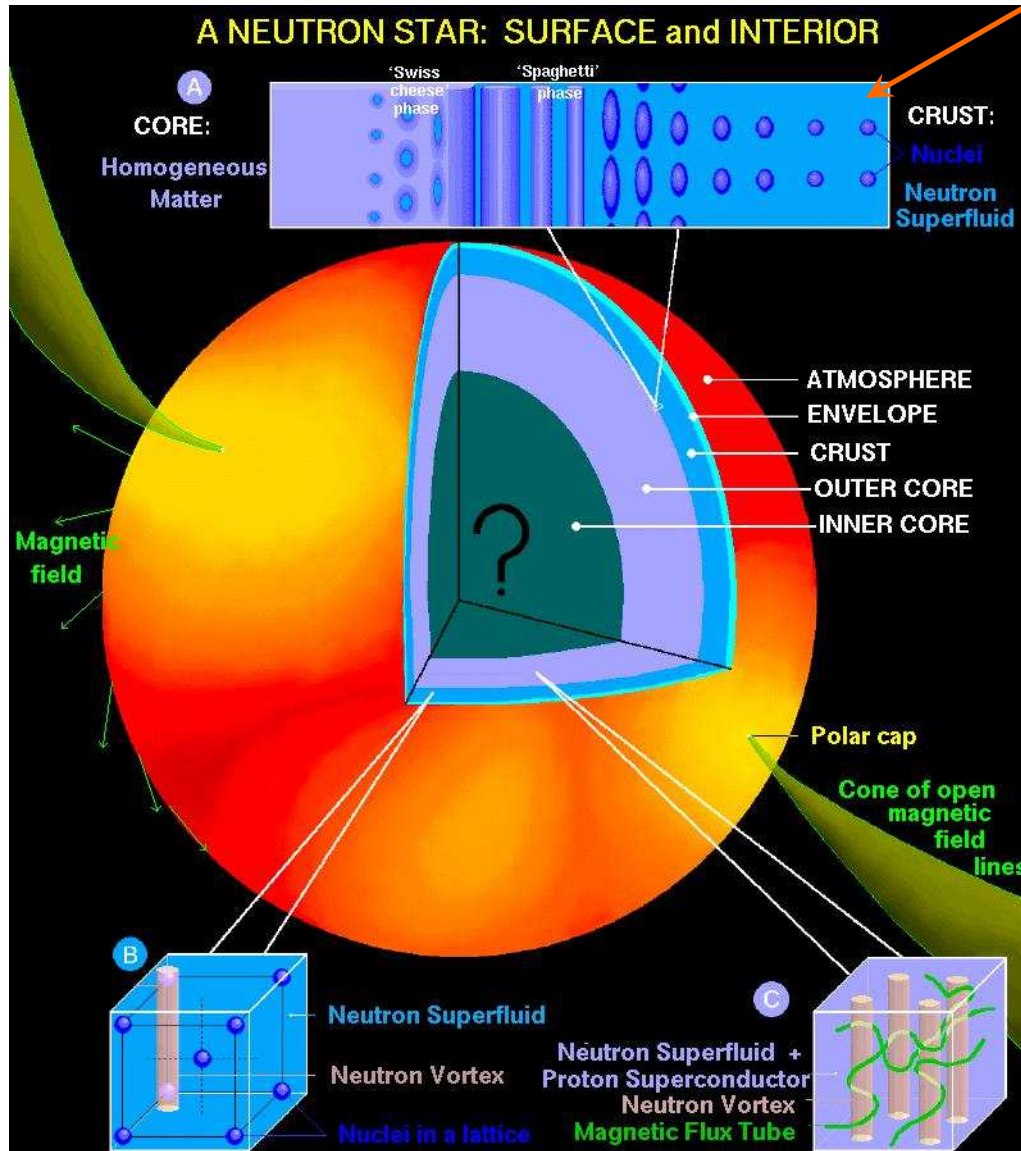
Biología de Sistemas

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Astrofísica Nuclear: Estrellas de Neutrones

(circa 2010)



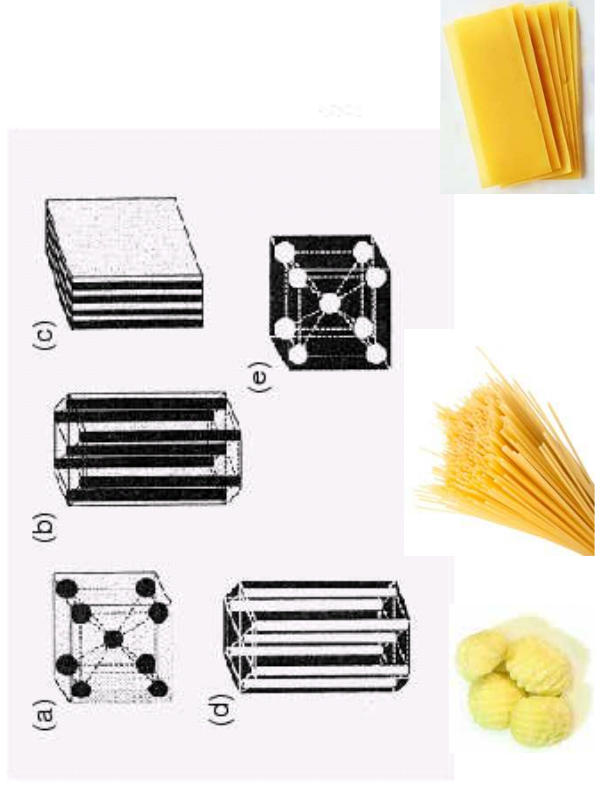
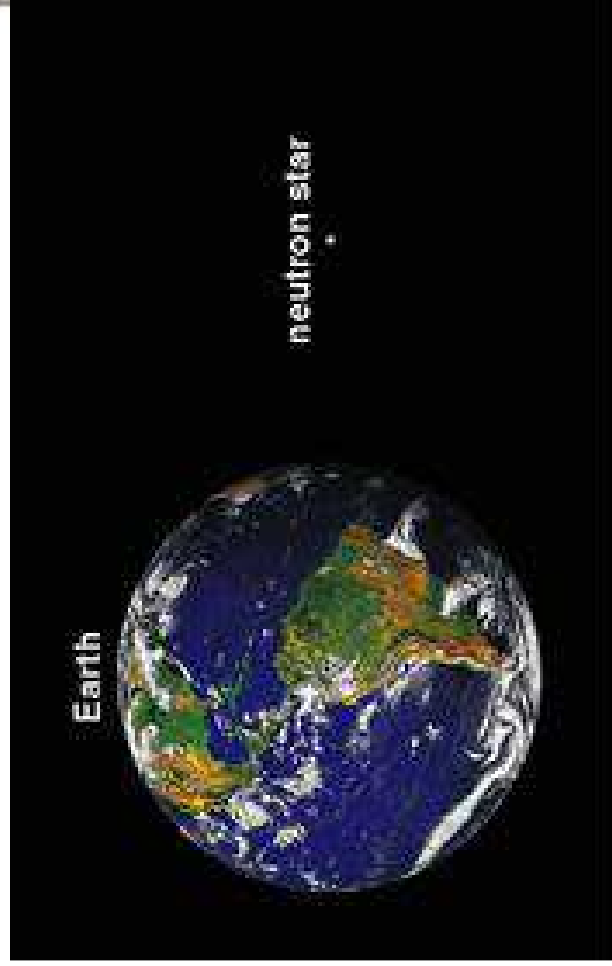
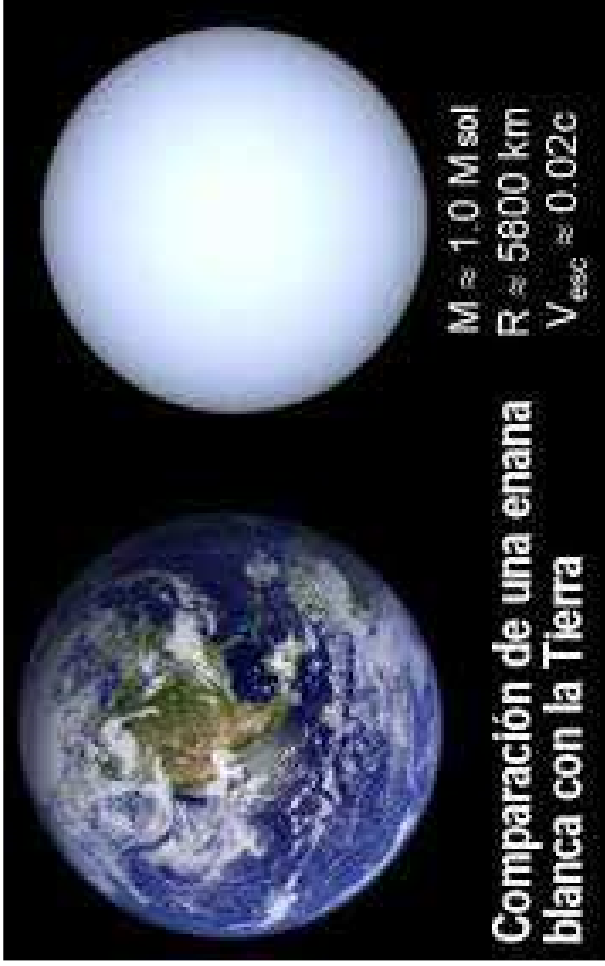
- $M \sim 1.4 M_{\odot} \rightarrow 2 M_{\odot}$
- $R \sim 10 \text{ km}$
- $\rho \sim 10^6 \text{ g/cm}^3 \rightarrow 8 \times 10^{14} \text{ g/cm}^3$
- $x \sim 0,5 \rightarrow 0,2 (?)$
- $T \sim 1 \text{ MeV} (10^{10} \text{ K})$

Laboratorio para NEOS:

$$E(\rho, T, \delta) \sim E(T, \rho, \delta = 0) +$$

$$+ E_{Sym}(T, \rho)\delta^2,$$

$$\delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$$



Algunas preguntas a responder

Que es nuclear Pasta?

Cual es el efecto de las condiciones periódicas de contorno en la formación de pastas?

Rol de coulomb?

Rol del apantallamiento de Debye?

Transiciones de fase En las pastas?

Opacidad de Neutrinos y la topología de la corteza.

Nucleosíntesis (Elementos mas pesados que Hierro)?

Esta Presentación:

Neutron Star

El modelo CMD
(Illinois potential)

Sistemas Finitos
Reconocimiento de fragmentos
Comportamiento critico
Energia se simetria
Isoscaling

**Sistemas infinitos
(NS ∞)**

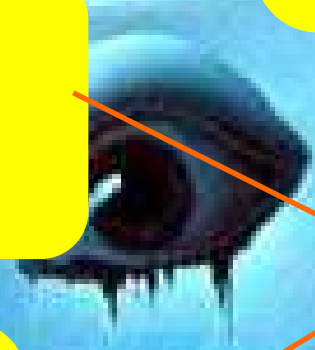
Topologia
Coulomb

Intermezzo

Sistemas 2D y 3D
Lennard Jones + Coulomb

Illinois potential

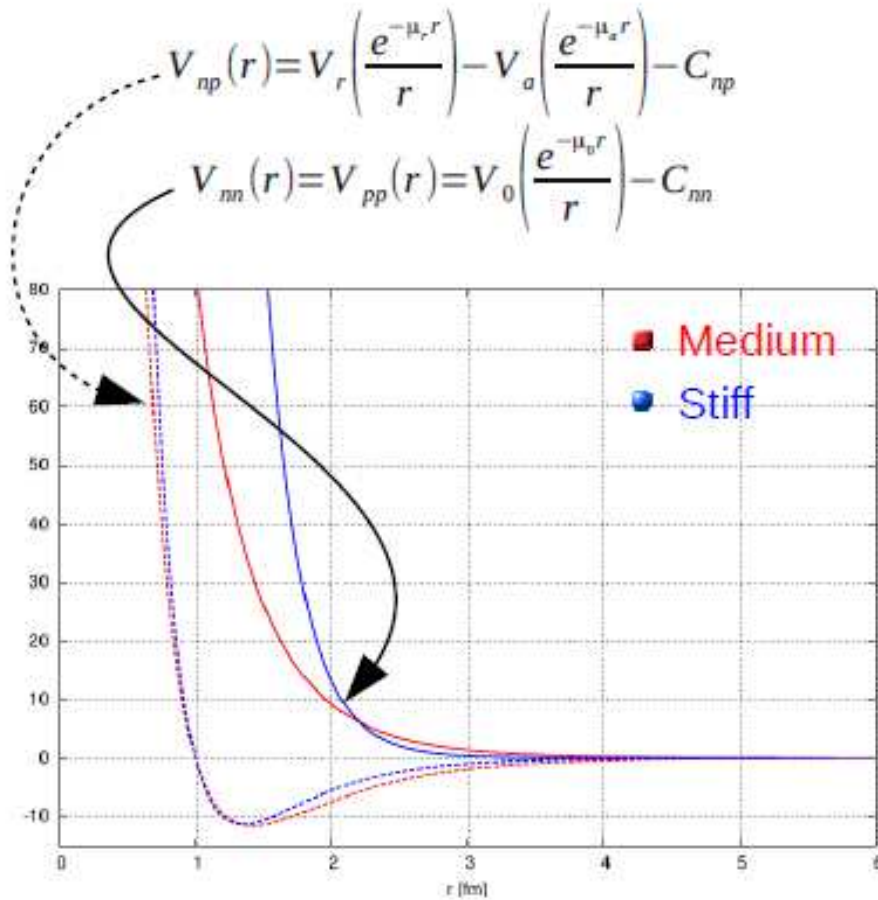
**Materia Nuclear
Materia NS**



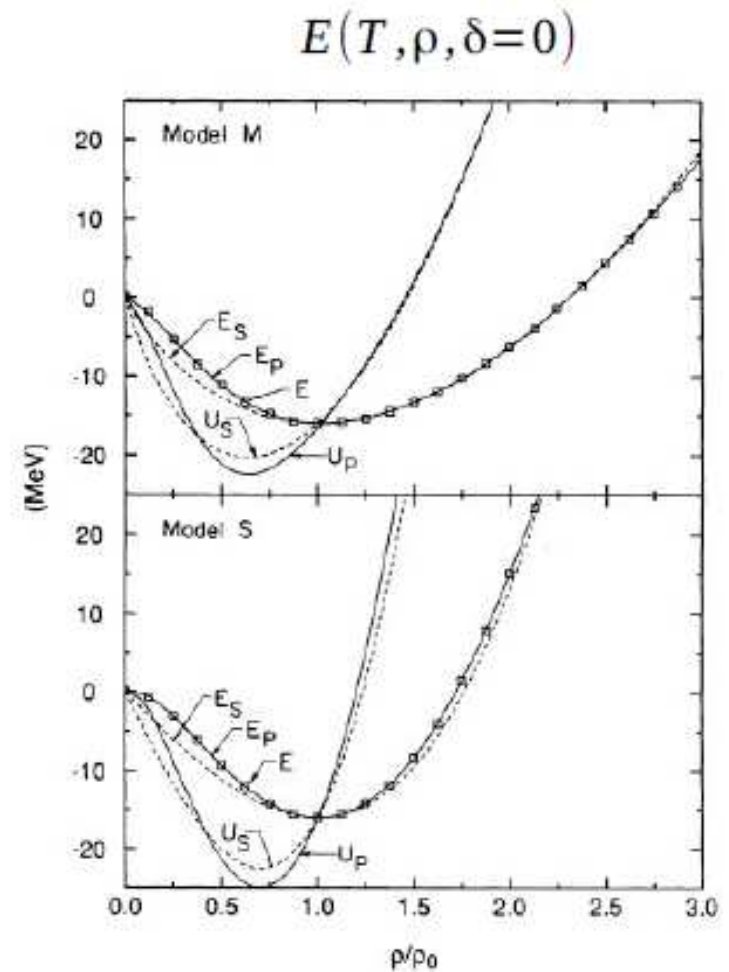
El modelo nuclear CMD
y
los núcleos finitos



Modelo CMD



- Simulaciones de Dinámica Molecular Clásica
- Mecánica Estadística de sistemas 'raros'
- Término de Simetría de la EOS Nuclear



Modelo CMD

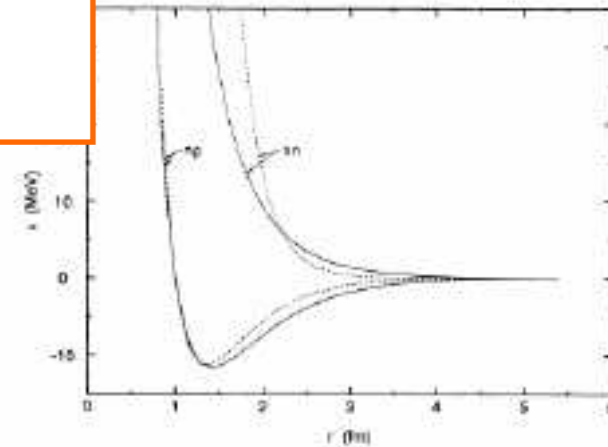
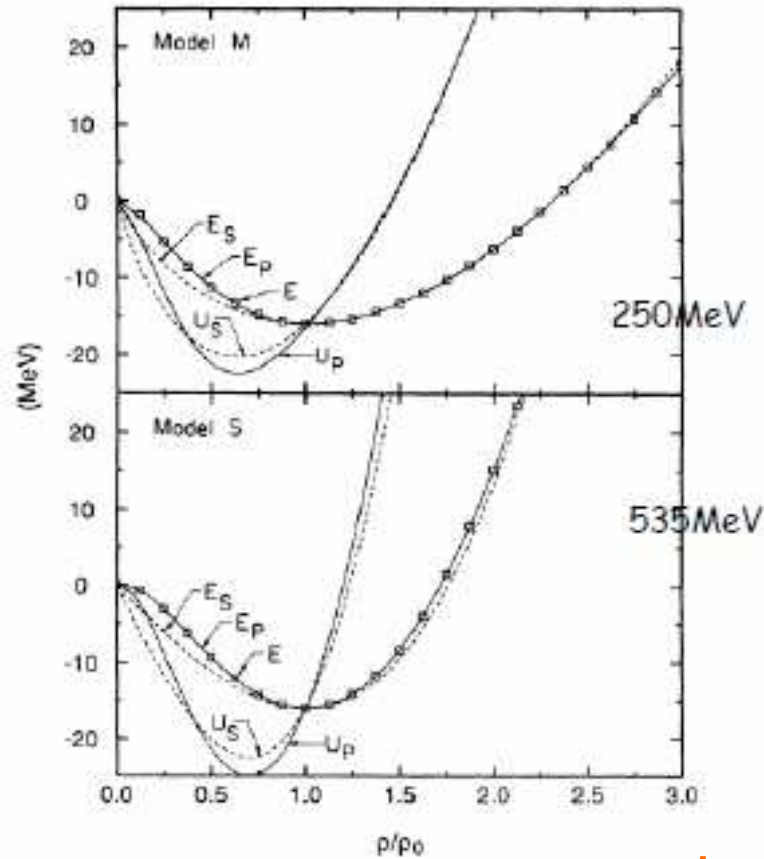


FIG. 1. The interparticle potentials v_m and v_n for the *M* (solid) and *S* (dashed) models.

TABLE II. Energies of nuclei (MeV/nucleon) for the *M* and *S* models compared to experimental binding energies.

<i>A</i>	<i>Z</i>	E_M	E_S	E_{exp}
2	1	-5.76	-5.67	-1.11
3	1	-7.14	-7.26	-2.83
4	2	-7.91	-7.05	-7.07
5	2	-7.47	-6.22	-5.47
16	8	-10.39	-10.49	-7.98
40	20	-10.44	-10.60	-8.55
90	40	-9.93	-10.25	-8.71
139	57	-9.12	-9.58	-8.38
197	79	-8.38	-8.84	-7.92
208	82	-8.16	-8.71	-7.69

PHYSICAL REVIEW C

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Publicacion original

Accuracy of the Vlasov-Nordheim approximation in the classical limit

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Department of Physics, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801

(Received 1 March 1990)

Modelo CMD

Resolver las ecuaciones de movimiento
(symplectic)

Reconocer Clusters
(MST, MSTE, ECRA)

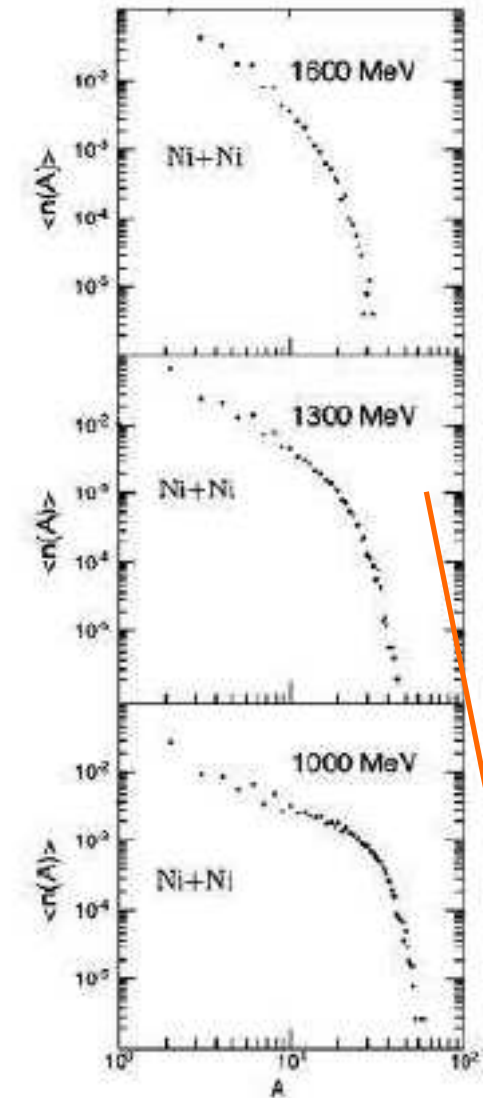
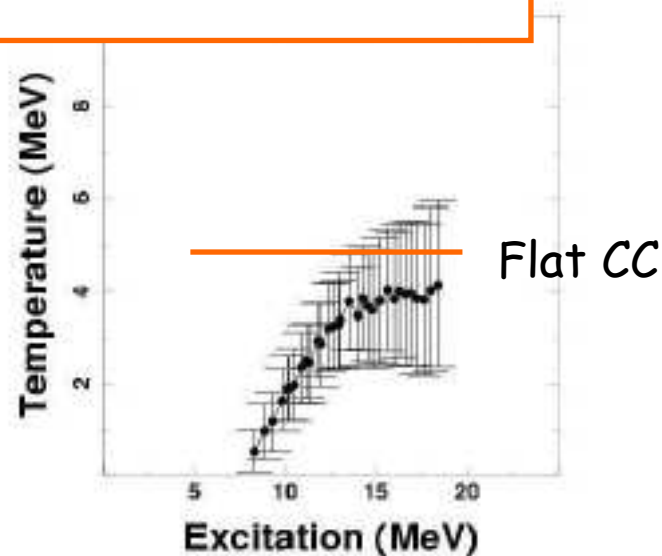
Comportamiento Critico

Curvas Calóricas

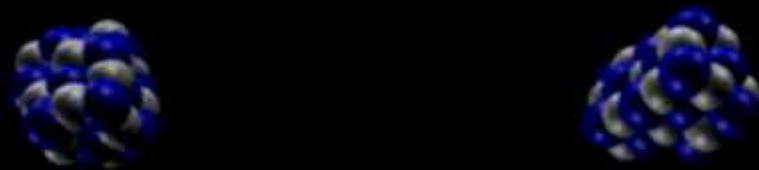
Calor especifico negativo

Isoscaling

, etc



Modelo CMD



Reconocimiento de Fragmentos

MST_algorithm

Given a set of particles i, j, k, \dots , clusters are defined such that :

$$i \in C \Leftrightarrow \exists j \in C / (r_i - r_j) < r_{clust}$$

With r_{clust} the clusterization radius

With periodic boundary conditions

MSTE algorithm

In this case clusters are defined in the following way:

$$i \in C \Leftrightarrow \exists j \in C / \left(\frac{p_{ij}^2}{4\mu} - v_{ij} \right) < 0$$

With μ the reduced mass and p_{ij} the relative momentum

Reconocimiento de Fragmentos

Early Cluster Formation Model

Given a set of particles (i,j,k..) clusters are defined as those partitions C_i that minimize the following expression:

$$E = \sum_{C_i} \left[\left(\sum_{i \in C_i} \frac{P_{ij}^2}{2m} \right)_{c.m.C_i} + \sum_{i < j \in C_i} v(r_{ij}) \right]$$

This is a highly self consistent problem that has been solved by devising a method in the spirit of simulated annealing (ECRA). For such a problem a Markov chain in the space of partitions is constructed

Modelo CMD

Table 1. Comparison of coefficients obtained for different models.

Coefficient	Stiff	Medium	Experimental
C_v	16.1	17.37	15.75
C_s	-11.73	-14.38	-17.8
C_c	-0.197	-0.226	-0.177
C_{sym}	-34.07	-25.08	-23.7

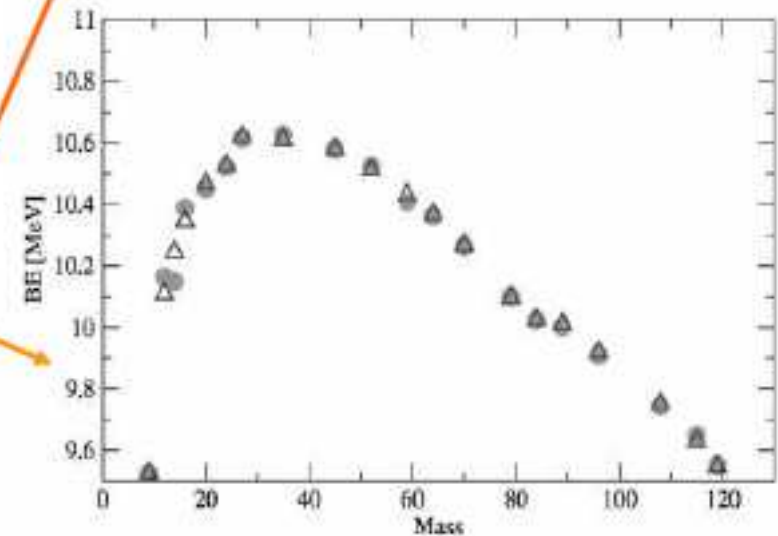
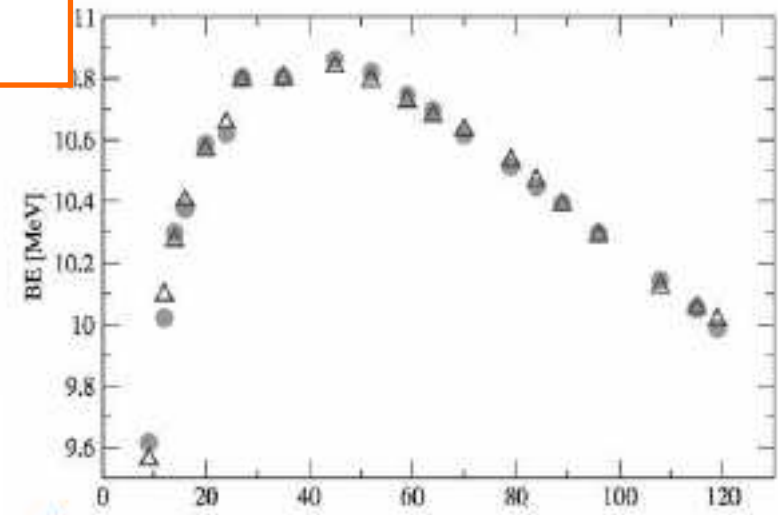
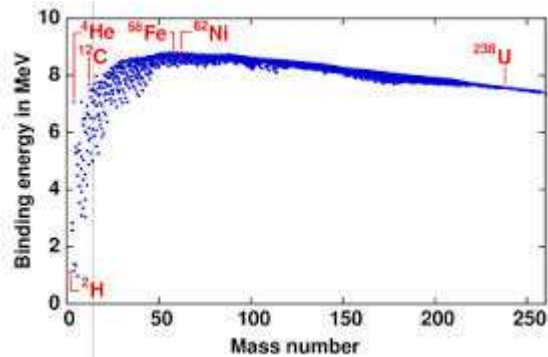


Figure 1. Energies obtained with the mass formula fit (triangles) for the stiff and m (top and bottom panels, respectively) together with the corresponding ground state using frictional molecular dynamics (circles).

$$E_B = C_v A - C_s A^{2/3} - C_{sym} \frac{(A - 2Z)^2}{A^{1/3}} - C_c \frac{Z(Z-1)}{A^{1/3}} + \delta(A, Z)$$

CMD and isoscaling

Comparison to other models

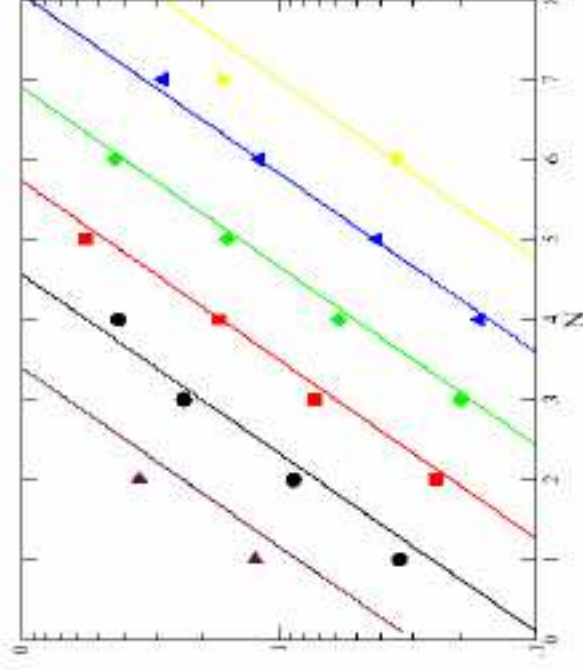
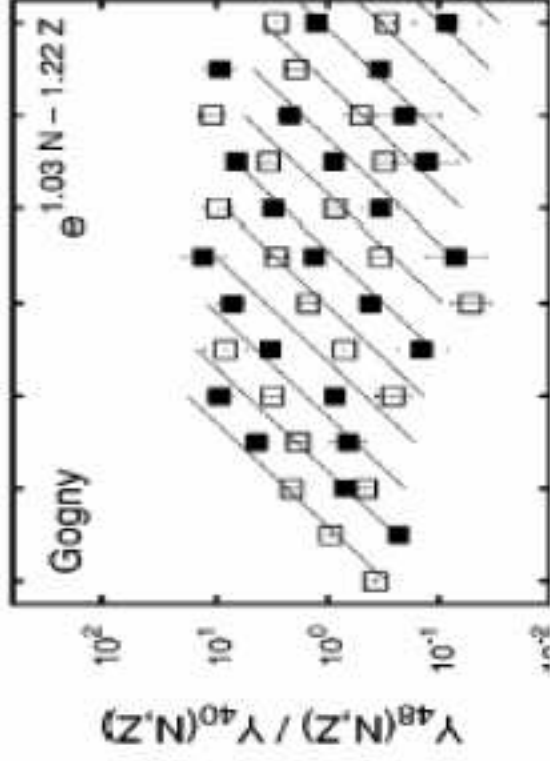


Compare $^{40}\text{Ca} - ^{48}\text{Ca}$ at 35 MeV/A
AMD-V: $\alpha=1.03$, $\beta=1.22$ CMD: $\alpha=1.07$, $\beta=1.22$

AMD



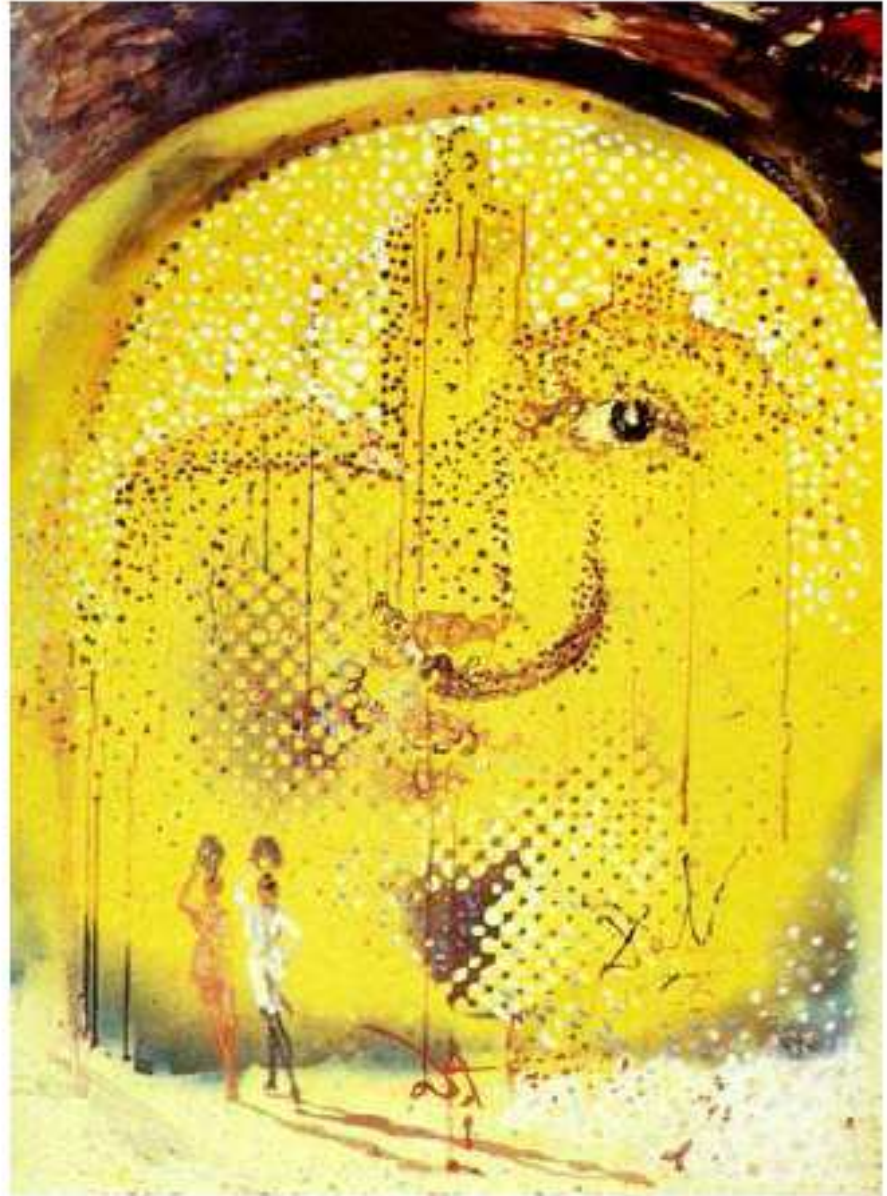
Classical
molecular dynamics



Molecular dynamics produces isoscaling

- Isoscaling is not quantal

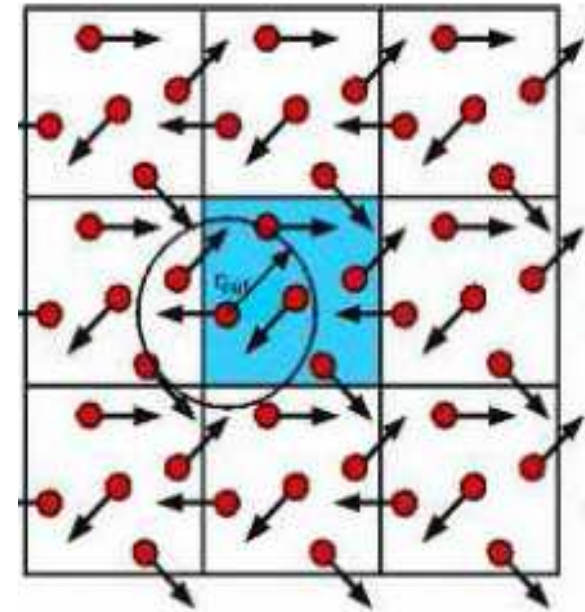
**De los núcleos
a las
Neutron Stars**



Nuestro trabajo : usar dinámica molecular con *CMD* para estudiar Estrellas de Neutrones

Resolver las ecuaciones dinámicas como antes
Pero

Usar condiciones periódicas de contorno



Para Coulomb usamos Aprox. Thomas Fermi

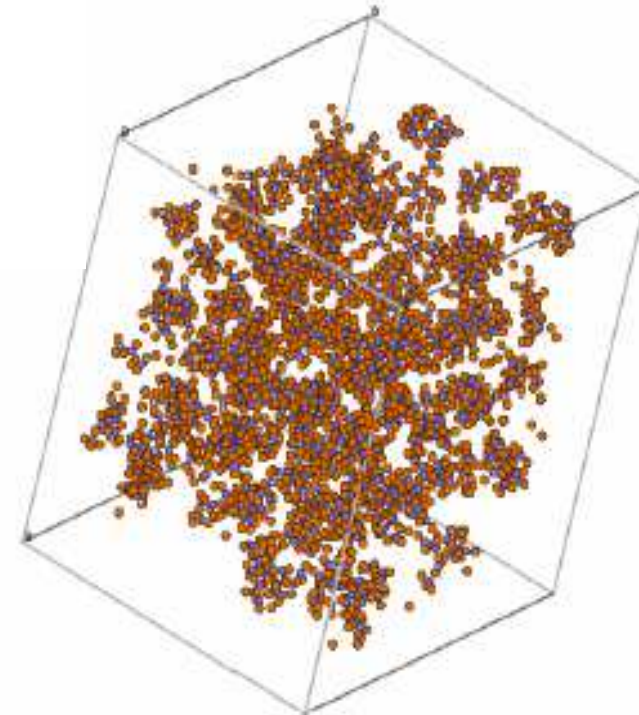
o
Debye

$$V_C(r) \approx \frac{e^2}{r}$$

$$V_C^{TF}(r) \approx \frac{e^2}{r} \exp(-r/\lambda)$$

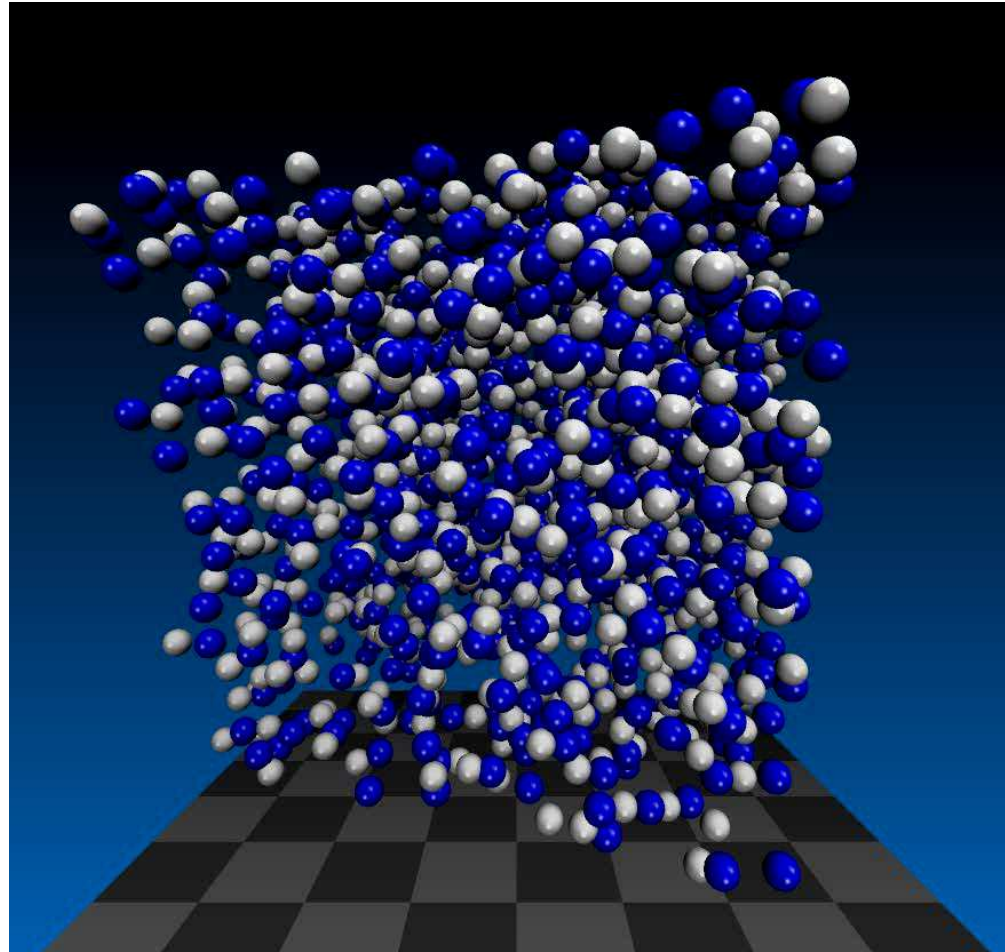
(sistema neutro)

$$\lambda_e = \pi^{1/2} / [2e(3\pi^2 n_e)^{1/3}]$$



5000 Partículas en celda elemental, T si going down.....

$$0.1\text{MeV} < T < 3\text{MeV}$$



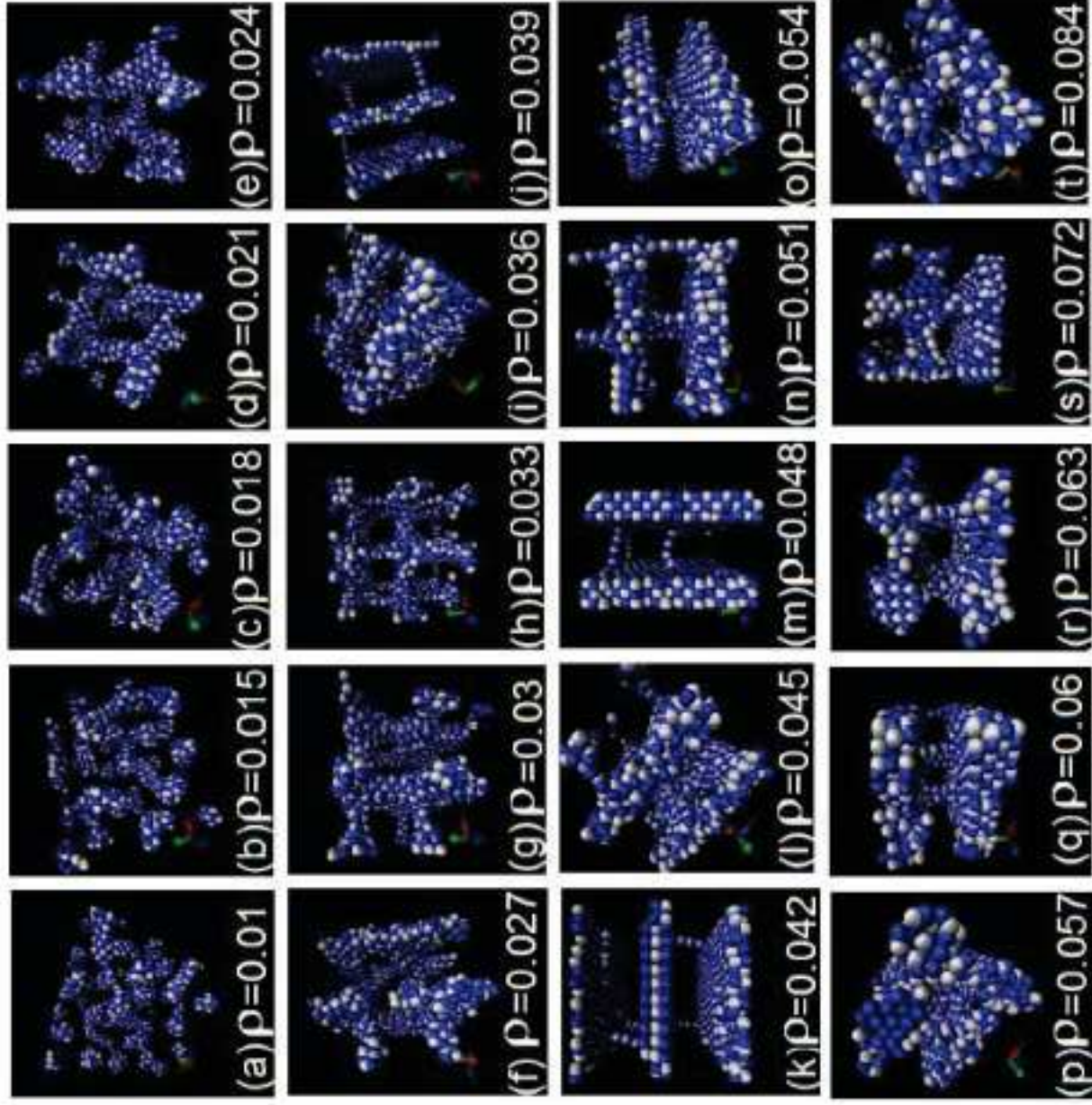
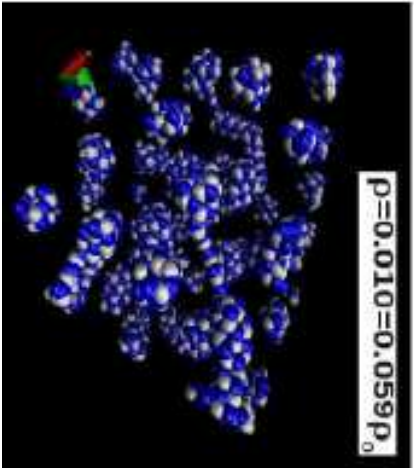
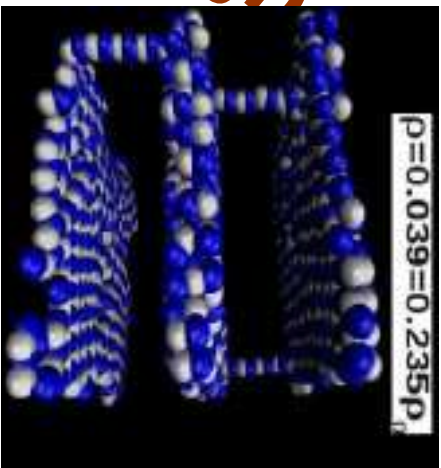


FIG. 1. (Color online) Smörgåsbord of pasta shapes corresponding to the densities shown and to $x = 0.5$ and $T = 0.1$ MeV.

Lasagn

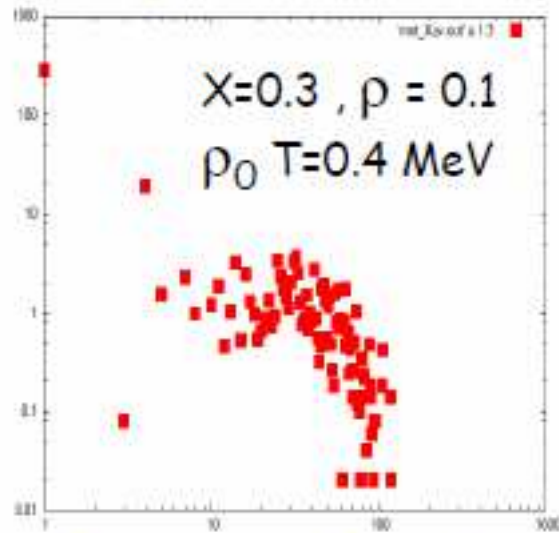
Spaghe

Gnocchi

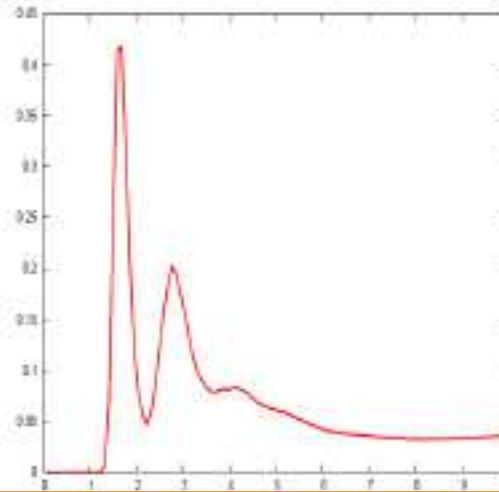


Pasta !

Study:

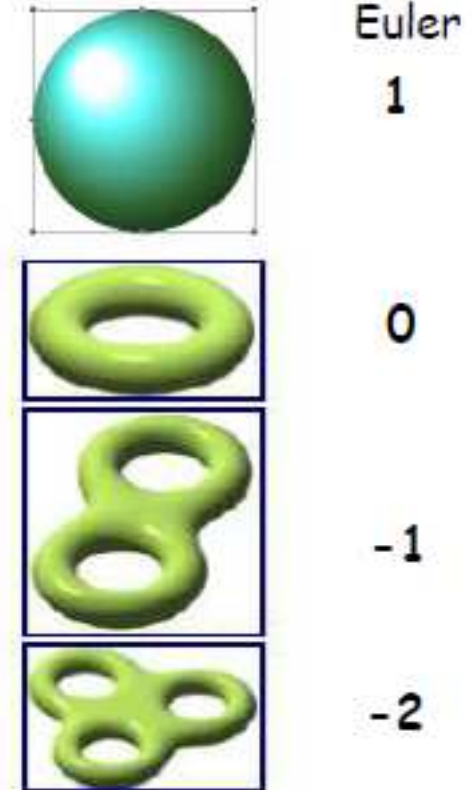


fragmentos y
distribución de masas



función de correlación
de pares

Topology (why?)



Funcionales de
Minkowski

Results: Nuclear Pasta!

...

Fragment Size Distribution (MST)

$$x=0.3$$

$$\rho=0.1\rho_0$$

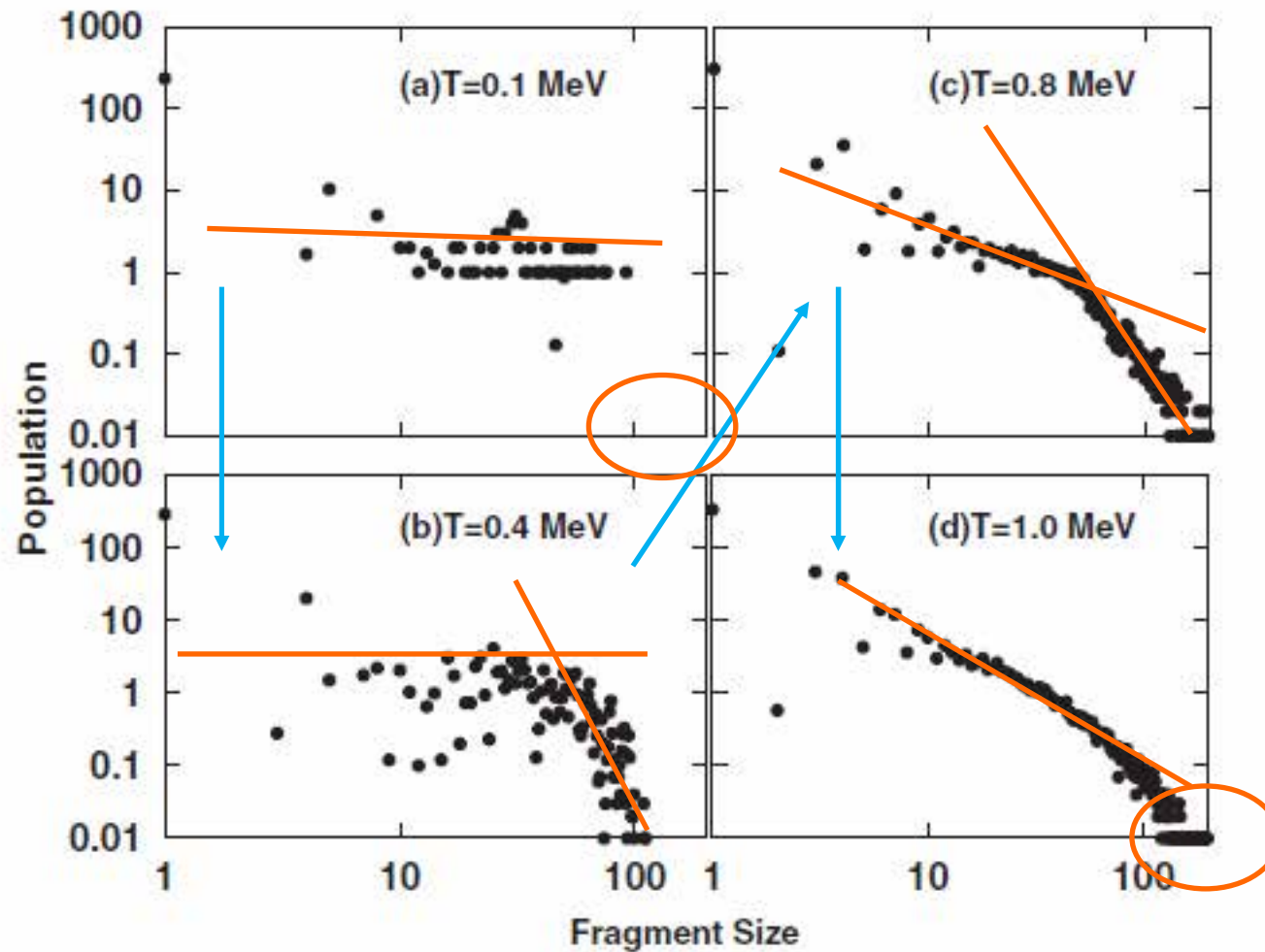


FIG. 2. Temperature evolution of the fragment size distribution obtained from 200 configurations with $x = 0.3$ and $\rho \approx \rho_0/10$.

Característica de Euler

Tomar una hoja de papel

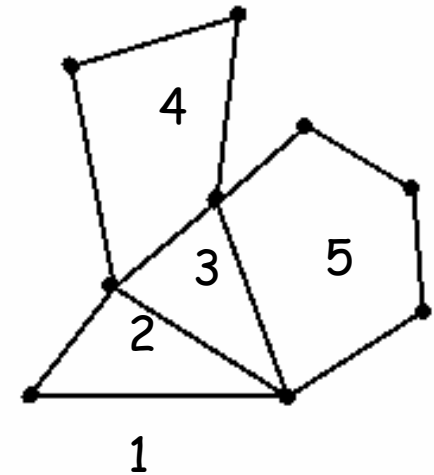
Dibujar puntos. Conectar punto con líneas de modo que : las líneas no pueden cortarse, todo punto queda conectado a todo otro punto

Cuente el numero d puntos (D), cuente las regiones determinadas por las Líneas (R) cuente las líneas (L);

$$D=9, L=12, R=5$$

$$\chi = D - L + R = 2 \dots$$

siempre






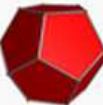
In 3D Convex Polyhedra





The Euler characteristic $\chi = V - E + F = 2$

V numbers of vertices (corners)

E numbers of edges

F numbers of faces

Name	Image	Vertices V	Edges E	Faces F	Euler characteristic: $V - E + F$
Tetrahedron		4	6	4	2
Hexahedron or cube		8	12	6	2
Octahedron		6	12	8	2
Dodecahedron		20	30	12	2
Icosahedron		12	30	20	2

Name	Image	Vertices V	Edges E	Faces F	Euler characteristic: $V - E + F$
Tetrahemihexahedron		6	12	7	1
Octahemioctahedron		12	24	12	0
Cubohemioctahedron		12	24	10	-2
Great icosahedron		12	30	20	2

In 3D non-convex polyhedra χ can have any value, and thus can be used to characterize the shape

Lo aplicamos a "Pastas"

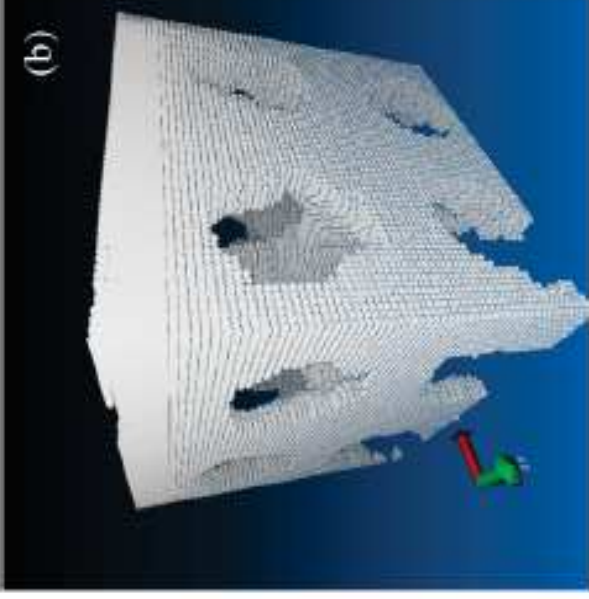
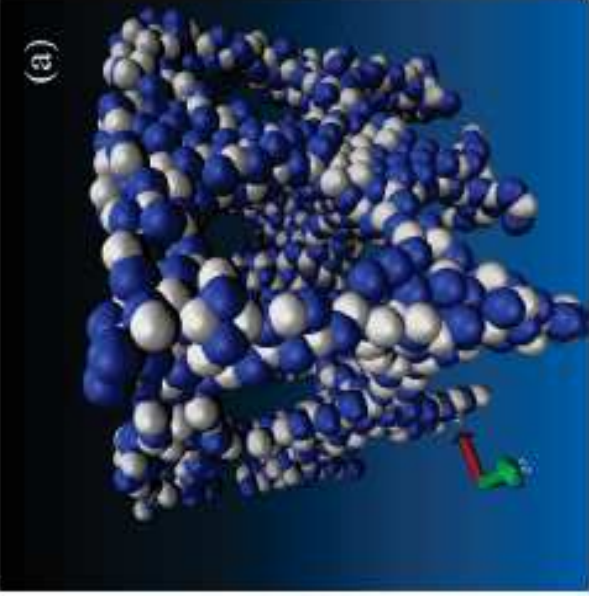


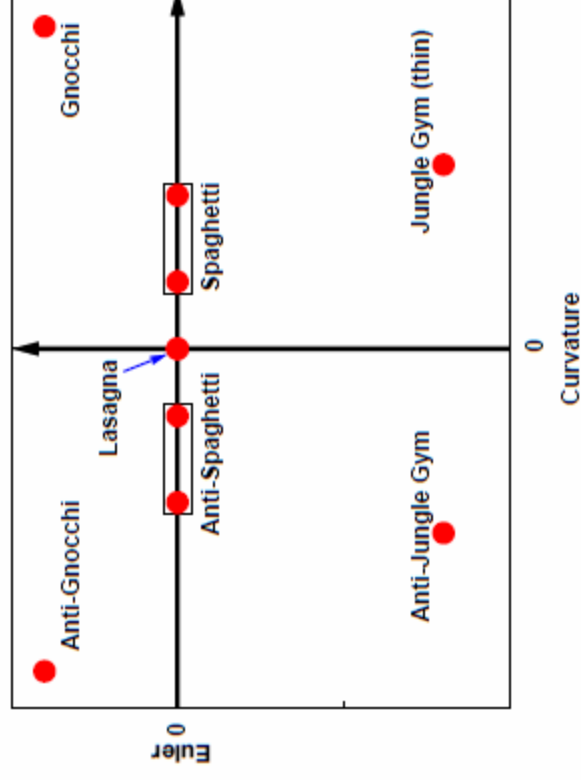
FIG. 8. (Color online) Sample transformation of a nuclear structure to a corresponding polyhedron. The structure corresponds to a case with $x = 0.5$, $\rho = 0.33 \text{ fm}^{-3}$, and $T = 0.1 \text{ MeV}$.

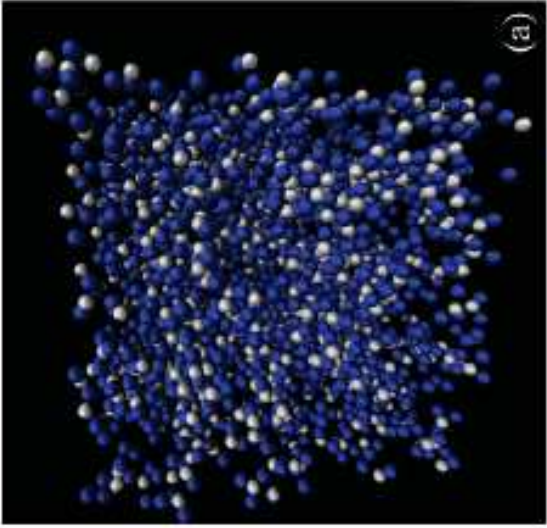
TABLE I: Classification Curvature - Euler

	Curvature < 0	Curvature ~ 0	Curvature > 0
Euler > 0	Anti-Gnocchi	Lasagna	Gnocchi
Euler ~ 0	Anti-Spaghetti		Spaghetti
Euler < 0	Anti-Jungle Gym		Jungle Gym

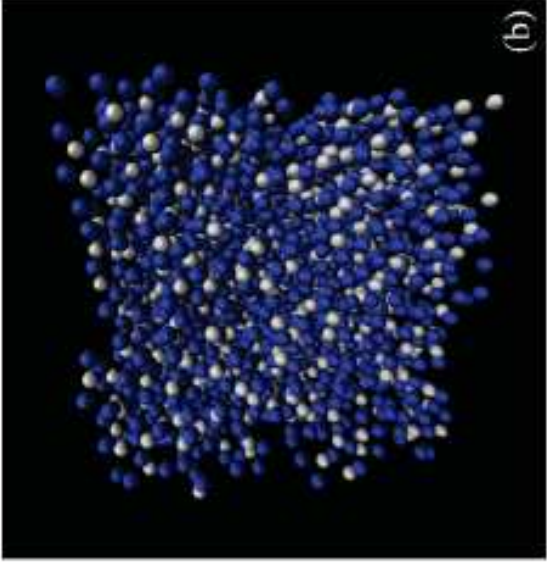
$$V = n_c, \quad A = -6n_c + 2n_f,$$

$$2B = 3n_c - 2n_f + n_c, \quad \chi = -n_c + n_f - n_e + n_v,$$



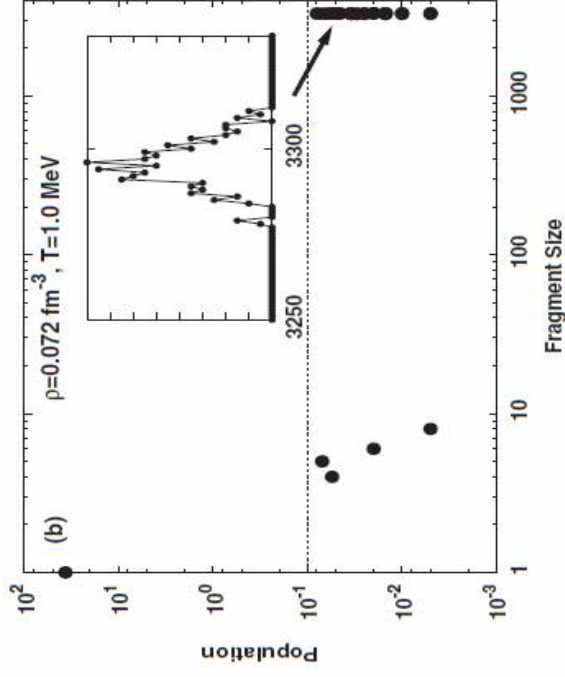
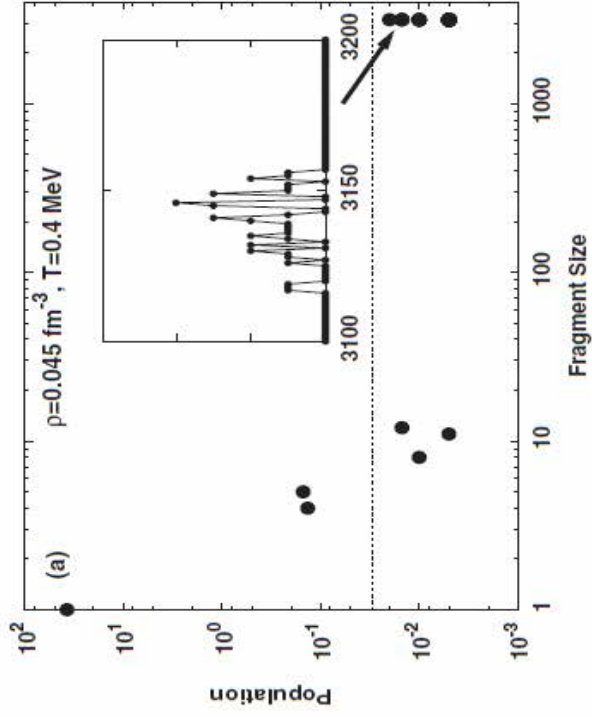


(a)

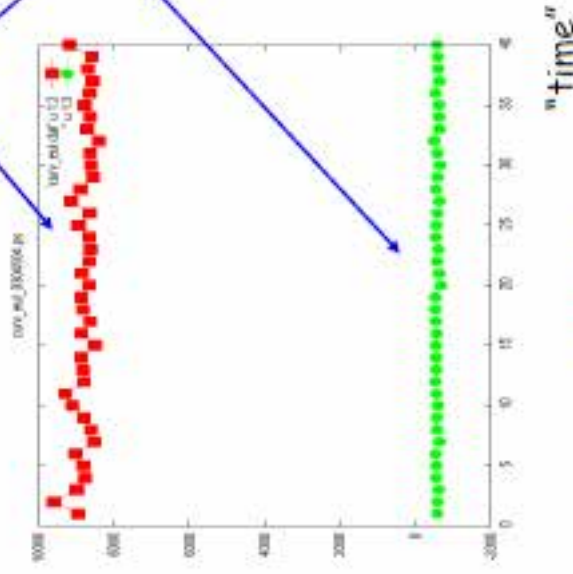


(b)

FIG. 9. (Color online) Spatial configurations formed under $T = 0.4$ MeV and $\rho = 0.045 \text{ fm}^{-3}$ (left) and $T = 1.0$ MeV and $\rho = 0.072 \text{ fm}^{-3}$ (right), both for $x = 0.3$.



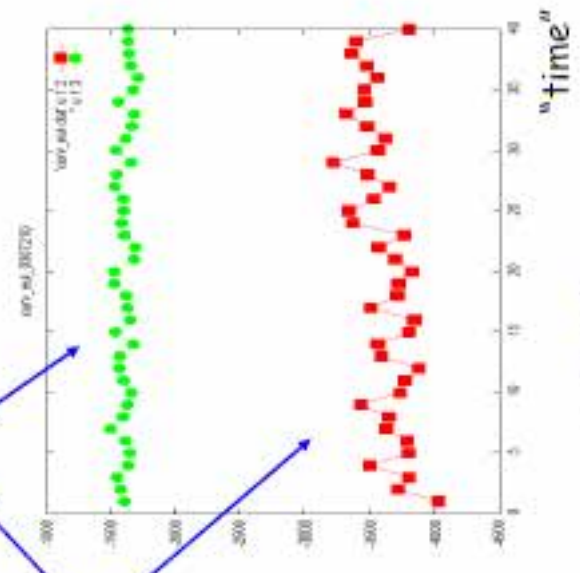
Mean curvature



$\rho = 0.045$

$X=0.3 \quad T=0.4$

Euler number (/2)



$\rho = 0.072$

$X=0.3 \quad T=1.0$

The role of coulomb

We analyze the behavior of a system driven by:

$$V(r_{ij}, K_i, K_j) = V_{Pandha}(r_{ij}, K_i, K_j) + \alpha \cdot V_{Coul}(r_{ij}, K_i, K_j)$$

with

$$0 \leq \alpha \leq 1$$


The role of coulomb

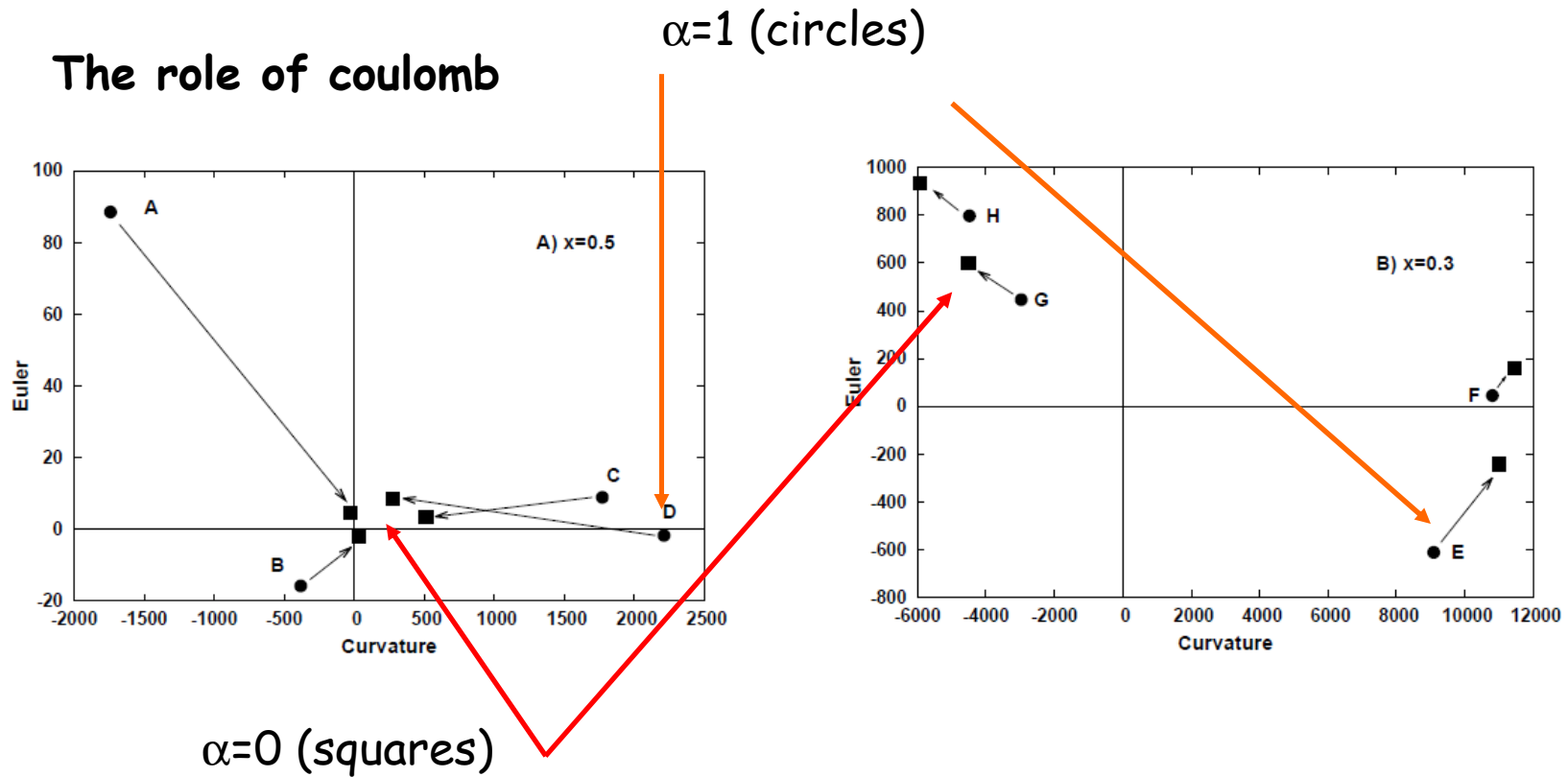


FIG. 7: Average values of the Curvature and Euler numbers of the structures listed in Table II; circles correspond to structures with Coulomb and squares to structures without Coulomb, arrows indicate the average displacement of the structures as α goes from 1 to 0.

TABLE II: Classification Curvature - Euler

	ρ (fm^{-3})	T (MeV)	$\alpha = 1$		$\alpha = 0$	
			Curvature	Euler	Curvature	Euler
A	0.072	1.0	-1740	89	-24	4.7
B	0.072	0.1	-382	-15.8	29.4	-1.9
C	0.015	0.1	1771	8.9	514	3.5
D	0.015	1.0	2212	-1.7	273	8.6
E	0.015	1.0	9090	-612	11011	-243
F	0.015	0.1	10804	454	11436	159
G	0.072	1.0	-2969	447	-4512	602
H	0.072	0.1	-4485	798	-5914	933


This suggests that there is pasta even without Coulomb!

Intermezzo : what is the role of Coulomb?

Illinois potential is rather complicated because :

$$V_{np}(r) = V_r [\exp(-\mu_r r) / r - \exp(-\mu_r r_c) / r_c] - V_a [\exp(-\mu_a r) / r - \exp(-\mu_a r_c) / r_c]$$

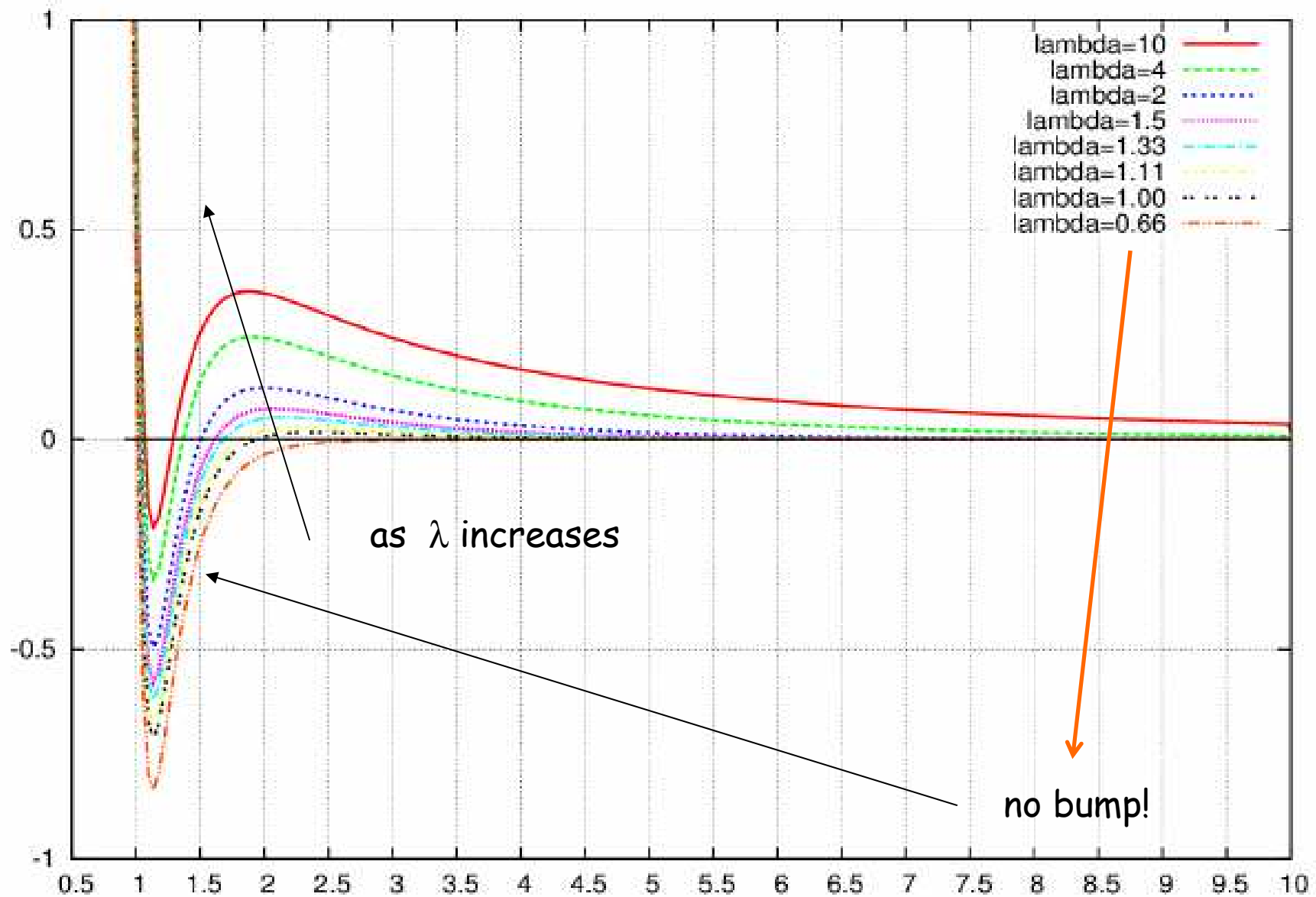
$$V_{mn}(r) = V_0 [\exp(-\mu_0 r) / r - \exp(-\mu_0 r_c) / r_c]$$

$$V_{pp}(r) = V_0 [\exp(-\mu_0 r) / r - \exp(-\mu_0 r_c) / r_c] + \exp\left(-\frac{r}{\lambda}\right) \cdot \left(\frac{v_c}{r}\right)$$


Lennard Jones + Coulomb

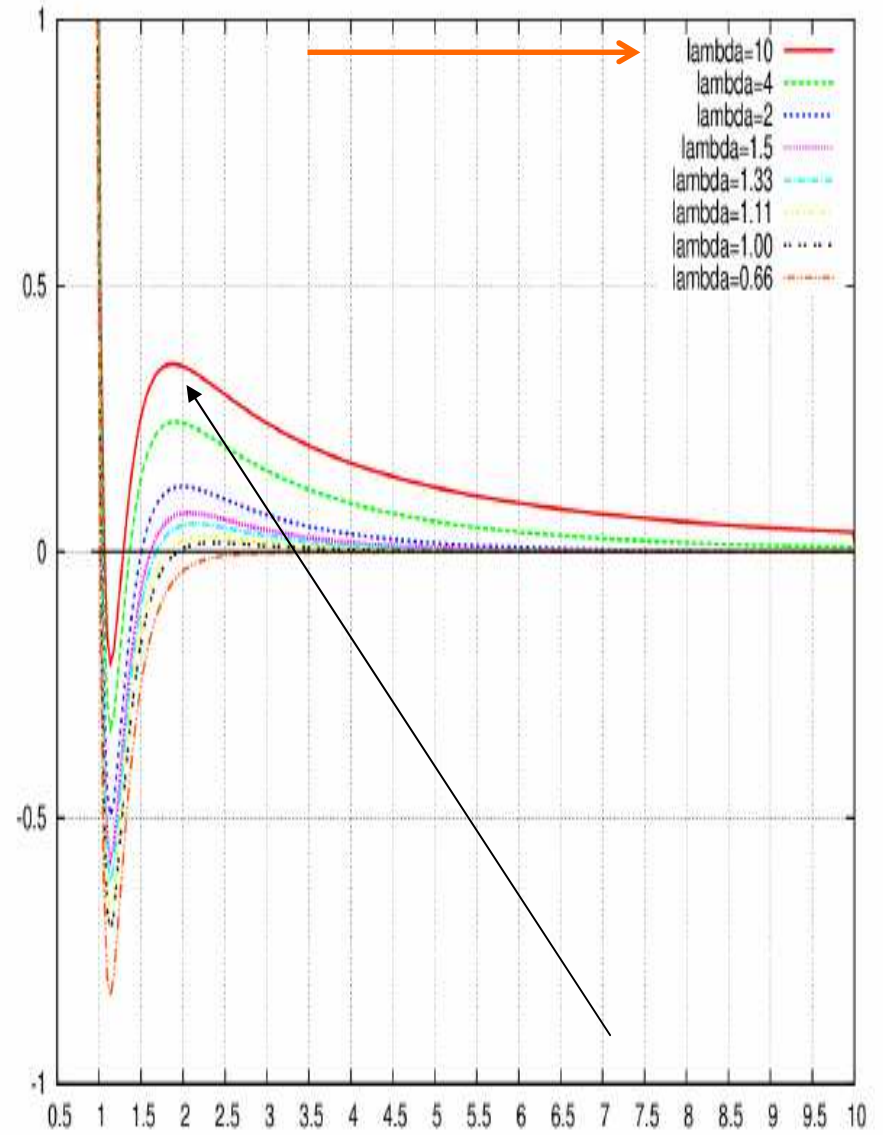
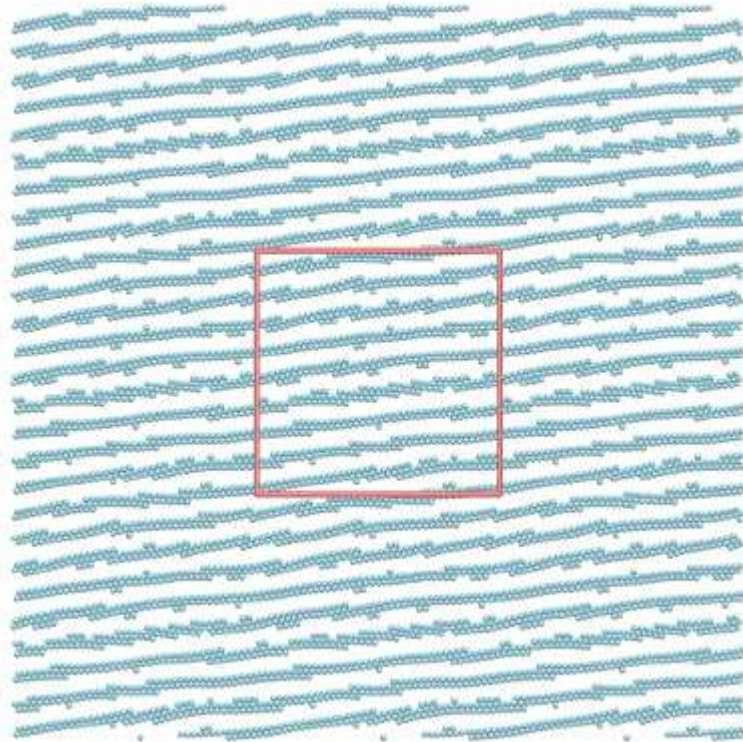
Lennard Jones +coulomb

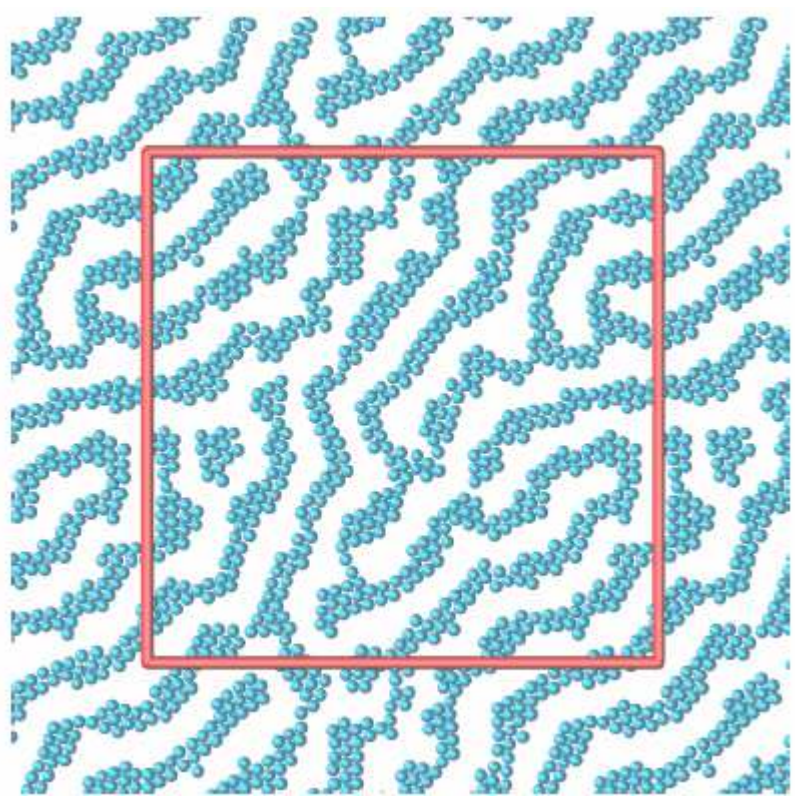
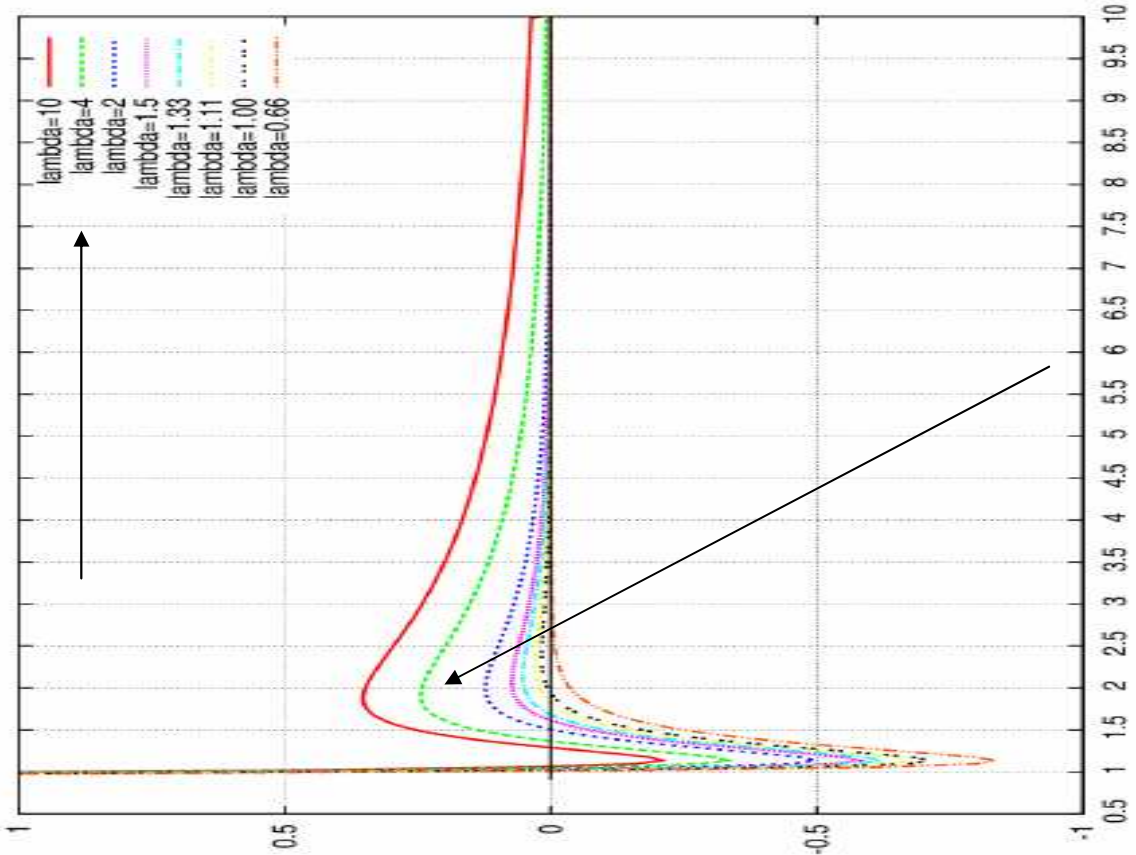
$$V(r) = V_0 \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right] + \exp\left(-\frac{r}{\lambda}\right) \cdot \left(\frac{v_c}{r} \right)$$

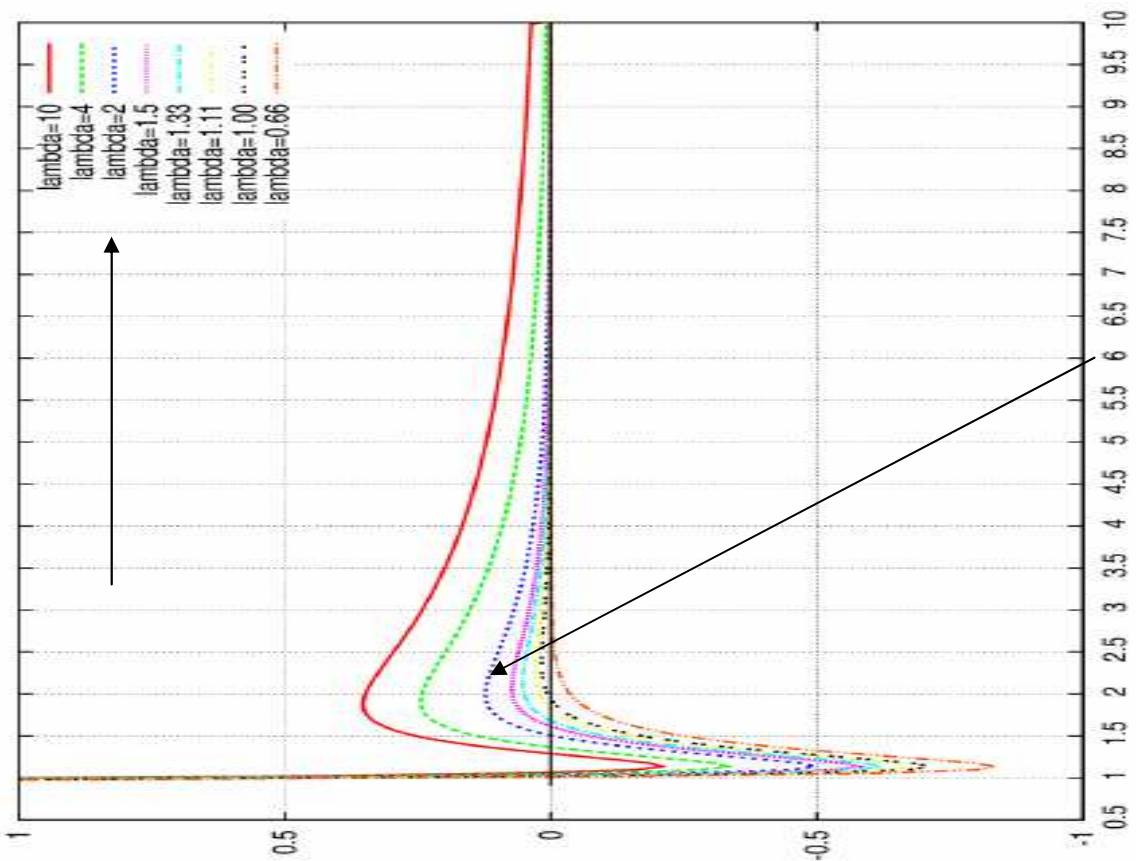
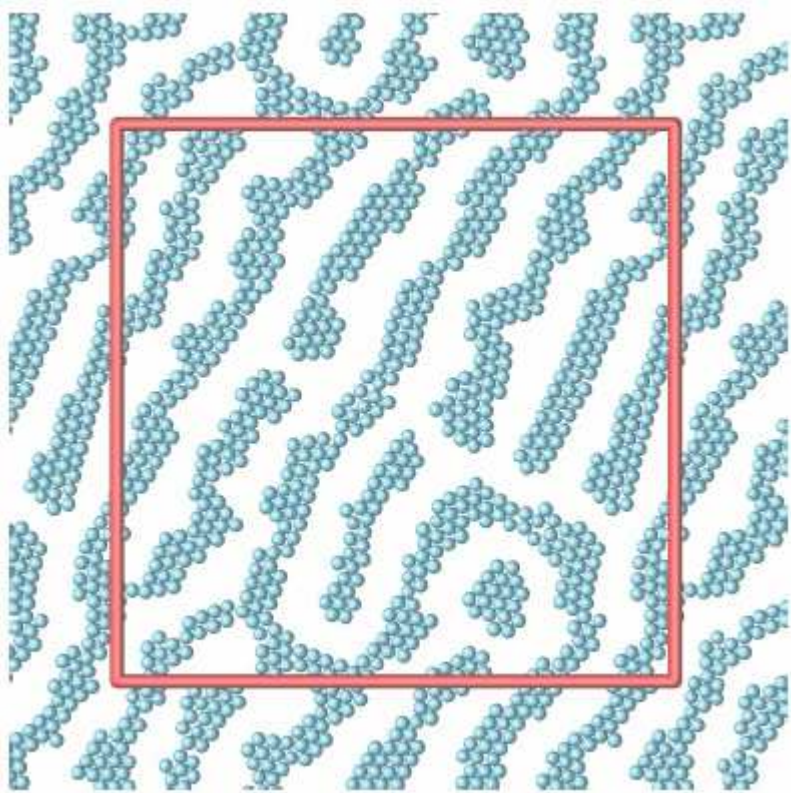


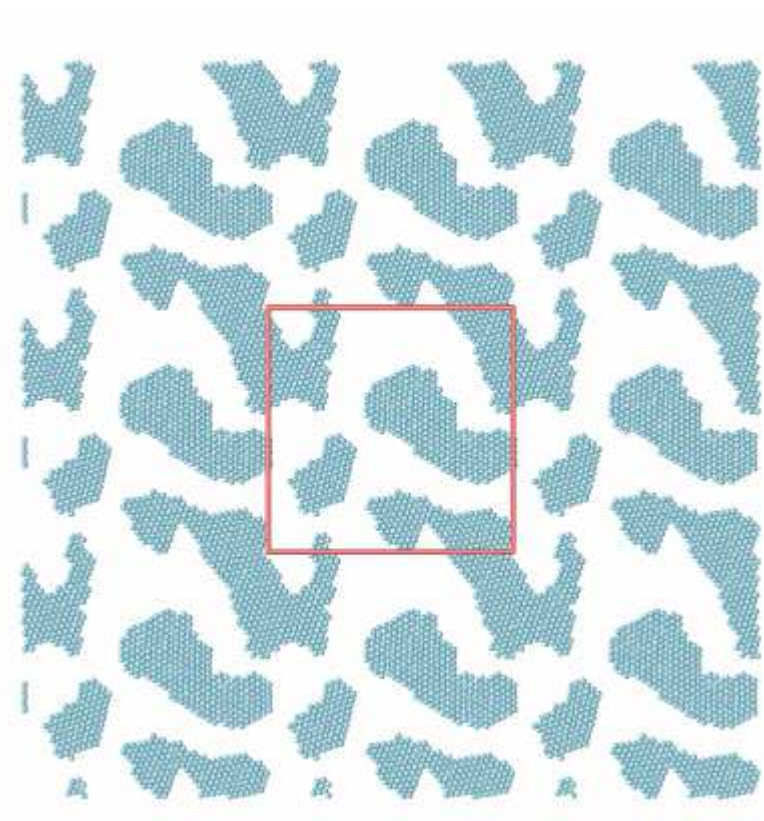
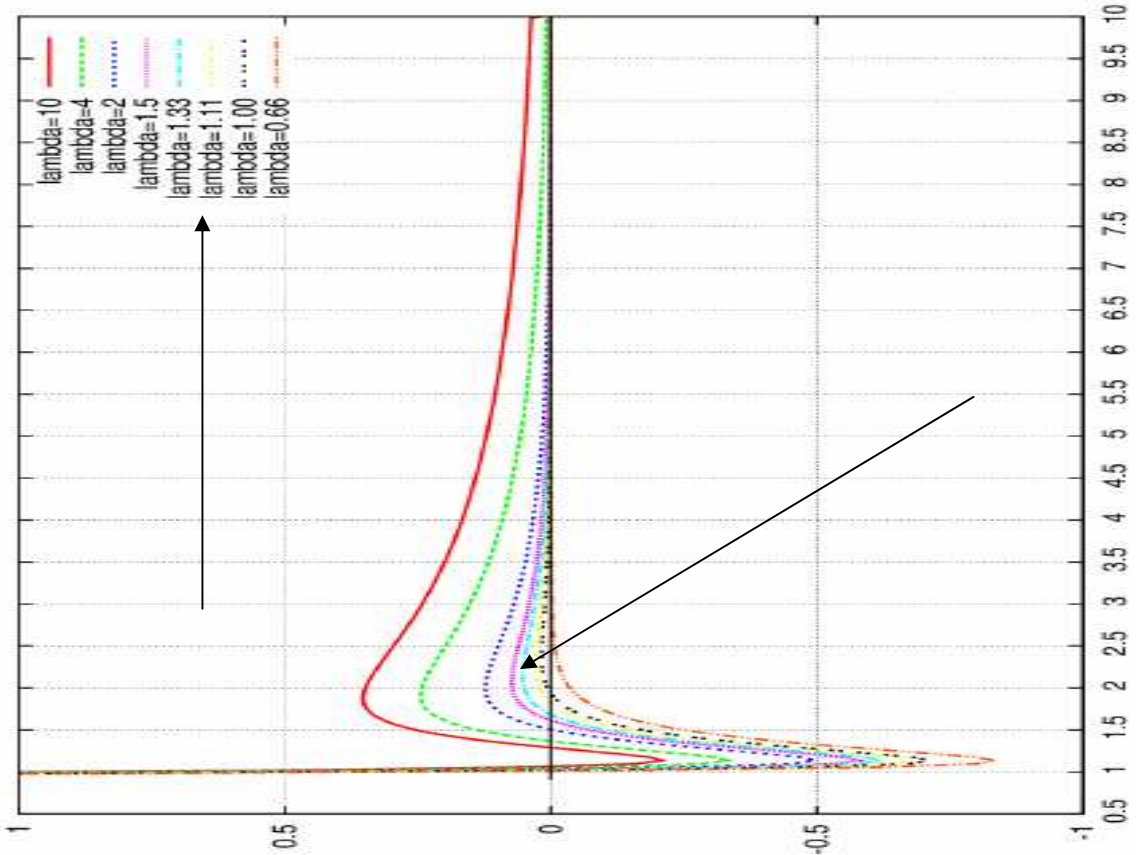
The role of coulomb

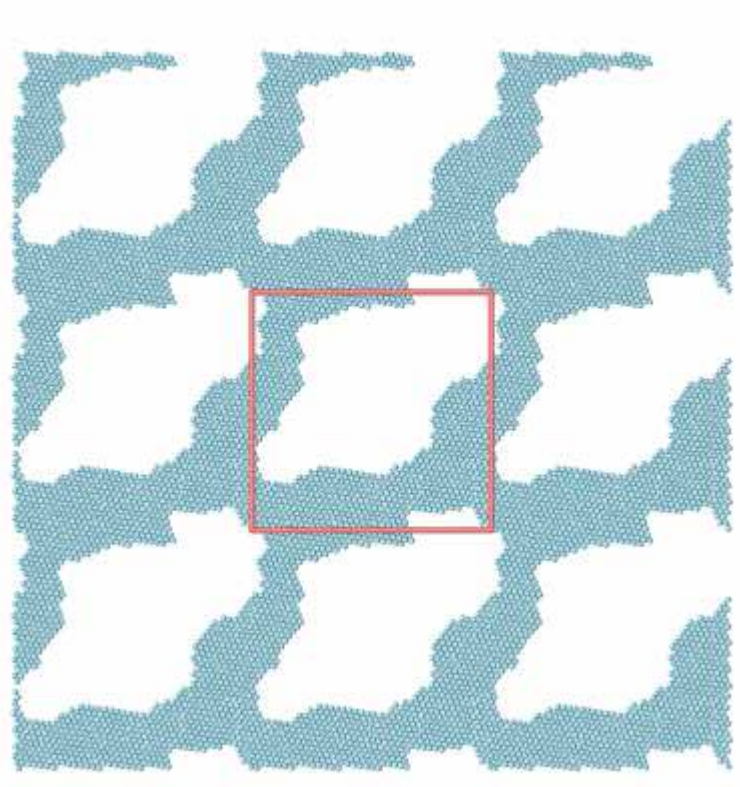
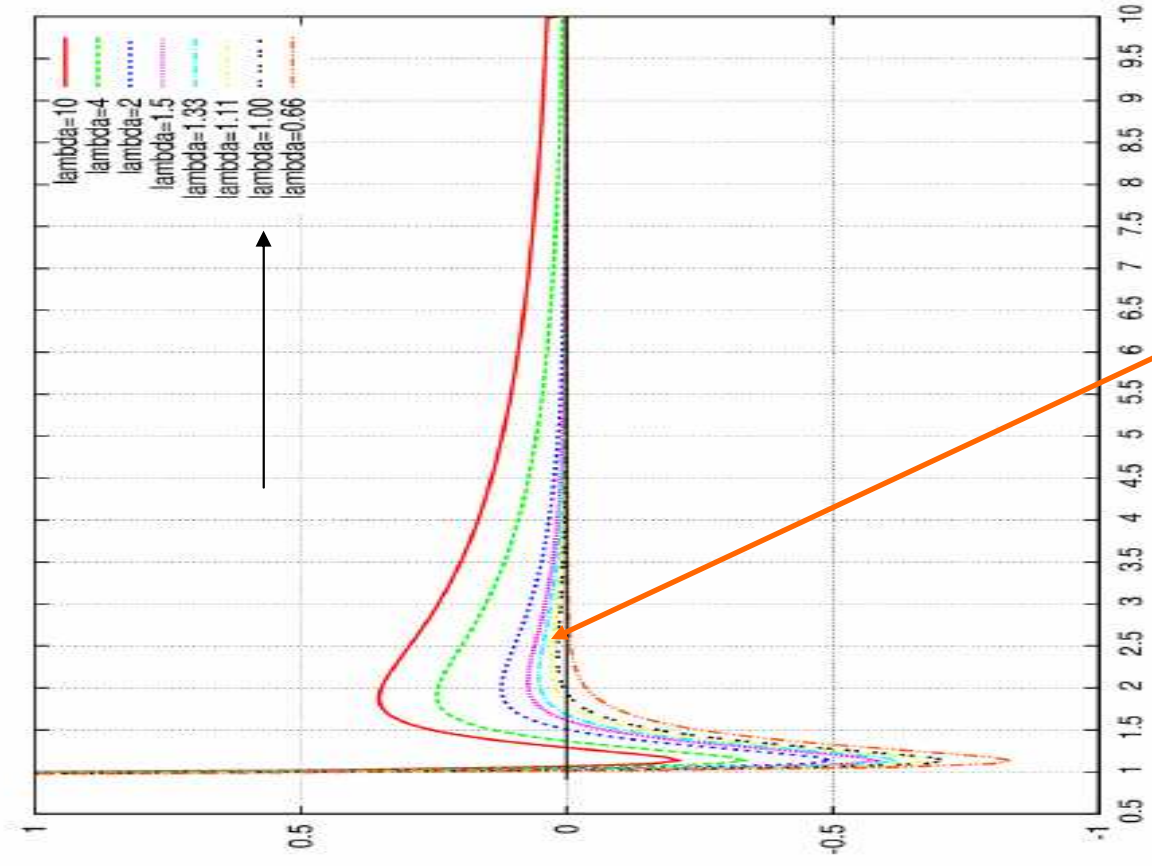
N=800

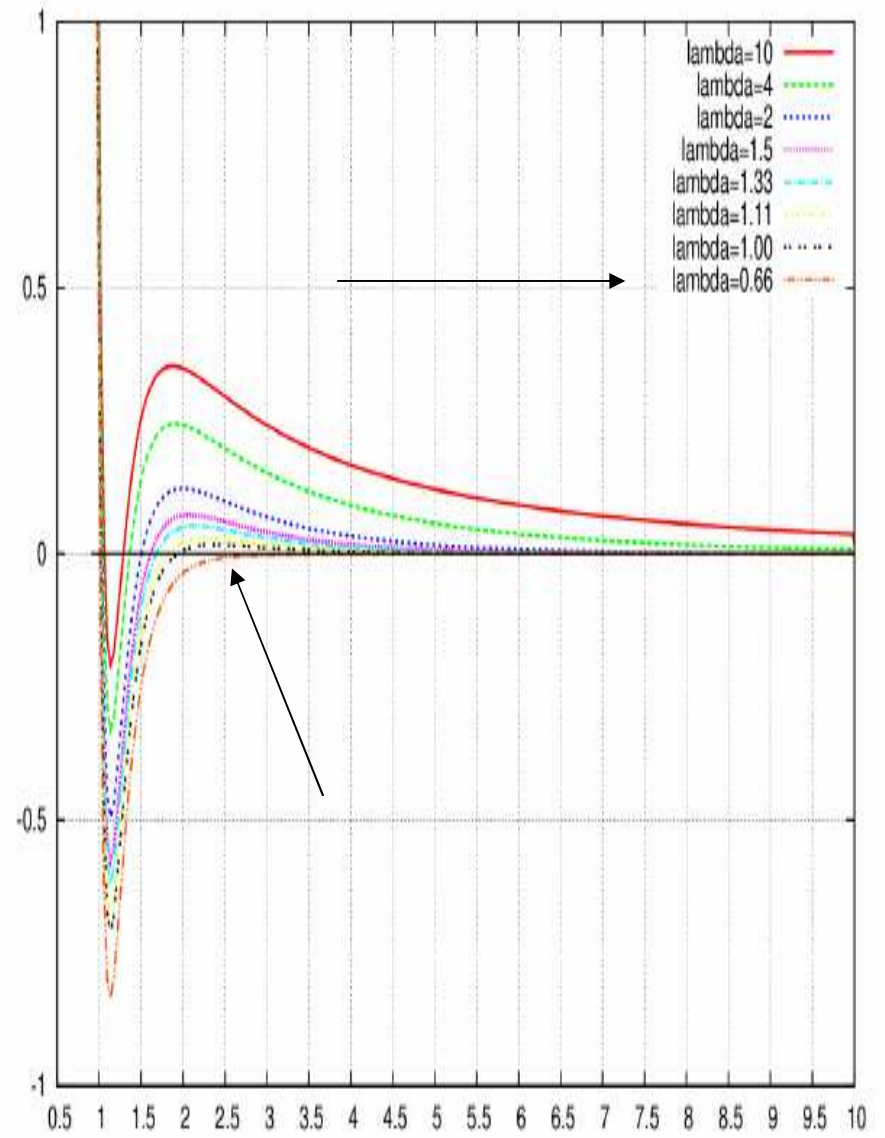
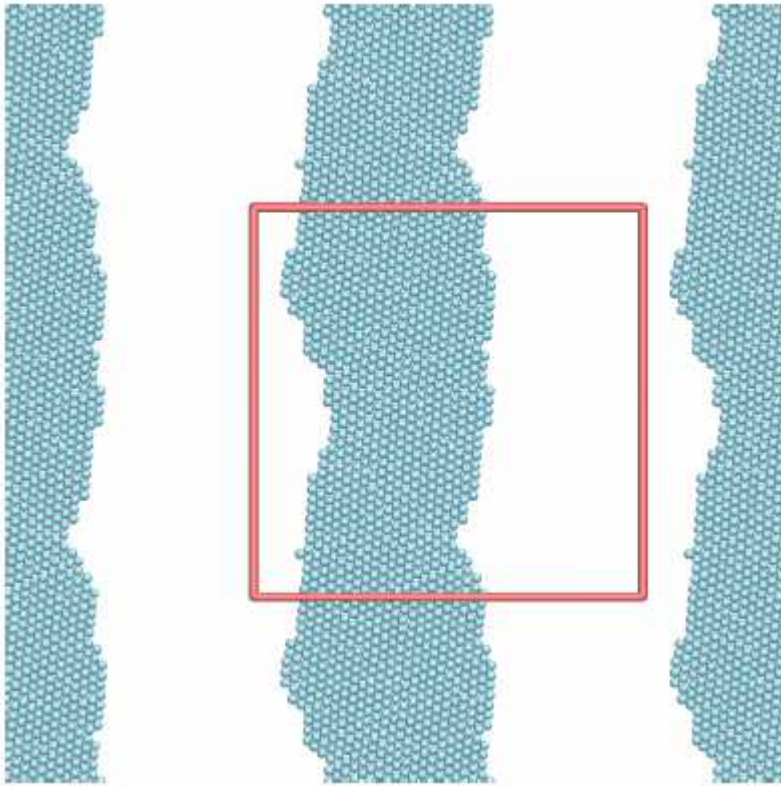






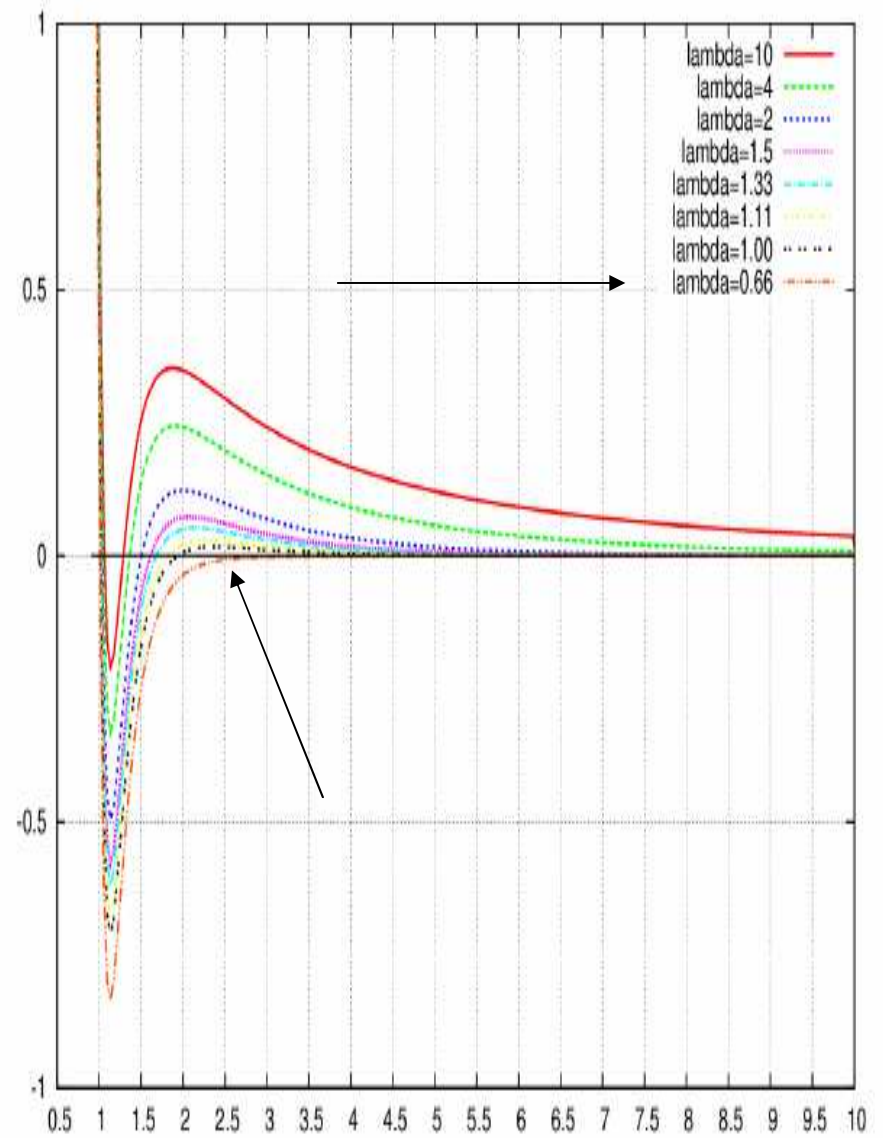
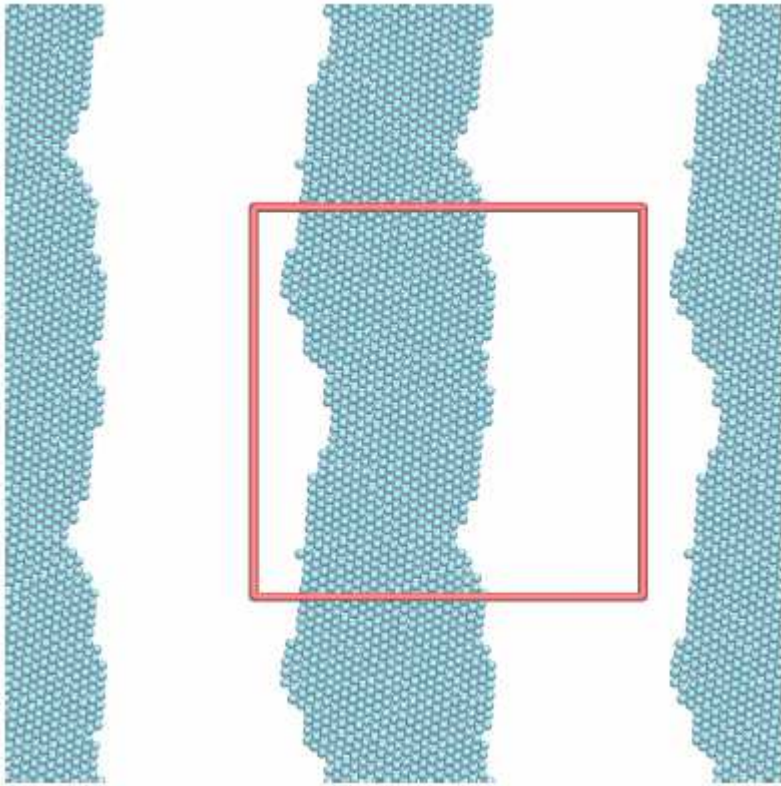






no bump, still get 'pasta'

Trivial infinite cluster
2 dimensions "no coulomb"



no bump, still get 'pasta'

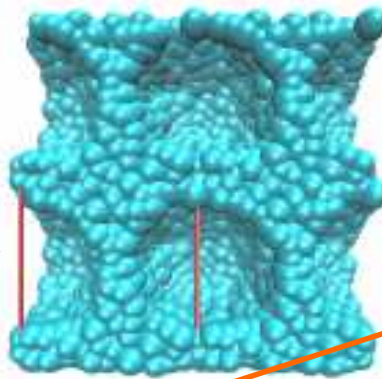
In 3 dimensions, no coulomb

$$V_{LJ}(r) = 4\epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right]$$

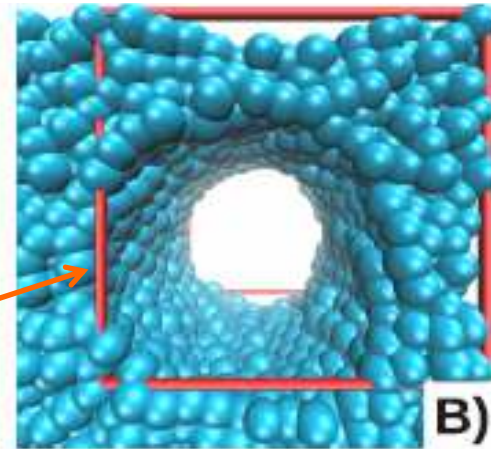
N=5000

Temperatura
muy baja

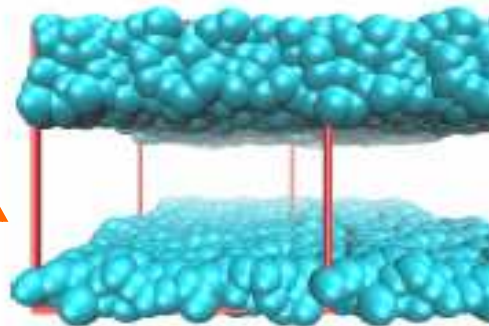
1 structure per
cell



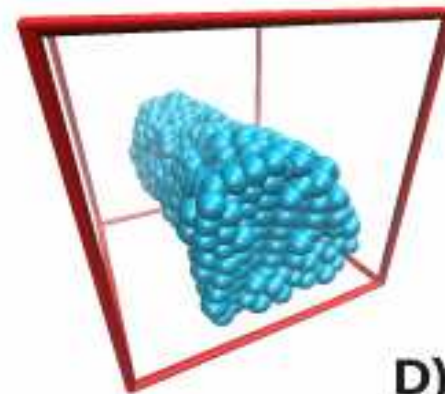
A)



B)



C)



D)

Finite size, periodic boundary conditions
and the appearance of "Pasta" without Coulomb



Finite size, periodic boundary conditions and the appearance of “Pasta” without Coulomb

The system is finite but not too small, particles interact by a short range potential

Given a configuration
we can write:

$$E = E_{bulk} + E_{surf}$$

$$A = E - TS$$

$$S_{sphere} = 4\pi \left(\frac{3}{4\pi} \right)^{\frac{2}{3}} \times u^{\frac{2}{3}} \times L^2$$

$$S_{rod} = 2 (\pi)^{\frac{1}{3}} \times u^{\frac{1}{3}} \times L^2$$

$$S_{slab} = 2 \times L^2$$

Surfaces

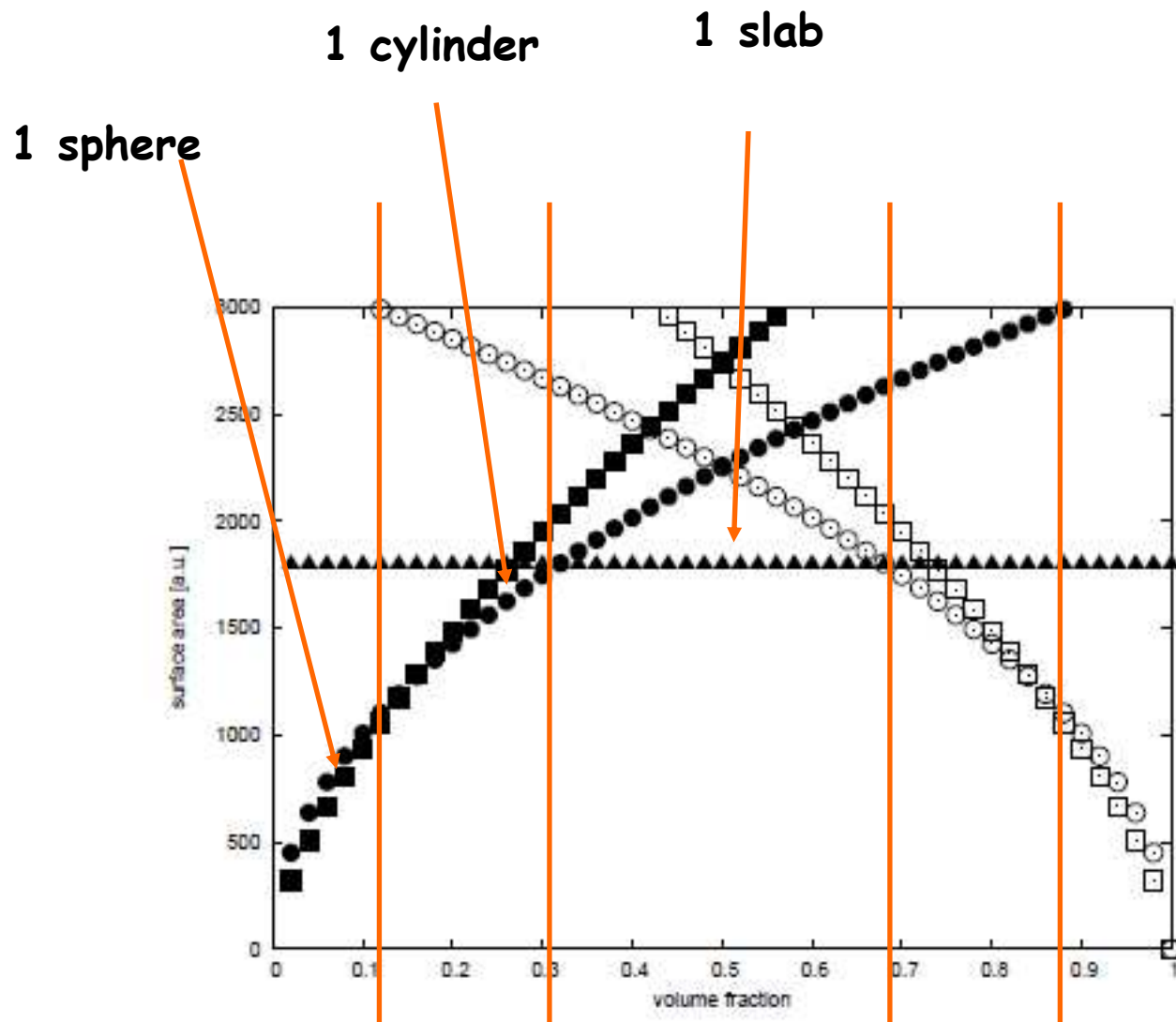


FIG. 12. Surface area of various simple shapes as a function of volume fraction for a cell of $L = 30$. The shapes are: Sphere (full squares), Cylindrical rod (full circles), slab (full triangles), cylindrical bubble (empty circles) and spherical bubble (empty squares)

En ausencia de Coulomb y debajo de una cierta T_c aparecen pseudo pastas , Una por celda, La Celda fija la escala
Las CPC siempre presentes!!!!

Nuclear Matter CMD



Back to CMD

Illinois Potential Medium

$$V_{np}(r) = V_r [\exp(-\mu_r r)/r - \exp(-\mu_r r_c)/r_c] - V_a [\exp(-\mu_a r)/r - \exp(-\mu_a r_c)/r_c]$$
$$V_{NN}(r) = V_0 [\exp(-\mu_0 r)/r - \exp(-\mu_0 r_c)/r_c],$$

Nuclear Matter

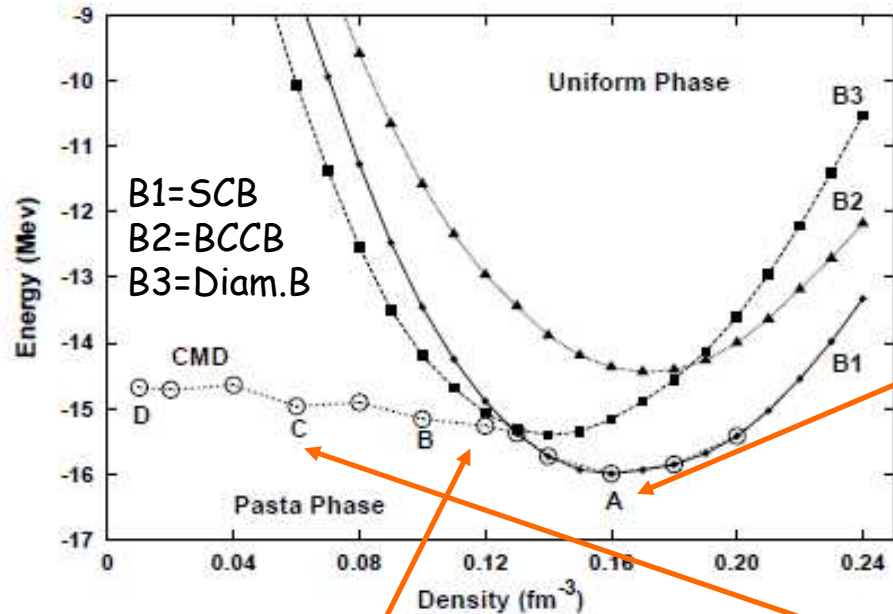


FIG. 1. Binding energies per nucleon for systems obtained with the Pandharipande medium potential for crystalline lattices with $B1$, $B2$ and $B3$ crystal geometries, and using molecular dynamics at $T = 0.001 \text{ MeV}$ (CMD). The structures corresponding to the four labeled points (“A” through “D”) are shown in Figure 2.

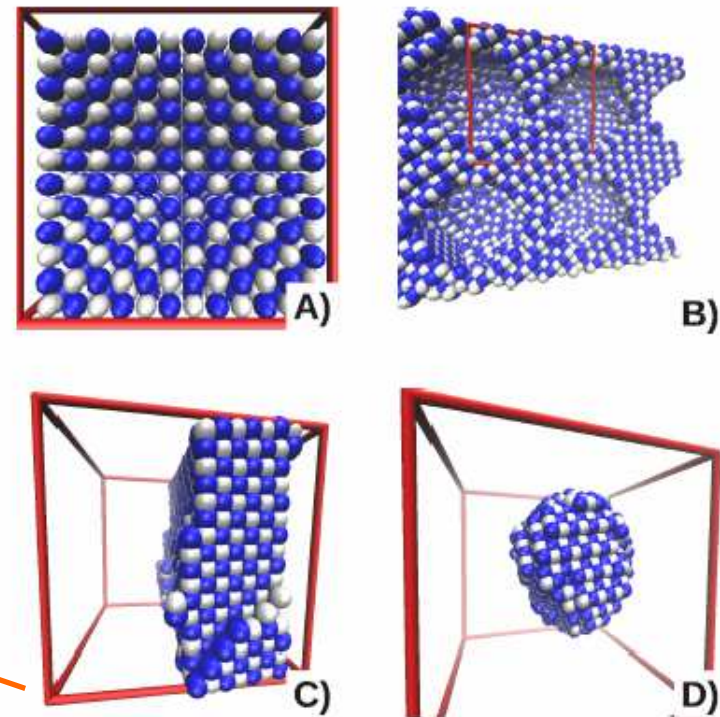


FIG. 2. (Color online) Structures corresponding to the labeled points of Figure 1. Point A corresponds to a formation in the regular ($B1$) lattice, while the rest of the points are non-homogeneous structures.

This work

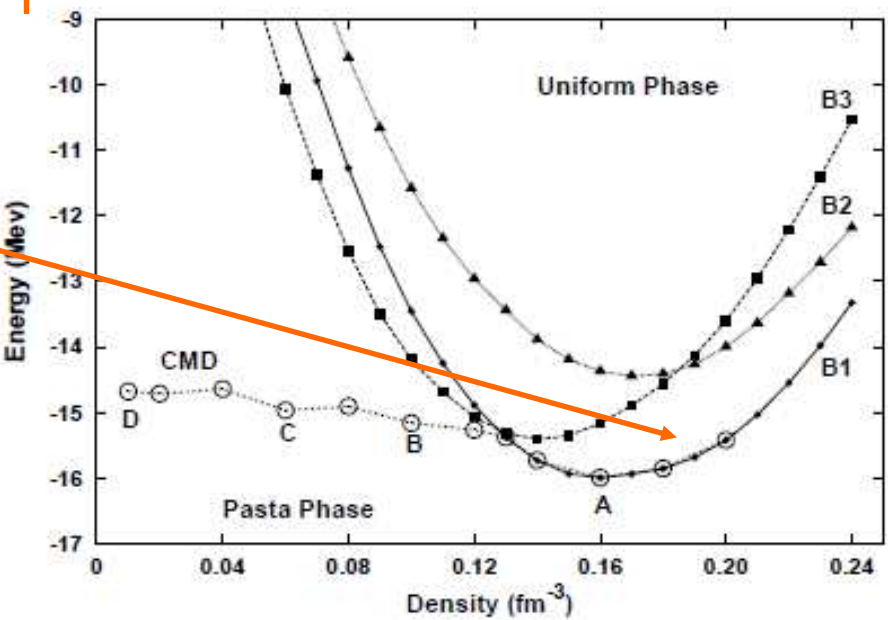
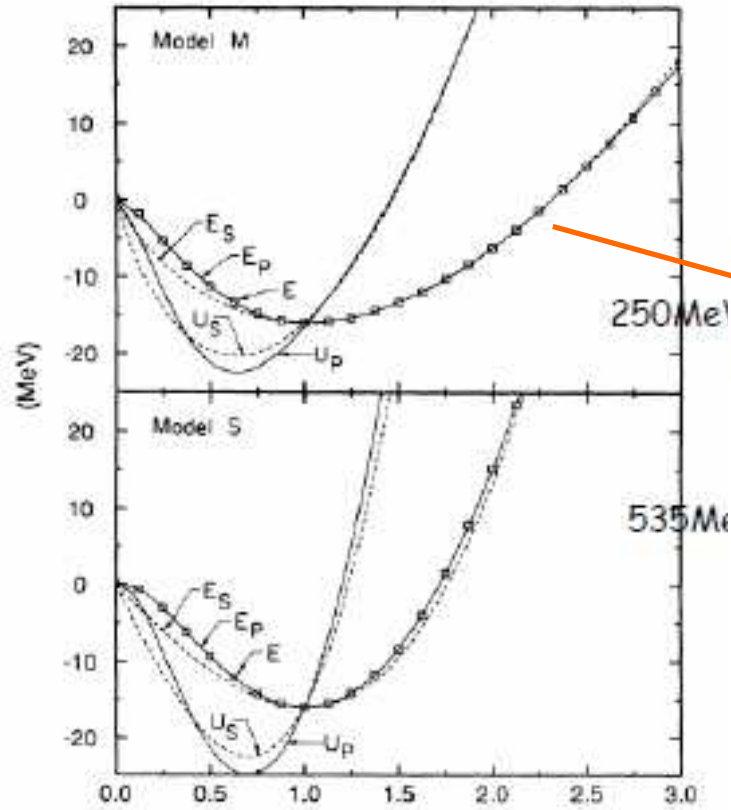


FIG. 1. Binding energies per nucleon for systems obtained with the Pandharipande medium potential for crystalline lattices with $B1$, $B2$ and $B3$ crystal geometries, and using molecular dynamics at $T = 0.001 \text{ MeV}$ (CMD). The structures corresponding to the four labeled points ("A" through "D") are shown in Figure 2.

original

Illinois Potential Medium

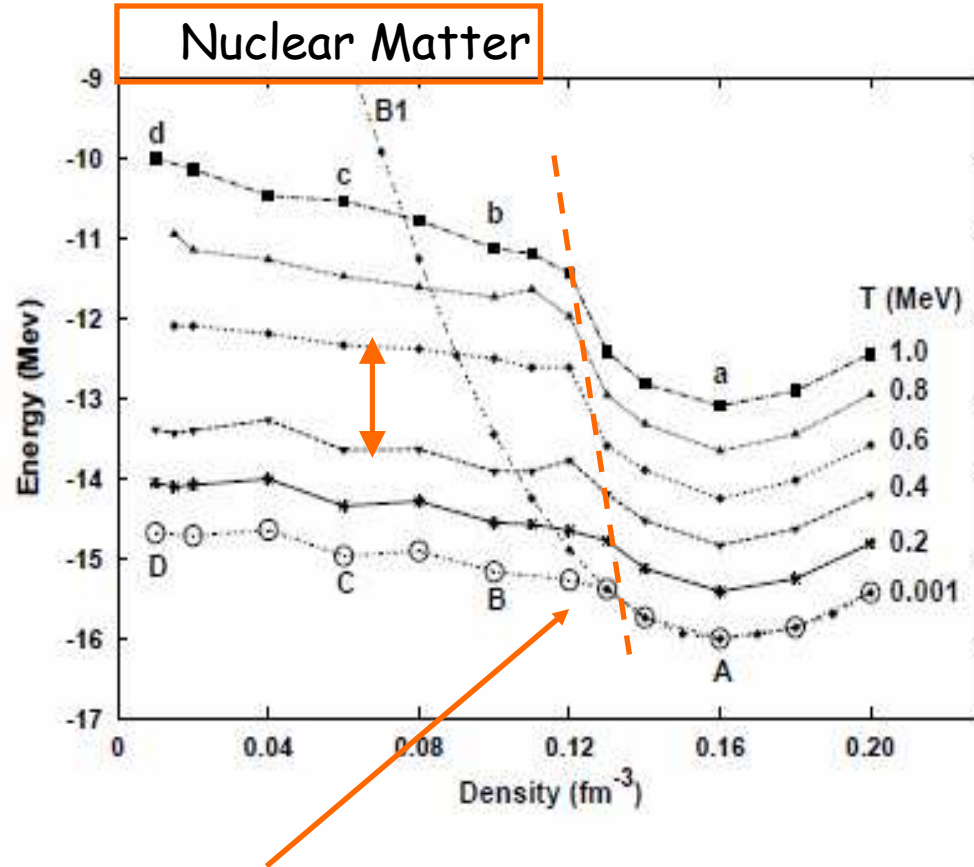


FIG. 3. Binding energies per nucleon for systems obtained with the Pandharipande medium potential at the listed temperatures.

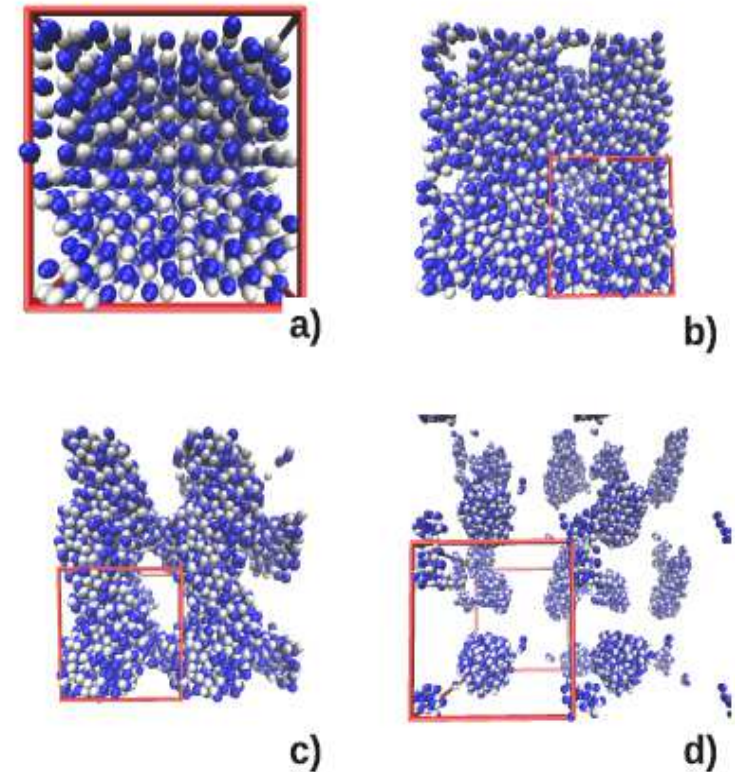
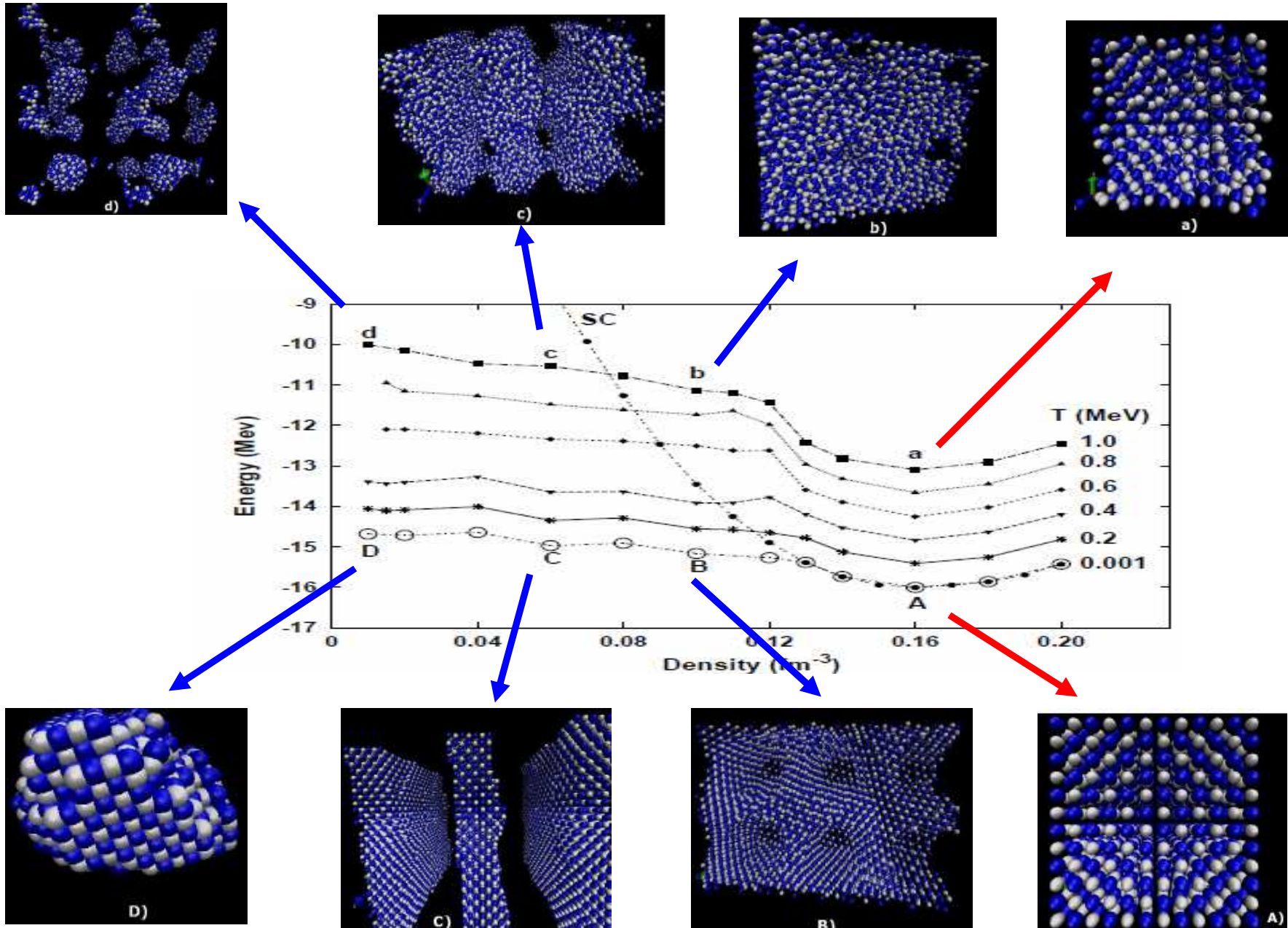


FIG. 4. (Color online) Structures corresponding to the labeled points of Figure 3 obtained with the Pandharipande medium potential at $T = 1.0 \text{ MeV}$.

Low T structures



Illinois potential + screened Coulomb NS Matter?

Ahora exploramos:

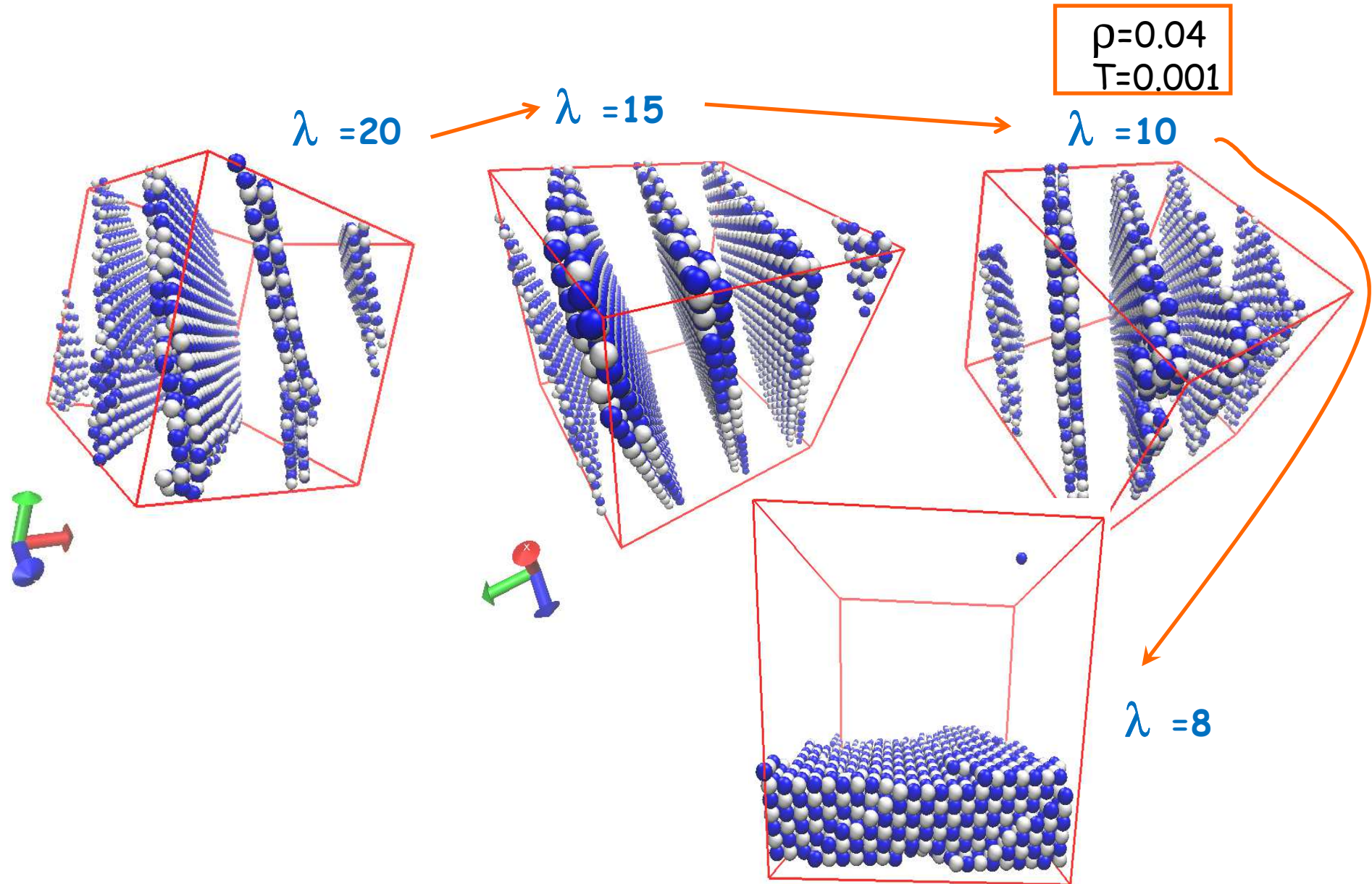
El efecto de variar λ

El efecto de variar T

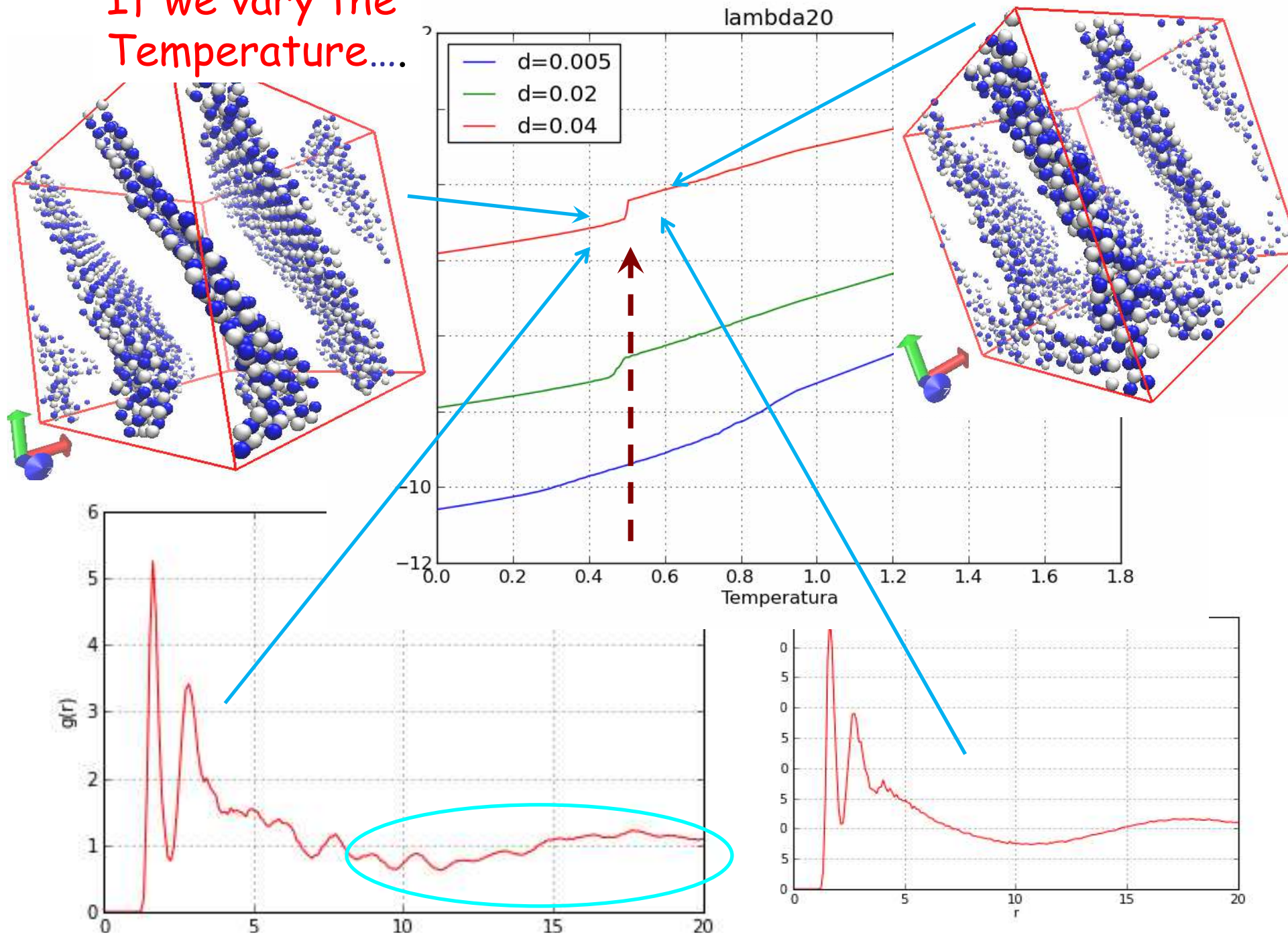
Clusters como funcion de T



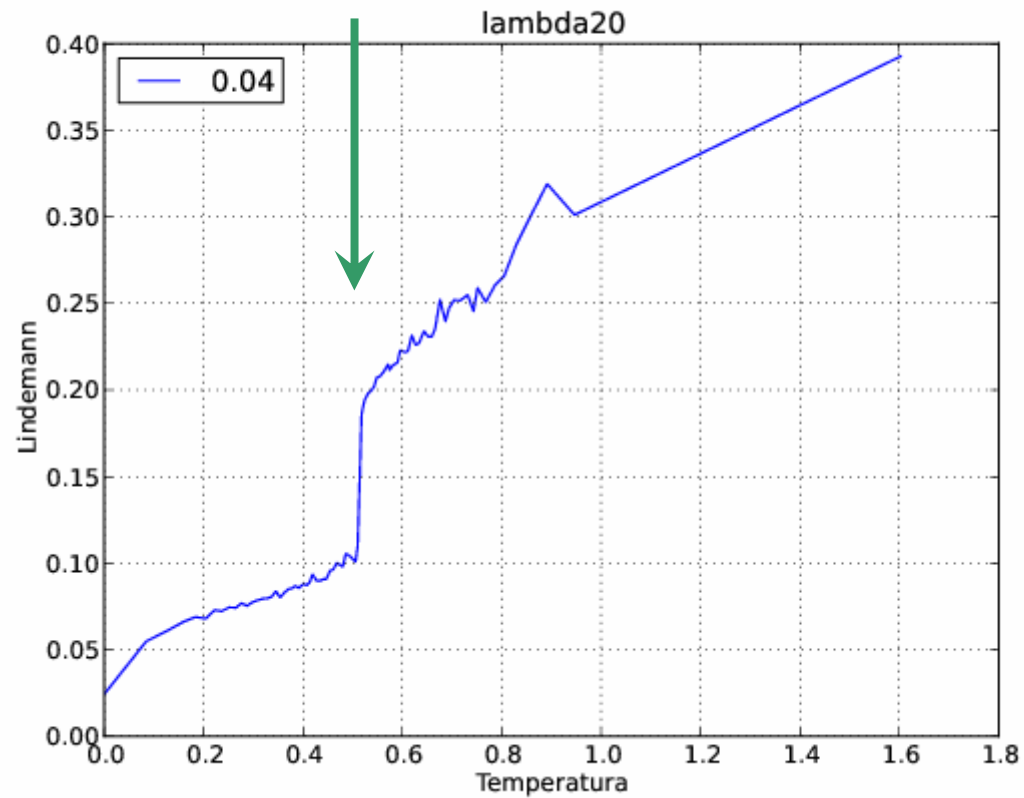
As before we fix the density and then we vary λ in order to see at which point the solution goes to a single structure per cell



If we vary the Temperature....



We now calculate the Lindemann coefficient



$$\Delta_L = \frac{1}{a'} \sqrt{\sum_i \left\langle \frac{\Delta r_i^2}{N} \right\rangle}$$

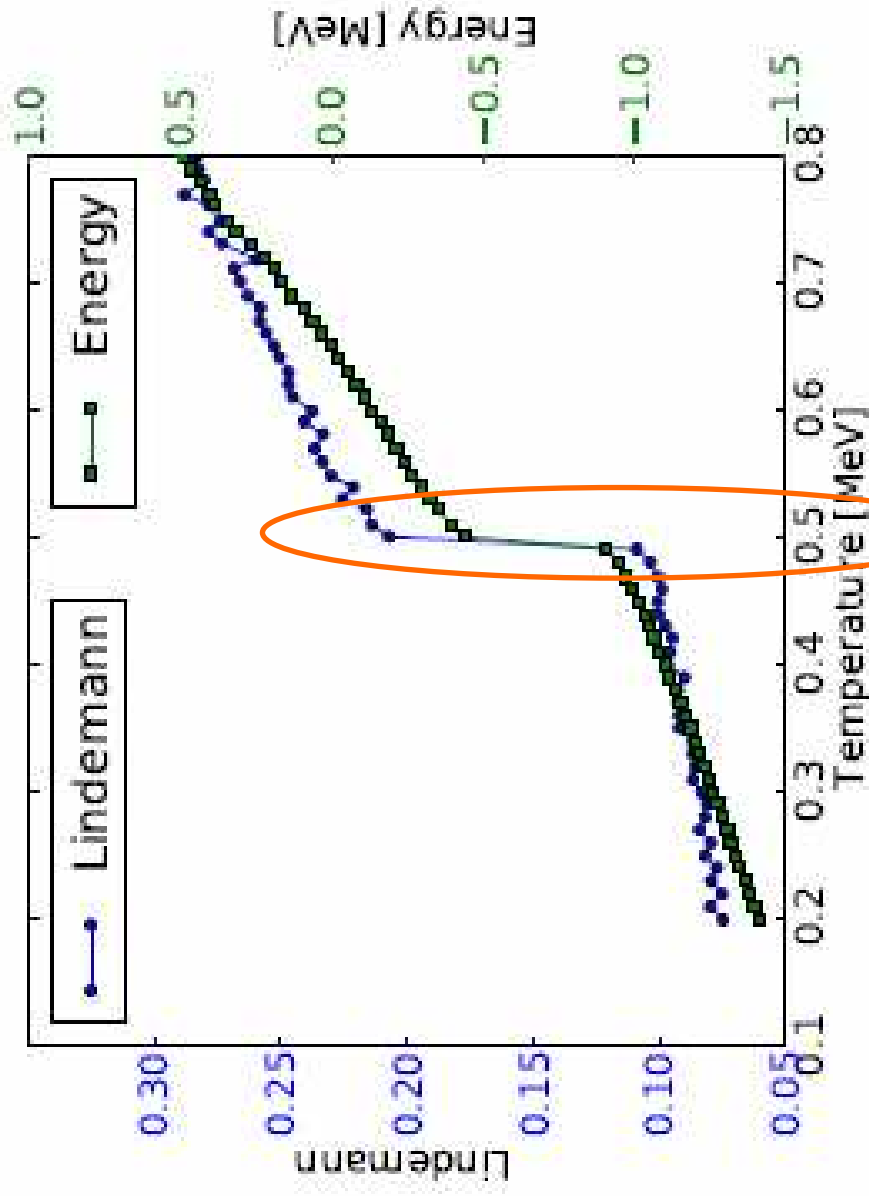
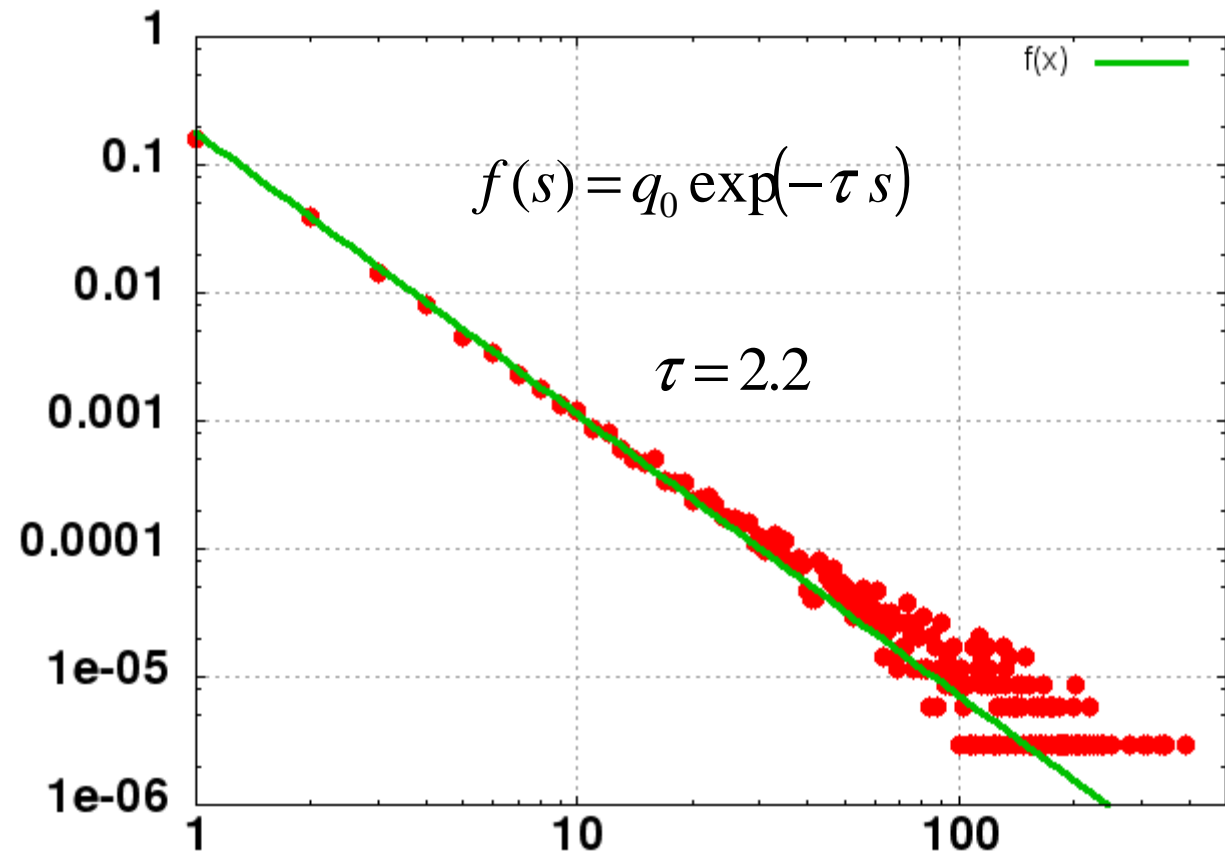
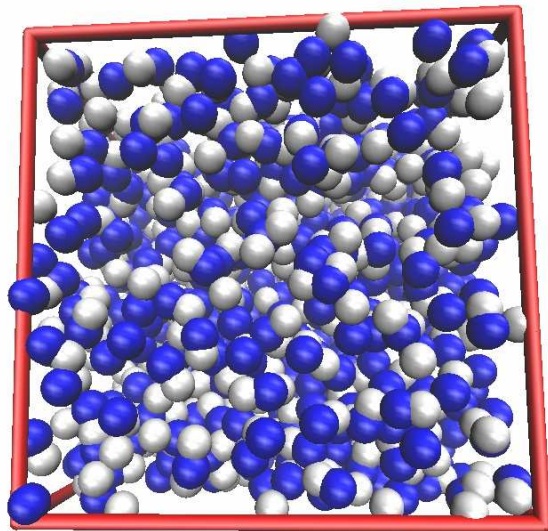


FIG. 2: Lindemann coefficient and energy as a function of temperature for a chosen density, $\rho = 0.05 \text{ fm}^{-3}$. The sudden change in their value is a signal of a solid-liquid phase transition. We can see that both discontinuities are at the same temperature

Clusters in 3D with T

T=4.5 MeV

MSTE(rc=5.4)



Cuando se "enciende" Coulomb y si $\lambda > \lambda_c$ aparecen
Las verdaderas pastas
Múltiples estructuras en una celda
El potencial fija la escala

Transiciones de Fase "dentro de las pastas" "sólido - líquido"

Clusters in Energy Space

Opacidad

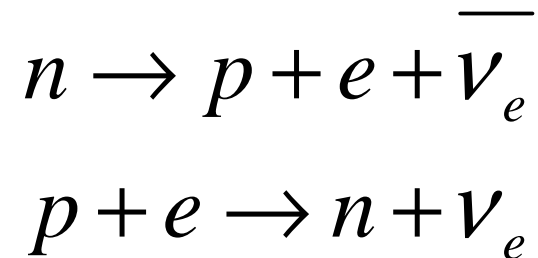
y

enfriamiento NS



Neutrino production in Neutron Stars

URCA process



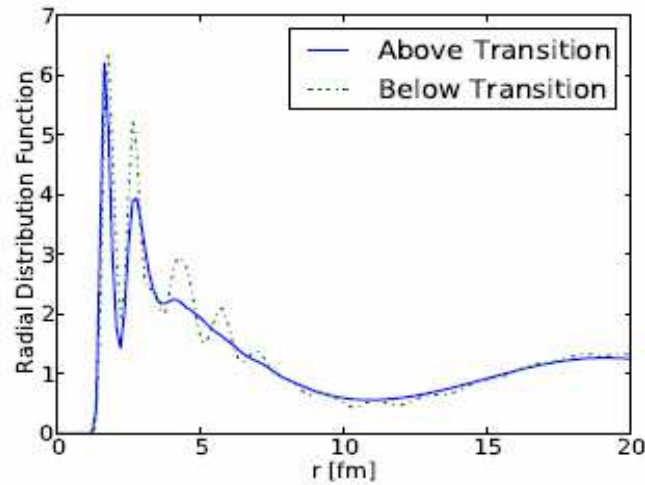
$$\lambda_\nu \approx 15 \text{ fm}$$

Correlación radial

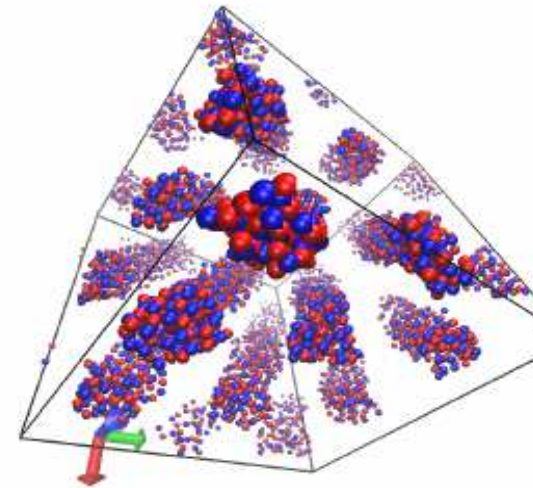
Structure Factor

$$S(q) = 1 + \rho \int_V dr e^{-i q r} g(r)$$

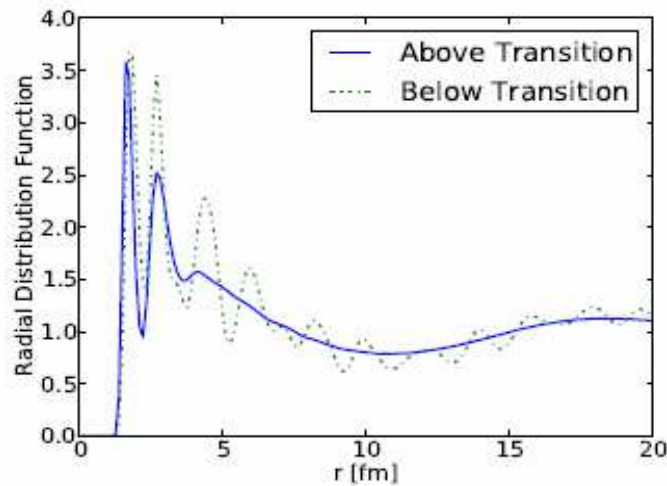
La función de correlación radial y la Temperatura



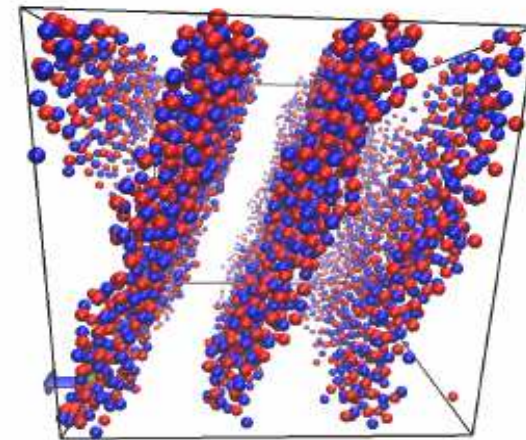
Radial distribution function for $\rho = 0.03 \text{ fm}^{-3}$



Snapshot of the system in the liquid phase for $\rho = 0.03 \text{ fm}^{-3}$

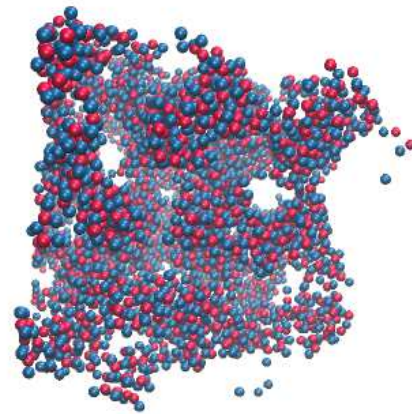


Radial distribution function for $\rho = 0.05 \text{ fm}^{-3}$

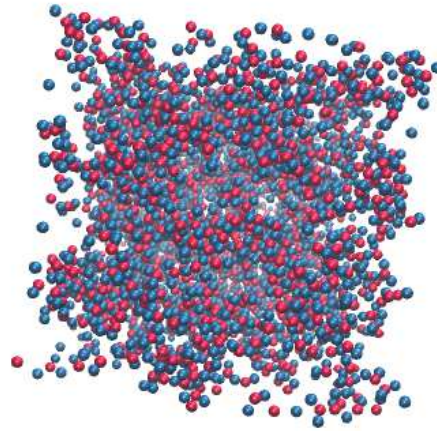


Snapshot of the system in the liquid phase for $\rho = 0.05 \text{ fm}^{-3}$

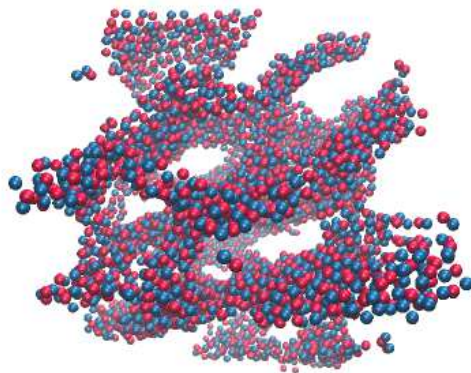
Morfología de la "pasta" y la Temperatura



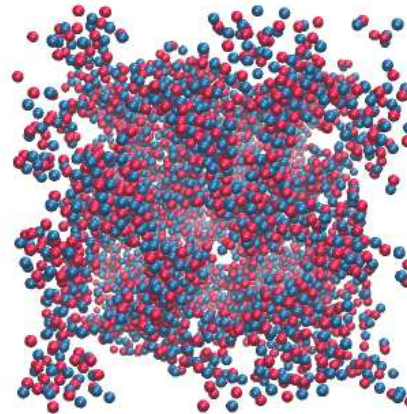
(a) $x = 0.4, T = 0.5 \text{ MeV}$



(b) $x = 0.4, T = 1.0 \text{ MeV}$



(c) $x = 0.5, T = 0.5 \text{ MeV}$



(d) $x = 0.5, T = 1.0 \text{ MeV}$

Figure 1: (Color online) Snapshots of a system with density $\rho = 0.04 \text{ fm}^{-3}$ for different values of proton fraction and temperature, generated with VMD [32]. Structures obtained at $T = 0.5 \text{ MeV}$ differ substantially. Nevertheless both show inhomogeneities.

Maximo Fragmento, X y la Temperatura (MSTE)

celda de ≈ 5500 particulas

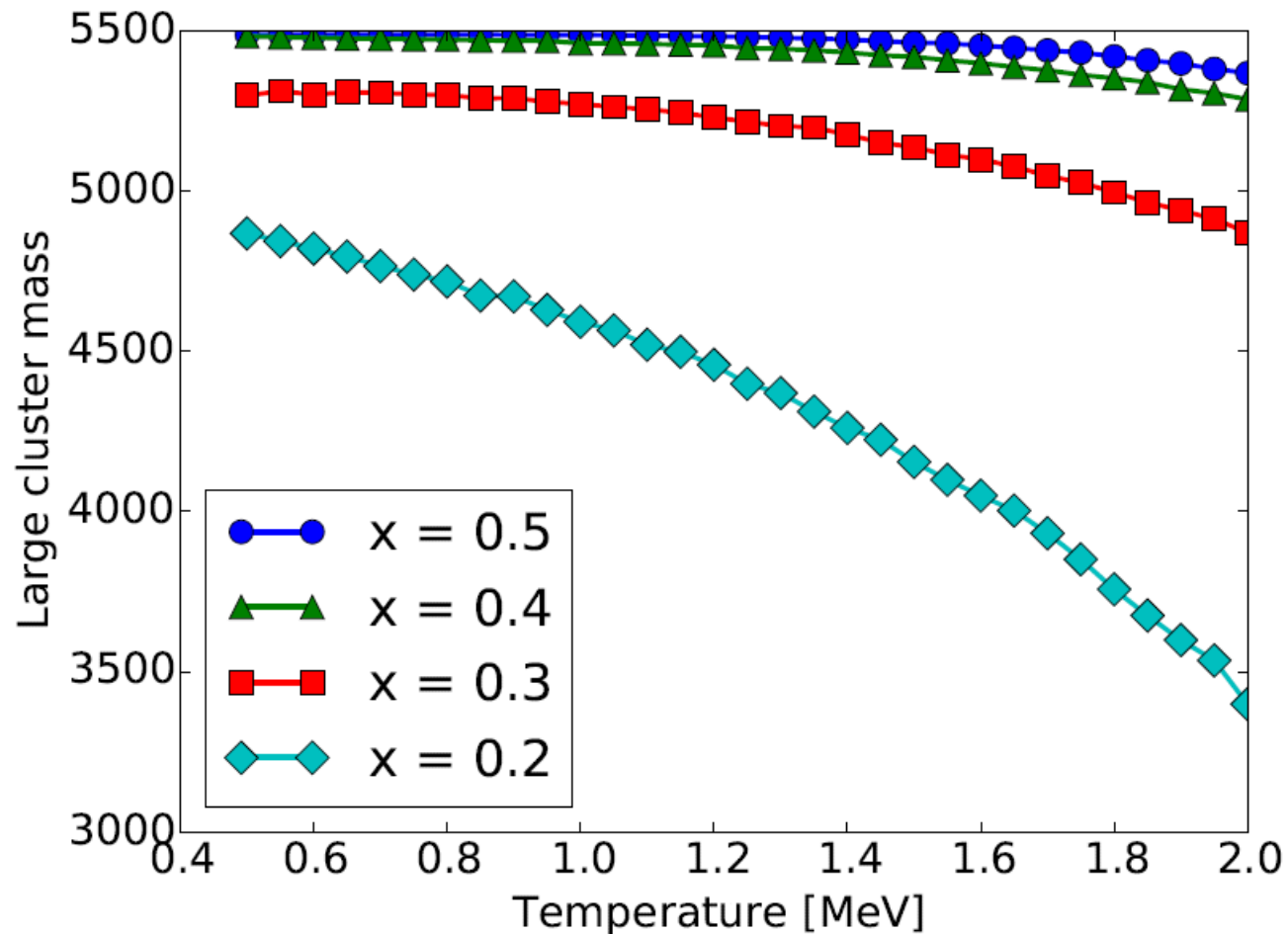
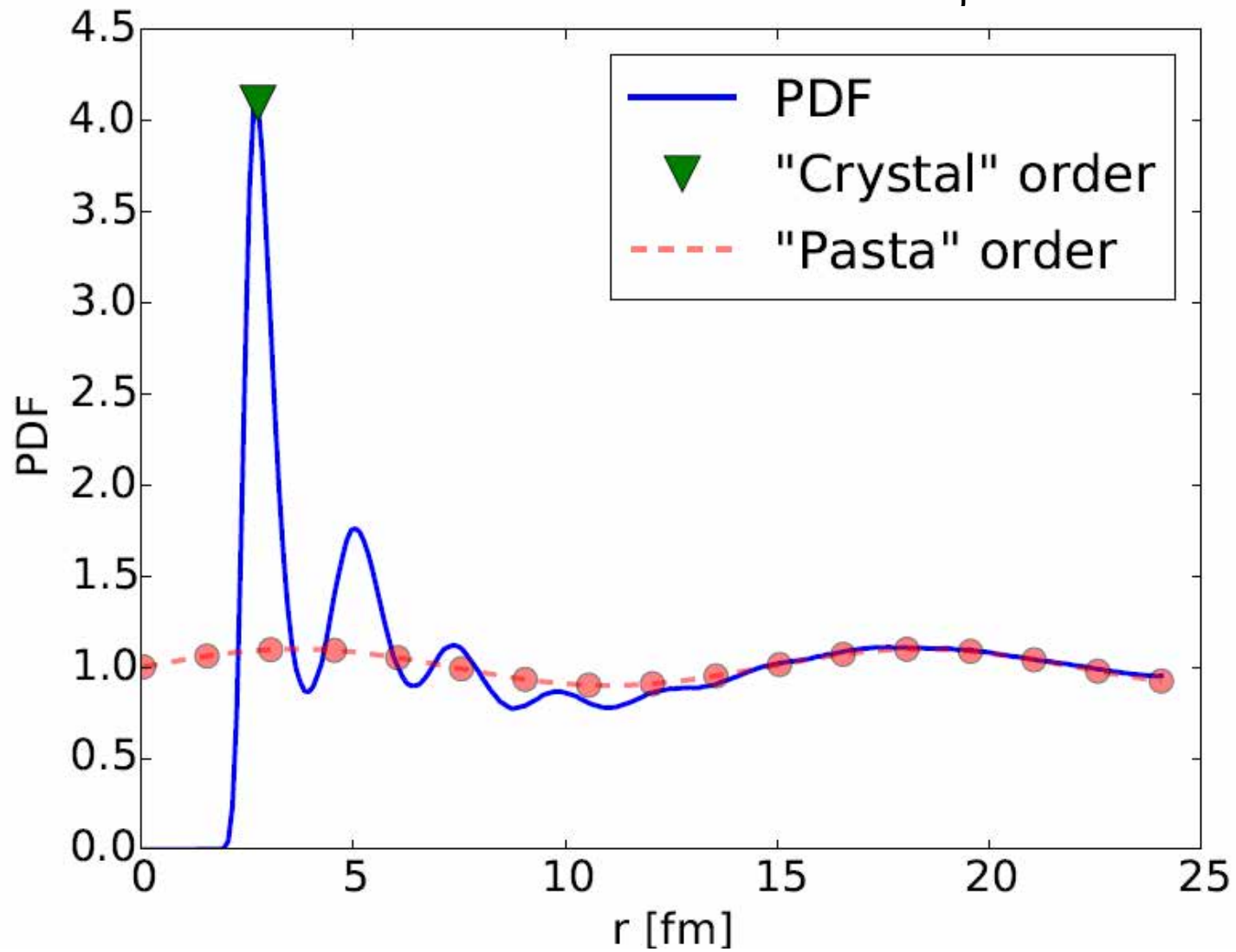


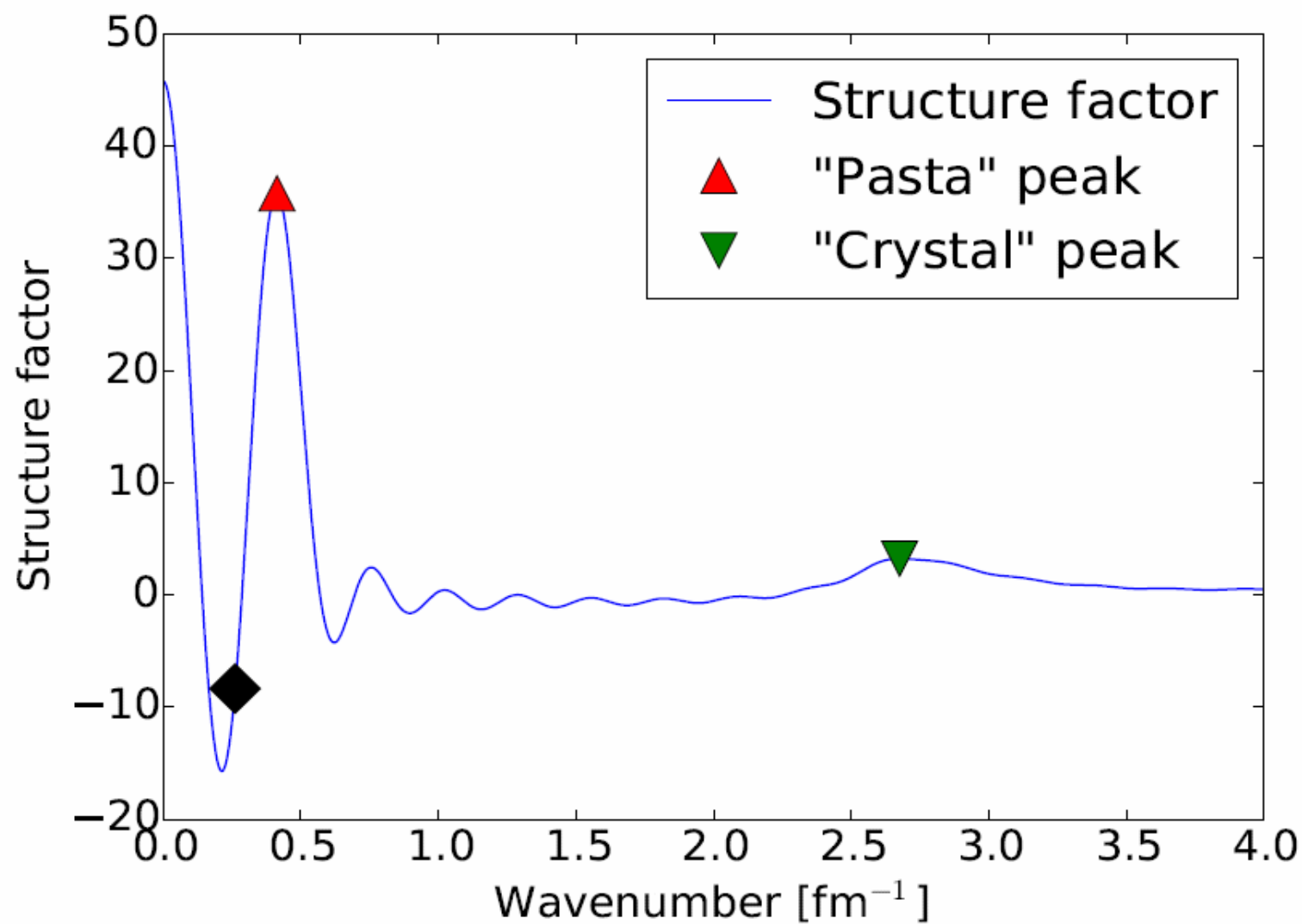
Figure 3: (Color online) Mass of the largest cluster for $\rho = 0.04 \text{ fm}^{-3}$ for different values of x .

La función de correlación radial

$X=0.5$
 $T=0.5$
 $\rho=0.05$

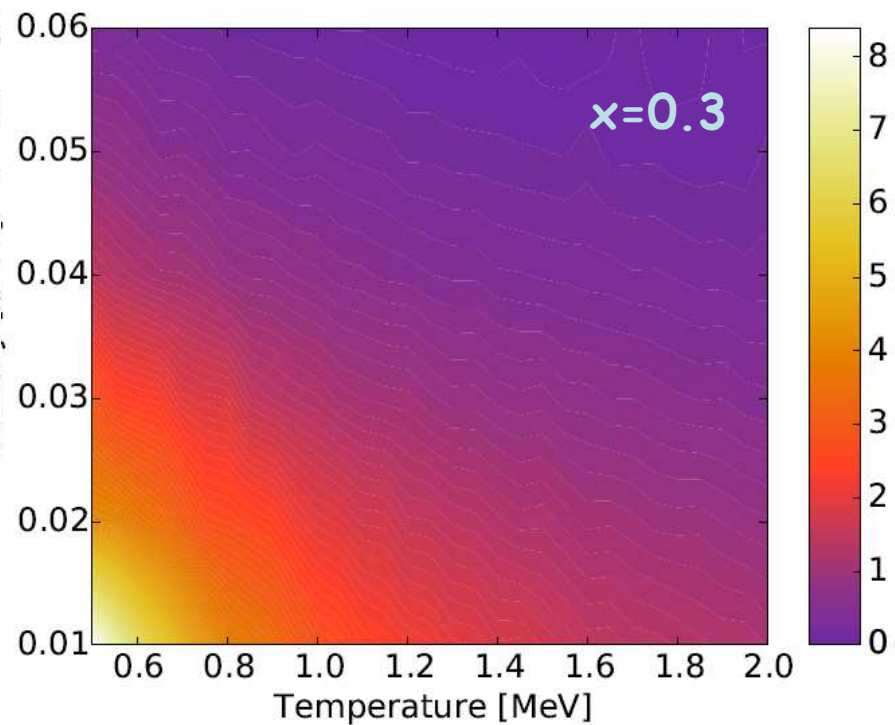
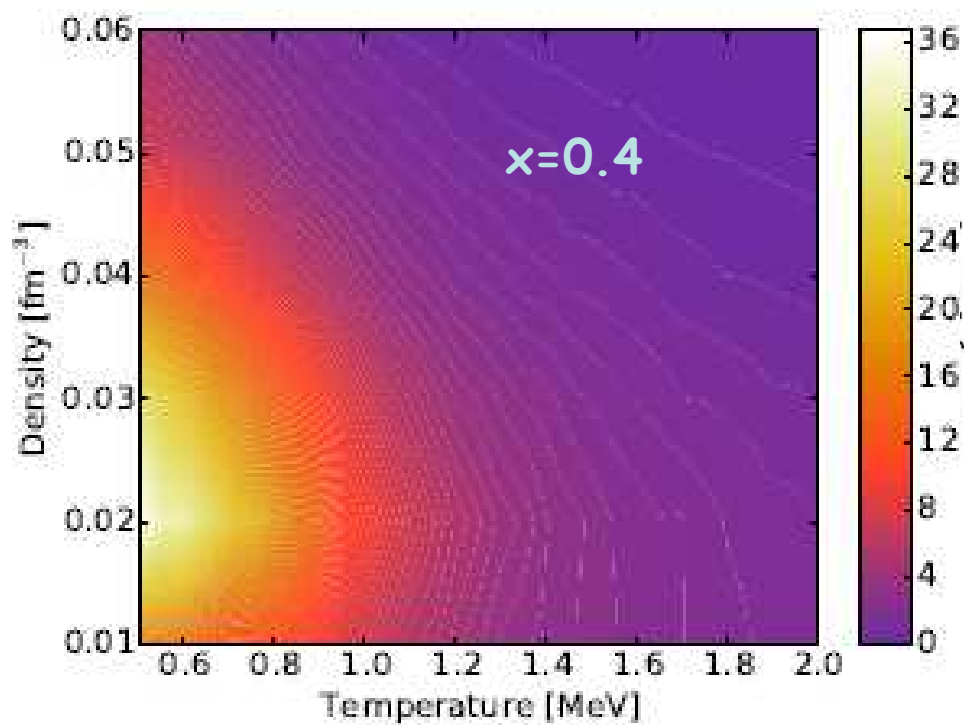
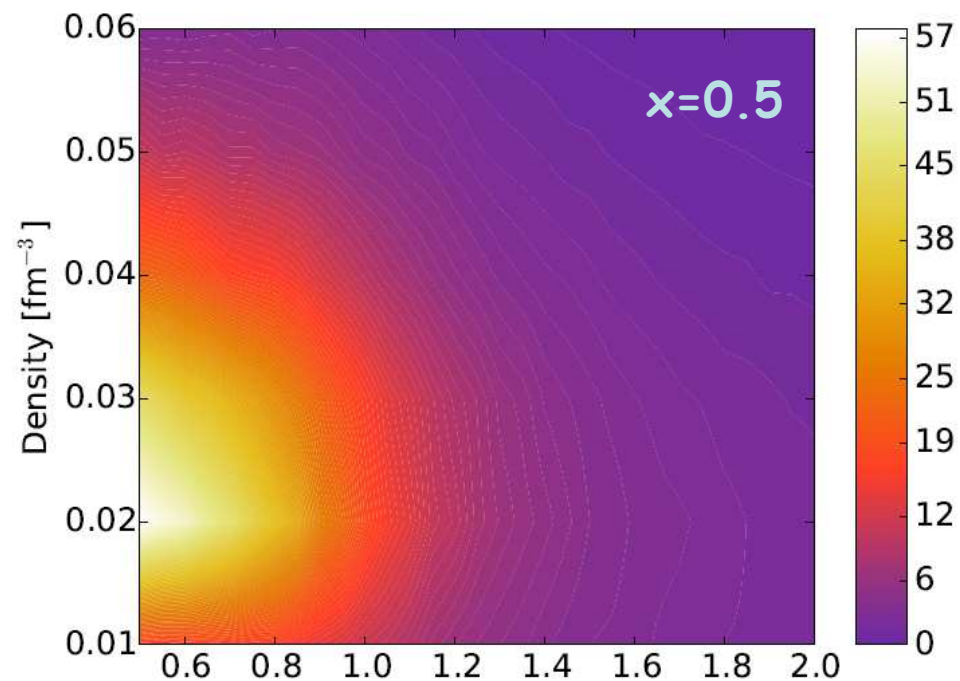


el factor de estructura



Absorption Peak

$$\lambda_{\nu_U}$$



Neutron Star Mergers

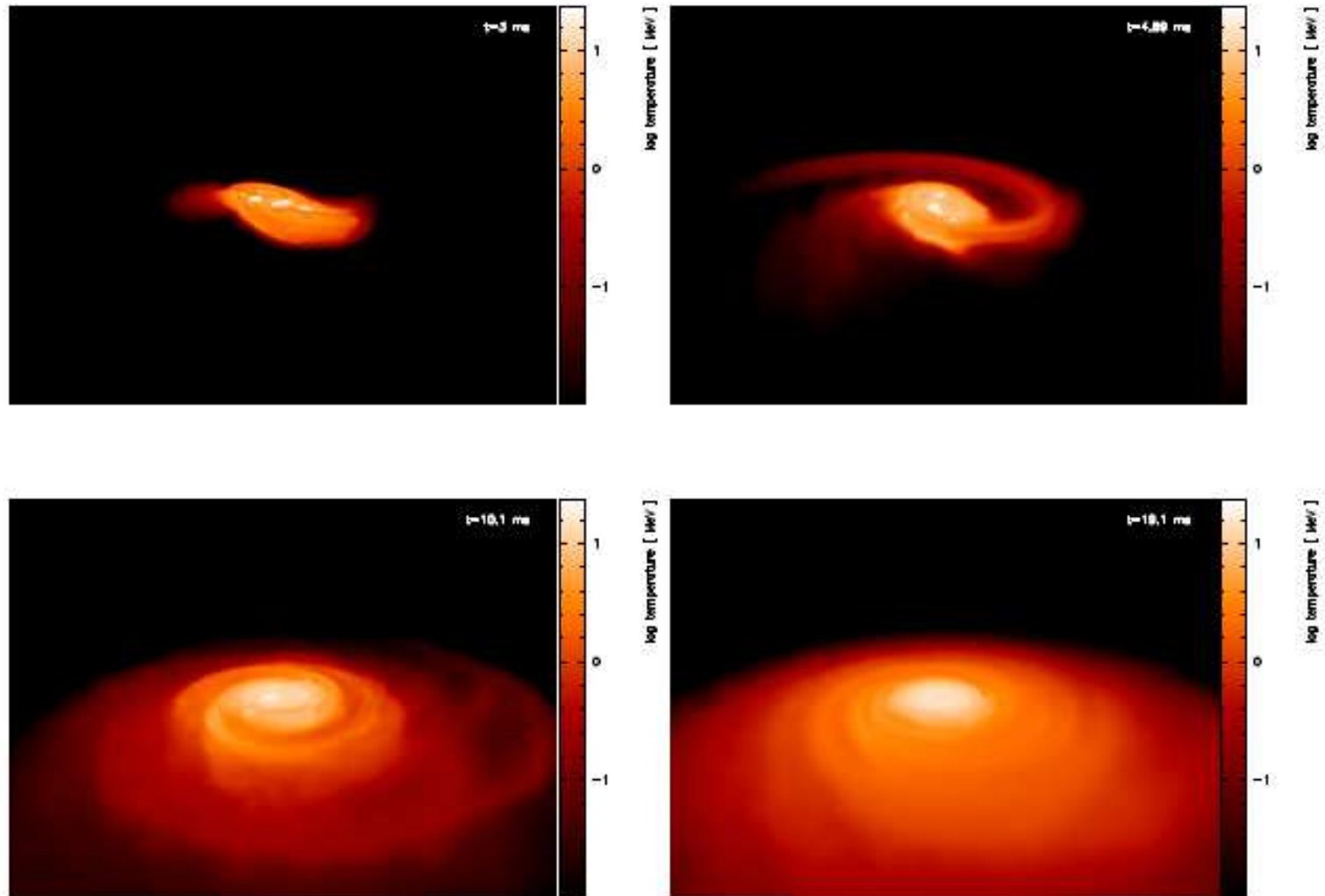
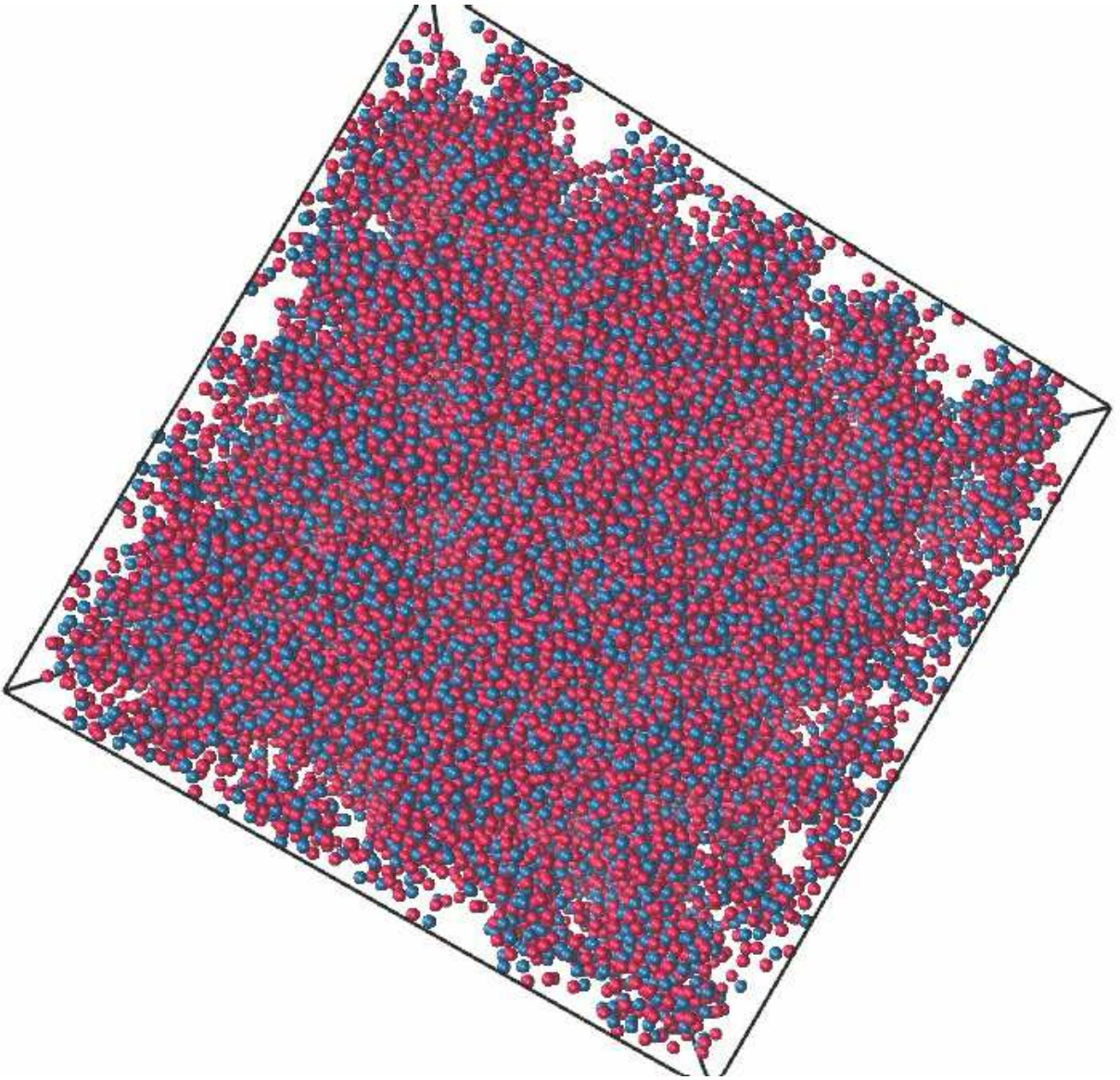
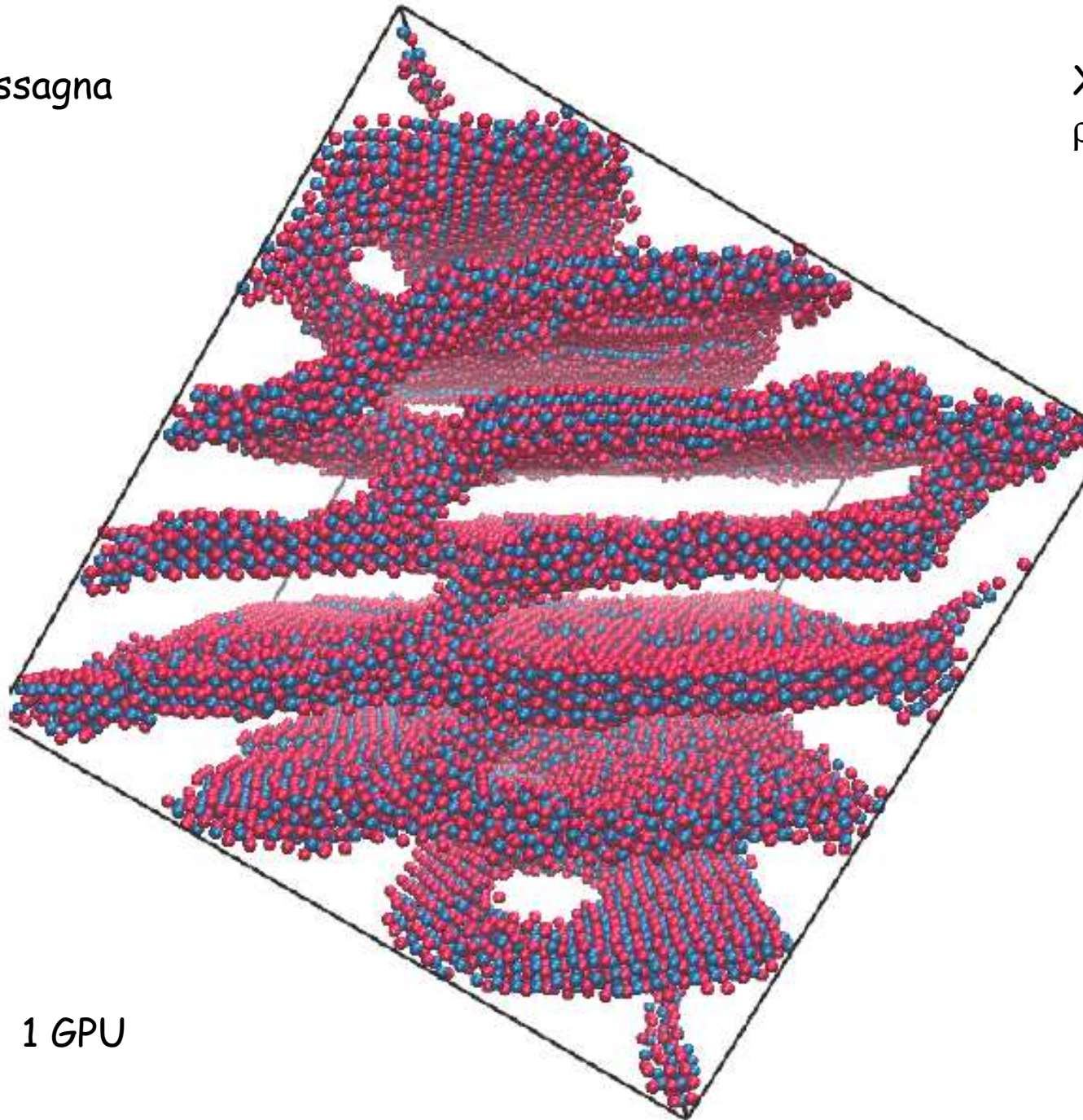


Figure 2. 3D rendering of the temperature distribution in the standard neutron star merger case (1.3 and $1.4 M_{\odot}$, no spin; run H). The upper half of the matter distribution has been "chopped off" to allow for a view into the stars. To enhance the contrast, the upper limit of the colourbar has been fixed to 20 MeV. In the various vortices that emerge due to Kelvin-Helmholtz instabilities peak temperatures in excess of 60 MeV are temporarily reached.



cuasi Lassa

$X=0.4$
 $\rho=0.04$



$N=51200!$ 1 GPU

Conclusiones :

- Sistemas con fuerzas competitivas sobrellevan una transición a Pastas a Temperaturas bajas
- Sistemas con interacciones del tipo $hc +$ atracción de corto rango desarrollan pseudo pastas (1 por celda)
- Estos sistemas son apropiadamente descritos vía
Minkowski functionals
 $g''(r) - S(q)$
Distribuciones de masa
- Para términos de Coulomb+debye screening debajo de λ_c pasa a 1 por celdas
- A baja temperatura las pastas son cristalinas
- Al subir la temperatura las pastas pierden su orden interno
- A temperaturas mayores aparecen pastas no tradicionales
- En todos los casos neutrino opacity, que disminuye con la temperatura
- Ambiente propicio para r-process?

Bibliografia Neutron Star unicamente

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"Neutron Star Crust" Editors: C.A. Bertulani and J. Piekarewicz

Nova Science Publishers, Inc. 2012

- "Topological characterization of neutron stars crust"

C.O.Dorso P.Gimenez Molinelli, & J.lopez PRC 86, 055805 (2012)

Physics (APS spotlighting exceptional research) synopses 2012

Physical Review C editor suggestion

- "Simulations of cold nuclear matter at sub-saturation densities"

P. A. Gimenez Molinelli, J. I. Nichols, J Lopez & C.O.Dorso Nucl.Phys.A 923(2014) 31

- "Effect of Coulomb screening length on nuclear ``pasta'' simulations"

P. N. Alcain, P. A. Gimenez Molinelli, J. I. Nichols, & C.O.Dorso Phys Rev C 89 (2014) 055801

- "Finite size effects in Neutron Star and Nuclear matter simulations"

P. A. Gimenez Molinelli & C.O.Dorso Nucl.Phys.A 933 (2015)306-324

- "Beyond Nuclear Pasta: Phase Transitions and neutrino opacity in non traditional Nuclear Pasta"

P.N. Alcain, P.Gimenez Molinelli & C.O.Dorso Phys.Rev.C90 065803 (2014)

- Neutron Star Opacity and Proton Fraction

P.N. Alcain & C.O.Dorso - arXiv:1412.6465, 2014 - arxiv.org

- Fragment production in expanding NS matter.

P.N. Alcain & C.O.Dorso en preparacion



Thank you